

# Neutron star constraints from theory, experiment and observations

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U.S. DEPARTMENT OF  
**ENERGY**

Office of Science

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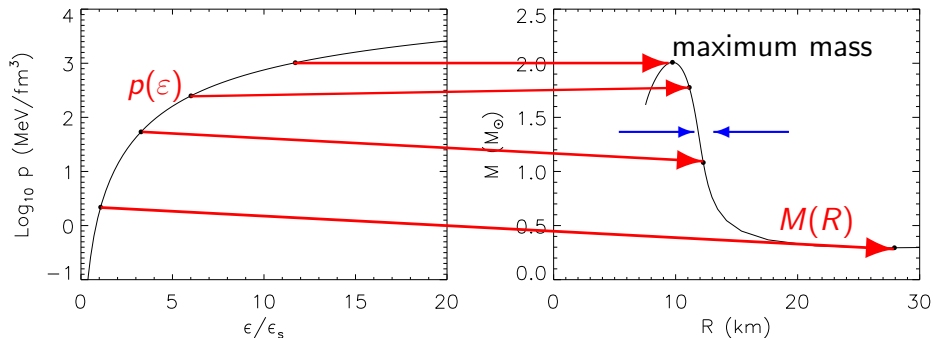
27 May, 2015, Thessaloniki  
Workshop on Binary Neutron Star Merges

- ▶ General Relativity Constraints on Neutron Star Structure
- ▶ The Neutron Star Radius and the Nuclear Symmetry Energy
- ▶ Nuclear Experimental Constraints on the Symmetry Energy
- ▶ Constraints from Pure Neutron Matter Theory
- ▶ Quark Matter in Neutron Stars
- ▶ Astrophysical Constraints on Masses and Radii

# Neutron Star Structure

## Tolman-Oppenheimer-Volkov equations

$$\frac{dp}{dr} = -\frac{G}{c^4} \frac{(mc^2 + 4\pi pr^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$
$$\frac{dm}{dr} = 4\pi \frac{\epsilon}{c^2} r^2$$



Equation of State



Observations

# Neutron Star Structure

Newtonian Gravity:

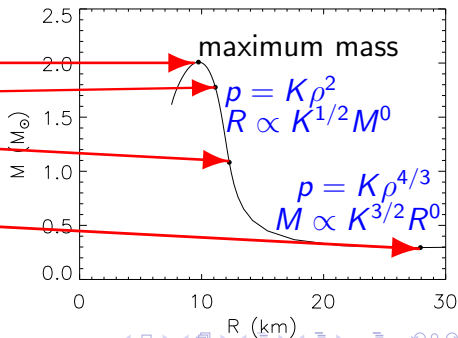
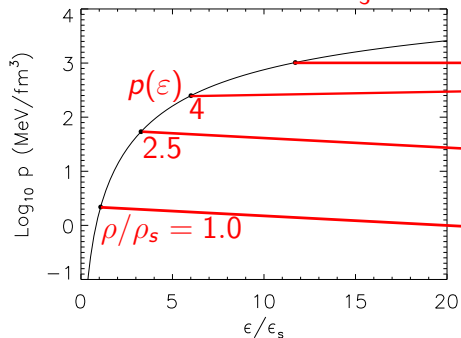
$$\frac{dp}{dr} = -\frac{Gm\rho}{r^2}; \quad \frac{dm}{dr} = 4\pi r^2 \rho; \quad \rho c^2 = \varepsilon$$

Newtonian Polytrope:

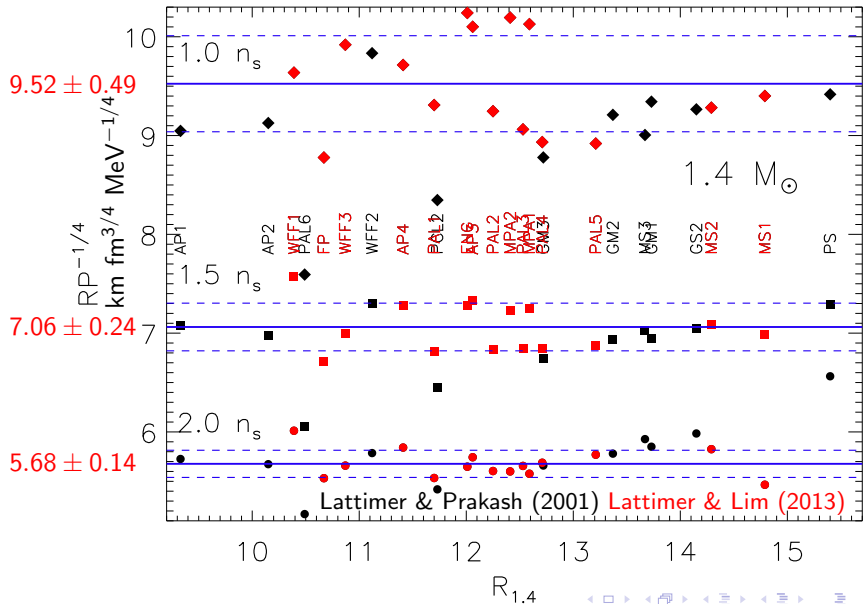
$$p = K\rho^\gamma; \quad M \propto K^{1/(2-\gamma)} R^{(4-3\gamma)/(2-\gamma)}$$

$$\rho < \rho_s: \gamma \simeq \frac{4}{3};$$

$$\rho > \rho_s: \gamma \simeq 2$$



# The Radius – Pressure Correlation



# Nuclear Symmetry Energy and Pressure

Defined as the difference between energies of pure neutron matter ( $x = 0$ ) and symmetric ( $x = 1/2$ ) nuclear matter.

$$S(\rho) = E(\rho, x = 0) - E(\rho, x = 1/2)$$

Expanding around the saturation density ( $\rho_s$ ) and symmetric matter ( $x = 1/2$ )

$$E(\rho, x) = E(\rho, 1/2) + (1-2x)^2 S_2(\rho) + \dots$$

$$S_2(\rho) = S_v + \frac{L}{3} \frac{\rho - \rho_s}{\rho_s} + \dots$$

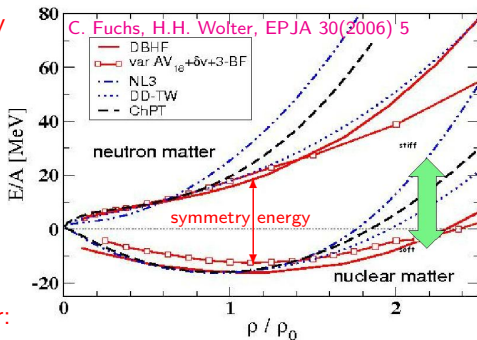
$$S_v \simeq 31 \text{ MeV}, \quad L \simeq 50 \text{ MeV}$$

Connections to pure neutron matter:

$$E(\rho_s, 0) \approx S_v + E(\rho_s, 1/2) \equiv S_v - B, \quad \rho(\rho_s, 0) = L\rho_s/3$$

Neutron star matter (in beta equilibrium):

$$\frac{\partial(E + E_e)}{\partial x} = 0, \quad \rho(\rho_s, x_\beta) \simeq \frac{L\rho_s}{3} \left[ 1 - \left( \frac{4S_v}{\hbar c} \right)^3 \frac{4 - 3S_v/L}{3\pi^2 \rho_s} \right]$$



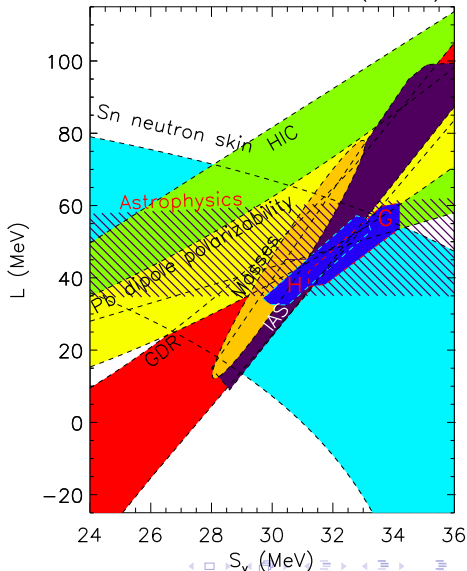
# Experimental and Neutron Matter Constraints

Gandolfi, Carlson & Reddy (2011);  
Hebeler & Schwenk (2011)

H&S: Chiral Lagrangian

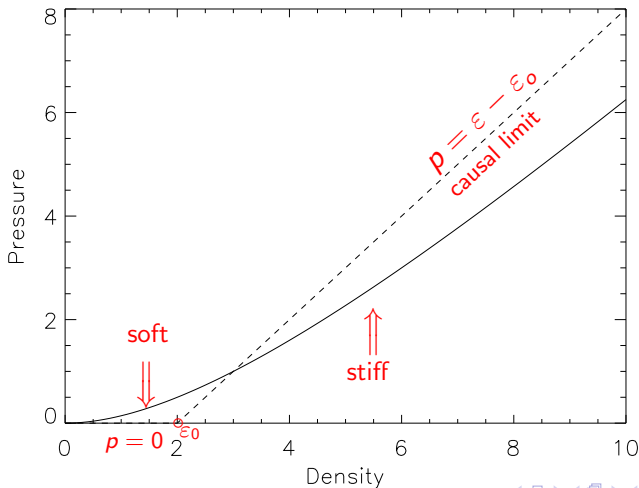
GC&R: Quantum Monte Carlo

$S_v - L$  constraints from  
Hebeler et al. (2012)



# Extremes of Compaction of Neutron Stars

- ▶ The most compact and massive configurations occur when the low-density equation of state is "soft" and the high-density equation of state is "stiff" (Koranda, Stergioulas & Friedman 1997).



$\varepsilon_0$  is the only  
EOS parameter

The TOV  
solutions scale  
with  $\varepsilon_0$

$$w = \varepsilon/\varepsilon_0$$

$$y = p/\varepsilon_0$$

$$x = r\sqrt{G\varepsilon_0}/c^2$$

$$z = m\sqrt{G^3\varepsilon_0}/c^2$$



# Extremal Properties of Neutron Stars

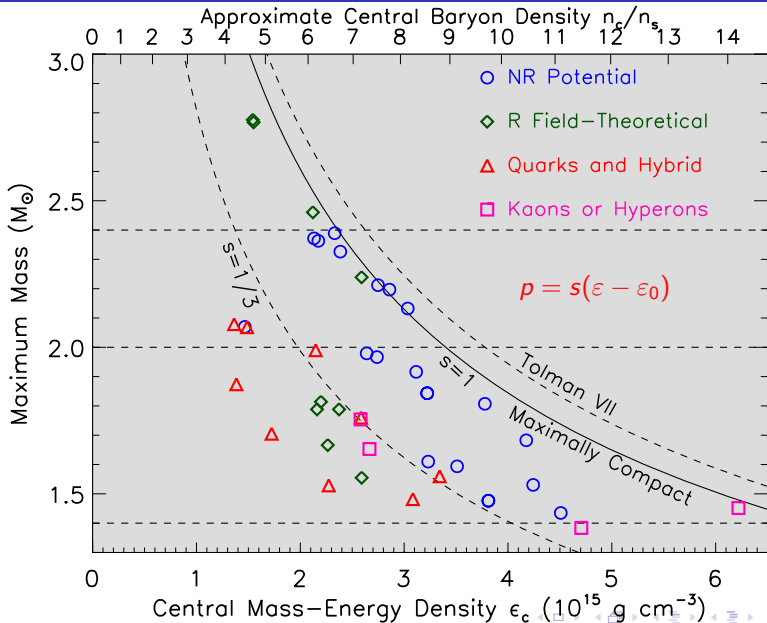
The maximum mass configuration is achieved when  
 $x_R = 0.2404$ ,  $w_c = 3.034$ ,  $y_c = 2.034$ ,  $z_R = 0.08513$ .

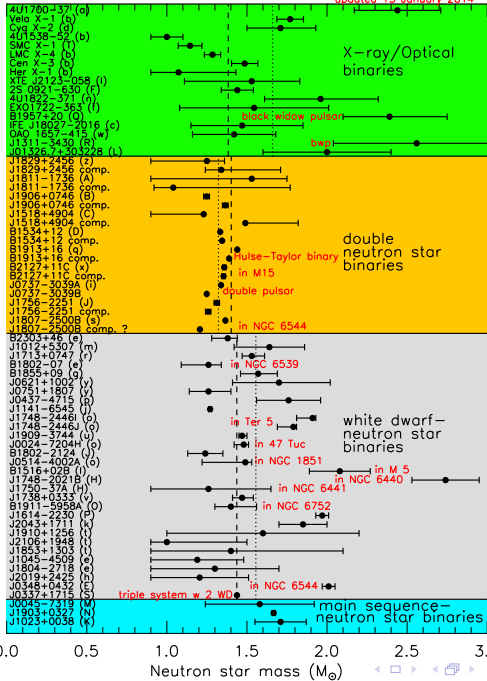
A useful reference density is the nuclear saturation density  
(interior density of normal nuclei):

$$\rho_s = 2.7 \times 10^{14} \text{ g cm}^{-3}, \quad n_s = 0.16 \text{ baryons fm}^{-3}, \quad \varepsilon_s = 150 \text{ MeV fm}^{-3}$$

- ▶  $M_{\max} = 4.1 (\varepsilon_s/\varepsilon_0)^{1/2} M_\odot$  (Rhoades & Ruffini 1974)
- ▶  $M_{B,\max} = 5.41 (m_B c^2/\mu_o)(\varepsilon_s/\varepsilon_0)^{1/2} M_\odot$
- ▶  $R_{\min} = 2.82 GM/c^2 = 4.3 (M/M_\odot) \text{ km}$
- ▶  $\mu_{b,\max} = 2.09 \text{ GeV}$
- ▶  $\varepsilon_{c,\max} = 3.034 \varepsilon_0 \simeq 51 (M_\odot/M_{\text{largest}})^2 \varepsilon_s$
- ▶  $p_{c,\max} = 2.034 \varepsilon_0 \simeq 34 (M_\odot/M_{\text{largest}})^2 \varepsilon_s$
- ▶  $n_{B,\max} \simeq 38 (M_\odot/M_{\text{largest}})^2 n_s$
- ▶  $BE_{\max} = 0.34 M$
- ▶  $P_{\min} = 0.74 (M_\odot/M_{\text{sph}})^{1/2} (R_{\text{sph}}/10 \text{ km})^{3/2} \text{ ms} =$   
 $0.20 (M_{\text{sph,max}}/M_\odot) \text{ ms}$

# Maximum Energy Density in Neutron Stars

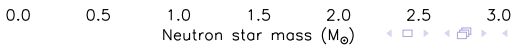




vanKerkwijk 2010  
 Romani et al. 2012

Although simple average mass of w.d. companions is  $0.23 M_{\odot}$  larger, weighted average is  $0.04 M_{\odot}$  smaller

Demorest et al. 2010  
 Antoniadis et al. 2013  
 Champion et al. 2008



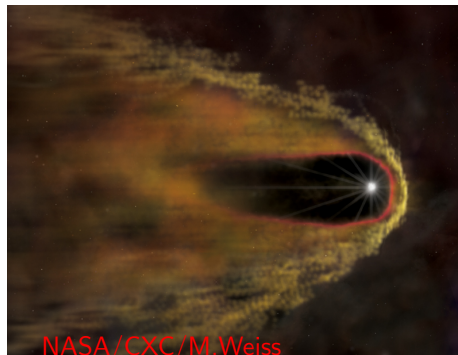
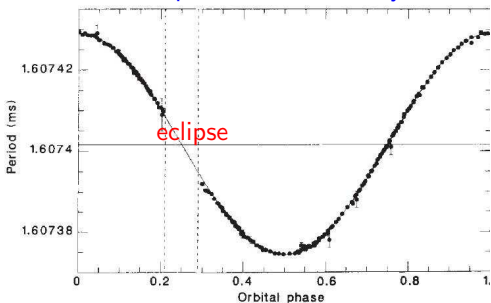
# What is the Maximum Mass?

- ▶ PSR J1614+2230 (Demorest et al. 2010)  
 $M = 1.97 \pm 0.04 M_{\odot}$ ; a nearly edge-on system with well-measured Shapiro time delay.
- ▶ PSRJ0548+0432 (Antoniadis et al. 2013)  
 $M = 2.01 \pm 0.04 M_{\odot}$ ; measured using optical data and theoretical properties of companion white dwarf.
- ▶ B1957+20 (van Kerkwijk 2010)  $M = 2.4 \pm 0.3 M_{\odot}$ ; black widow pulsar (BWP).
- ▶ PSR J1311-3430 (Romani et al. 2012)  
 $M = 2.55 \pm 0.50 M_{\odot}$ ; BWP.
- ▶ PSR J1544+4937 (Tang et al. 2014)  
 $M = 2.06 \pm 0.56 M_{\odot}$ ; BWP.
- ▶ PSR 2FGL J1653.6-0159 (Romani et al. 2014)  
 $M > f(M_2)/\sin^3 i \gtrsim 1.96 M_{\odot}$ ; largest  $f(M_2)$ .
- ▶ PSR J1227-4859 (de Martino et al. 2014)  
 $M = 2.2 \pm 0.8 M_{\odot}$ ; reback pulsar.

# Black Widow Pulsar PSR B1957+20

1.6ms pulsar in circular 9.17h orbit with  $\sim 0.03 M_{\odot}$  companion.  
Pulsar is eclipsed for 50-60 minutes each orbit; eclipsing object has a volume much larger than the companion or its Roche lobe.  
It is believed the pulsar is ablating the companion leading to mass loss and an eclipsing plasma cloud. Companion nearly fills its Roche lobe.  
Ablation by pulsar leads to eventual disappearance of companion.  
The optical light curve does not represent the center of mass of the companion, but the motion of its irradiated hot spot.

pulsar radial velocity



NASA/CXC/M.Weiss

# Black Widow Pulsar PSR B1957+20

$$K_i = 2\pi \frac{a_i \sin i}{P}$$

$$q = \frac{M_P}{M_*} = \frac{a_*}{a_P} = \frac{K_*}{K_P}$$

$$K_* = K_{obs} \left(1 + \frac{R_*}{a_*}\right)$$

Radiated area of companion has a smaller orbit than the center of mass.

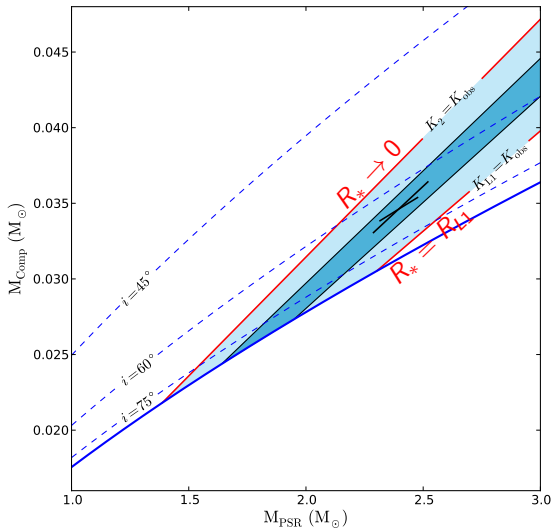
$$M_P = q(1+q)^2 \frac{P}{2\pi G} \left(\frac{K_P}{\sin i}\right)^3$$

Modeling of light curve shape suggests that

$$M_P > 1.8 M_\odot, \sin i < 66^\circ$$

Most probable values:

$$2.20 M_\odot < M_P < 2.55 M_\odot$$



van Kerkwijk 2010

# Causality + GR Limits and the Maximum Mass

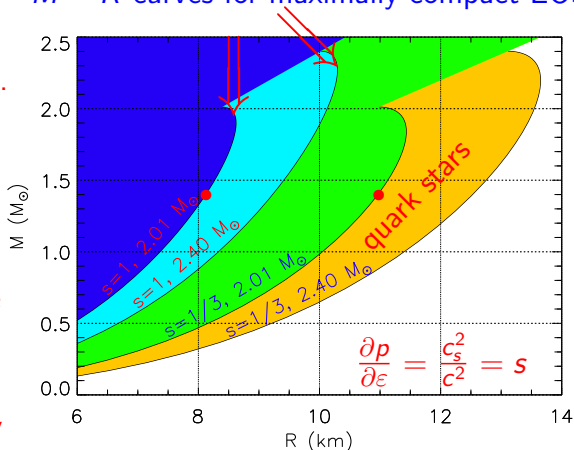
A lower limit to the maximum mass sets a lower limit to the radius for a given mass.

Similarly, a precise  $(M, R)$  measurement sets an upper limit to the maximum mass.

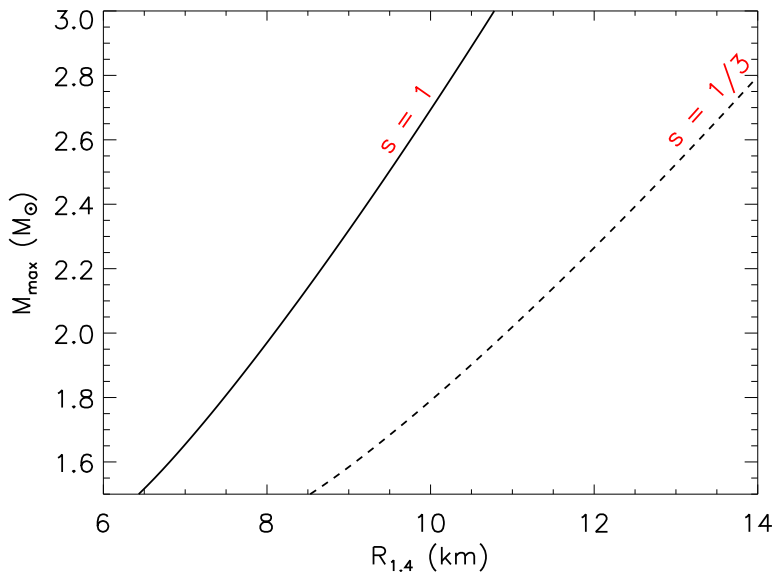
$1.4M_{\odot}$  stars must have  $R > 8.15M_{\odot}$ .

$1.4M_{\odot}$  strange quark matter stars (and likely hybrid quark/hadron stars) must have  $R > 11$  km.

$M - R$  curves for maximally compact EOS



# Maximum Mass and Neutron Star Radii





# What About Realistic EOSs?

It has been proposed that the effective sound speed limit is  $c/\sqrt{3}$  (Bedaque & Steiner 2015), in which case  $1.4M_{\odot}$  stars must have  $R_{1.4} > 11$  km.

Hybrid quark/hadron stars are realistically at least 1-2 km larger (Alford et al. 2015).

What additional constraints are imposed by our knowledge of the low-density equation of state?

# Chiral Lagrangian Neutron Matter Calculations

The study of Hebeler & Schwenk (2010) suggested moderate values  $40 \text{ MeV} < L < 60 \text{ MeV}$ , consistent with but at the lower boundary of the range favored by nuclear experiments.

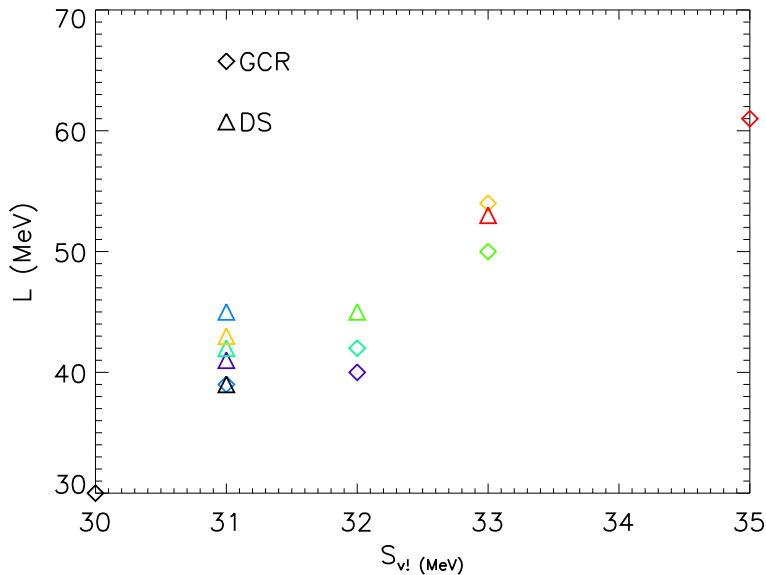
These results were in substantial agreement with the quantum Monte Carlo neutron matter calculations of Gandolfi, Carlson & Reddy (2012).

The chiral Lagrangian calculations have been refined and extended to matter with proton fractions up to and including symmetric matter (Drischler & Schwenk 2014).

The symmetry energy coefficients are found to be correlated with the saturation properties for a given parameter set.

There is a small quartic contribution to the symmetry energy.

# Neutron Matter Comparisons



# Extrapolation of Neutron Matter EOS to Higher Densities

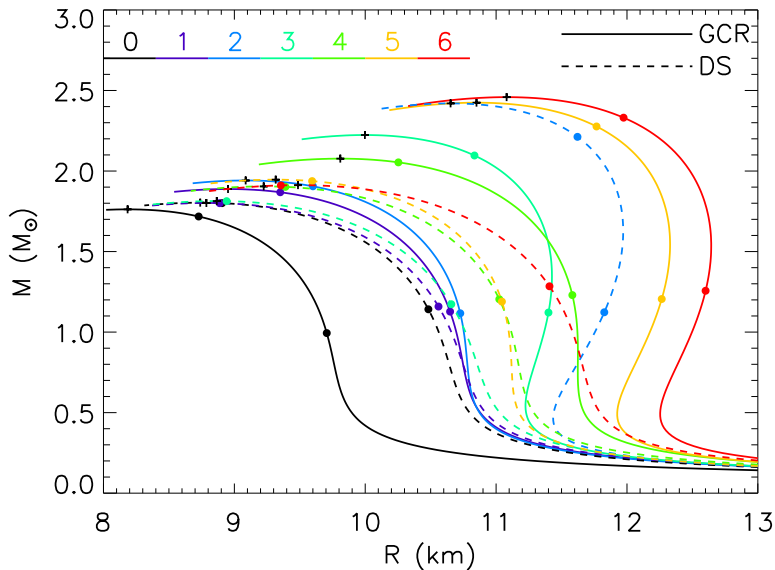
Gandolfi, Carlson & Reddy fit their QMC neutron matter equations of state to the 4-parameter fit:

$$E_n(u) = au^\alpha + bu^\beta$$

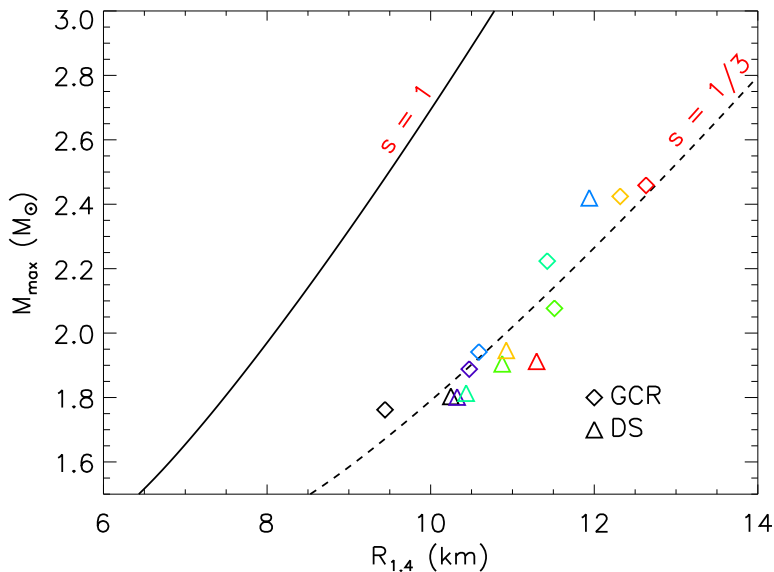
with  $u = n/n_s$  and  $n_s = 0.16 \text{ fm}^{-3}$ .

This can also be done with the chiral Lagrangian neutron matter equations of state computed by Drischler, Hebeler & Schwenk.

# Neutron Matter Extrapolations and $M - R$



# Neutron Matter Extrapolations and $M_{max} - R_{1.4}$

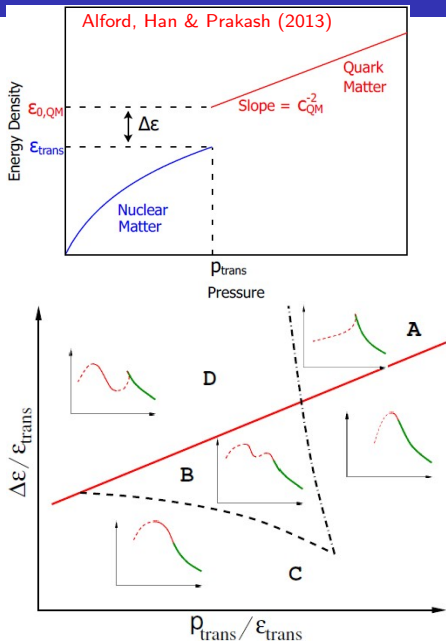


# First Order Phase Transition in Neutron Stars

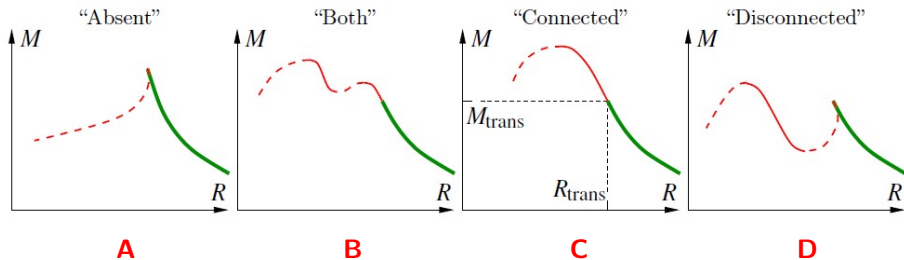
- ▶ Generic first order phase transition with 3 parameters:  $\Delta\varepsilon$ ,  $\varepsilon_t$  and  $P_t$ .
- ▶ Make 2 dimensionless parameter combinations:  $\Delta\varepsilon/\varepsilon_t$  and  $P_t/\varepsilon_t$ .
- ▶ Critical condition for existence of stable hybrid core connected to normal branch (**A**, **D**):

$$\frac{\Delta\varepsilon}{\varepsilon_t} \leq \frac{1}{2} + \frac{3}{2} \frac{P_t}{\varepsilon_t}.$$

- ▶ It is also possible to have a stable hybrid core disconnected from normal branch (**B**, **D**).
- ▶ Parametrize high-density phase with a constant sound speed  $c_{\text{QM}}^2 = dp/d\varepsilon \sim 1/3$ .



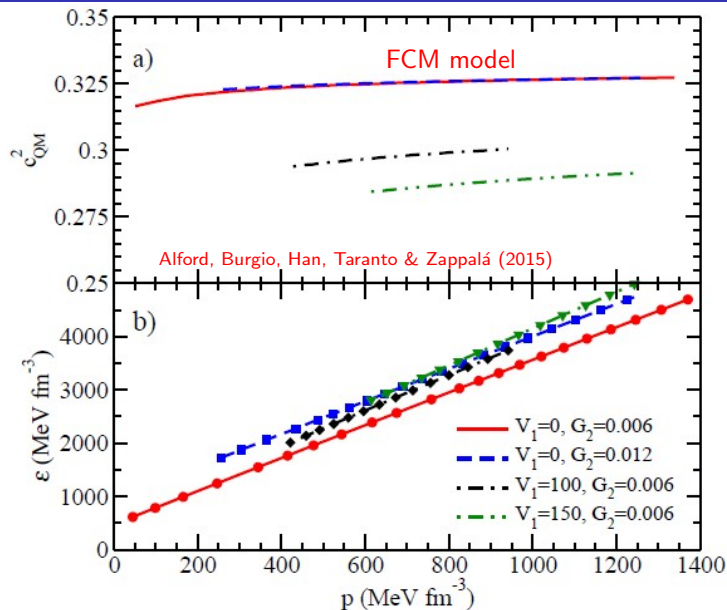
# Possible Hybrid Configurations



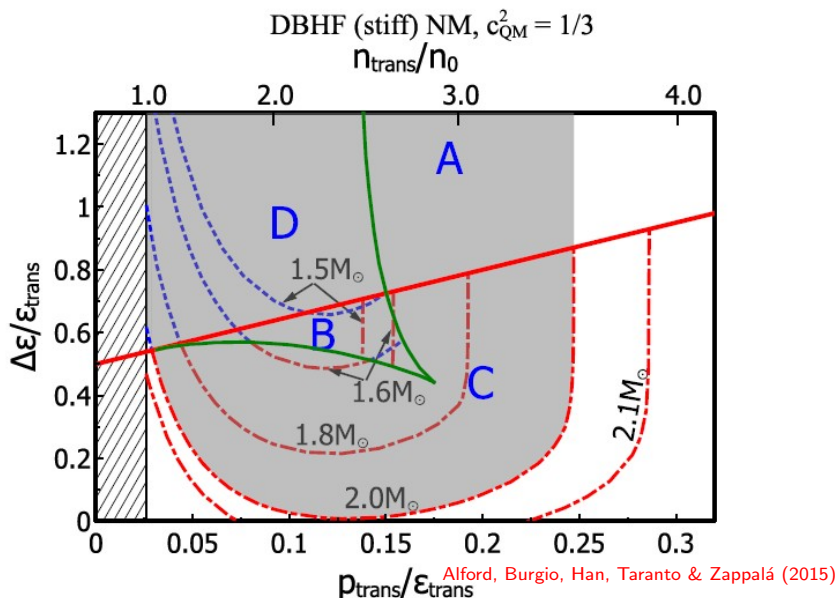
Alford, Han & Prakash (2013)



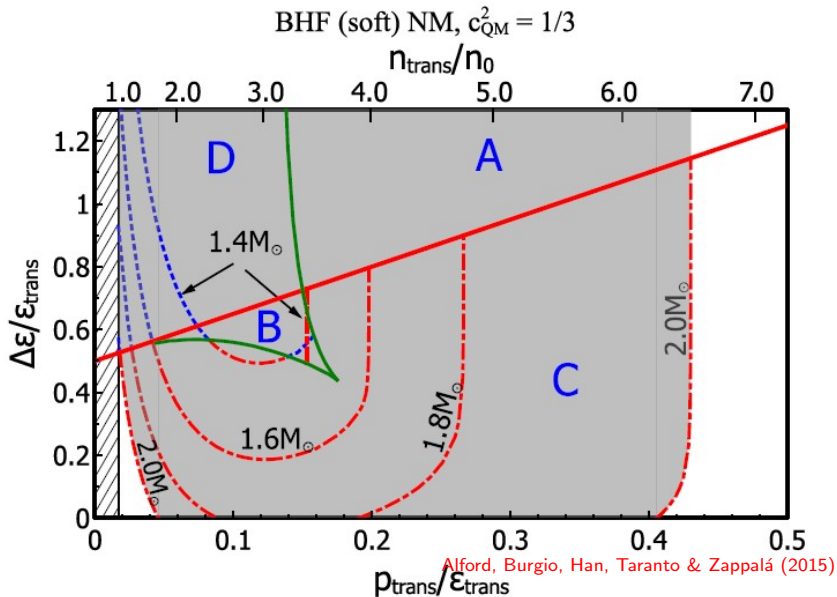
# Sound Speed in Quark Matter



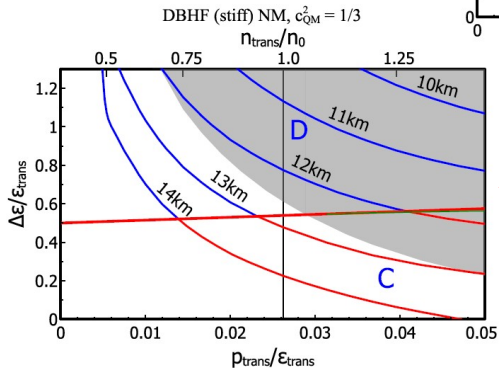
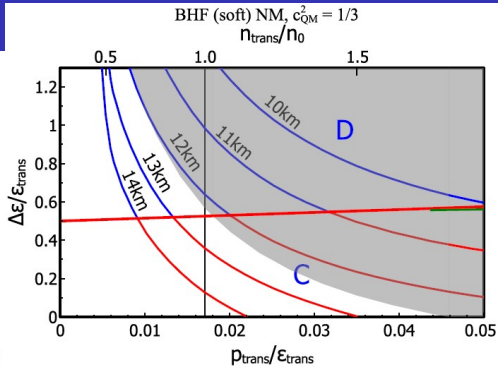
# Mass Constraint



# Mass Constraint

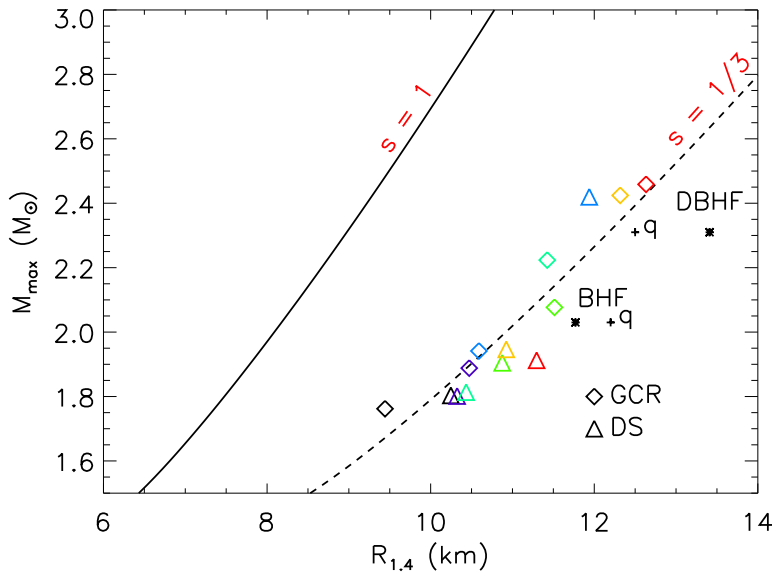


# Radius Constraint



Alford, Burgio, Han, Taranto & Zappalá (2015)

$$M_{max} - R_{1.4}$$



# Mass-Radius Diagram and Theoretical Constraints

GR:

$$R > 2GM/c^2$$

$P < \infty$ :

$$R > (9/4)GM/c^2$$

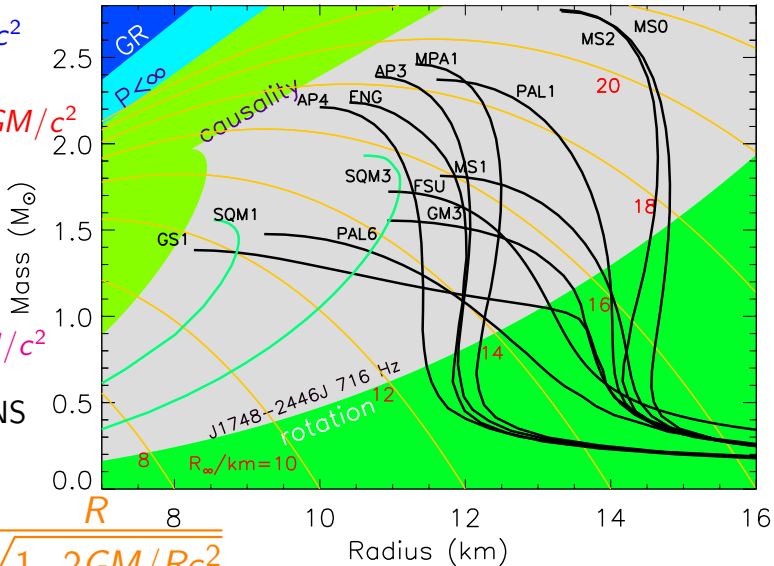
causality:

$$R \gtrsim 2.9GM/c^2$$

— normal NS

— SQS

$$R_\infty = \frac{R}{\sqrt{1 - 2GM/Rc^2}}$$

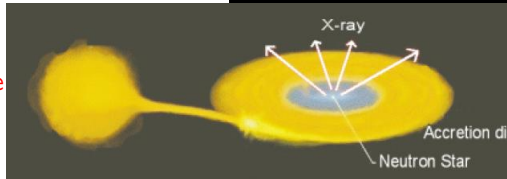


# Simultaneous Mass/Radius Measurements

- ▶ Measurements of flux  $F_\infty = (R_\infty/D)^2 \sigma T_{\text{eff}}^4$  and color temperature  $T_c \propto \lambda_{\text{max}}^{-1}$  yield an apparent angular size (pseudo-BB):

$$\frac{R_\infty}{D} = \frac{R}{D} \frac{1}{\sqrt{1 - 2GM/Rc^2}}$$

- ▶ Observational uncertainties include distance  $D$ , interstellar absorption  $N_H$ , atmospheric composition

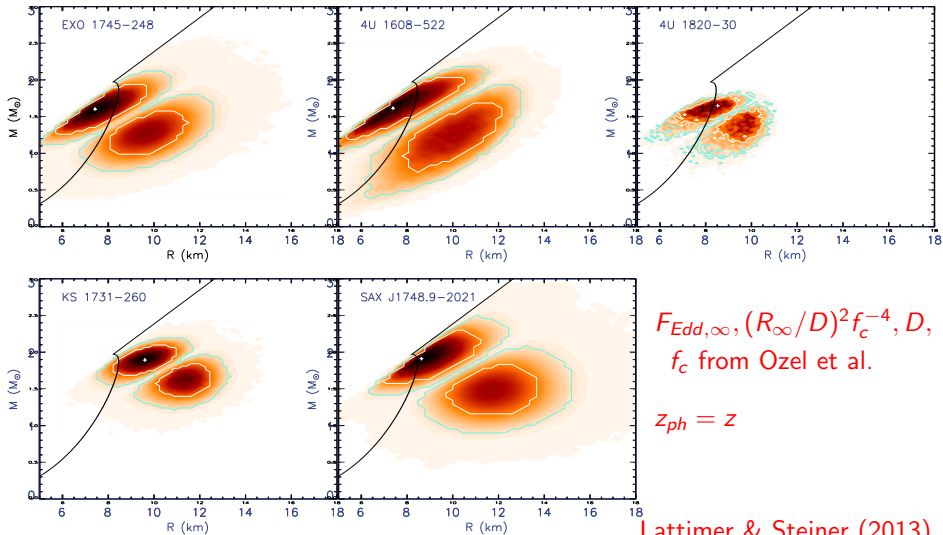


Best chances for accurate radius measurement:

- ▶ Nearby isolated neutron stars with parallax (uncertain atmosphere)
- ▶ Quiescent low-mass X-ray binaries (QLMXBs) in globular clusters (reliable distances, low  $B$  H-atmospheres)
- ▶ Bursting sources (XRBs) with peak fluxes close to Eddington limit (where gravity balances radiation pressure)

$$F_{\text{Edd}} = \frac{cGM}{\kappa D^2} \sqrt{1 - 2GM/Rc^2}$$

# M – R PRE Burst Estimates



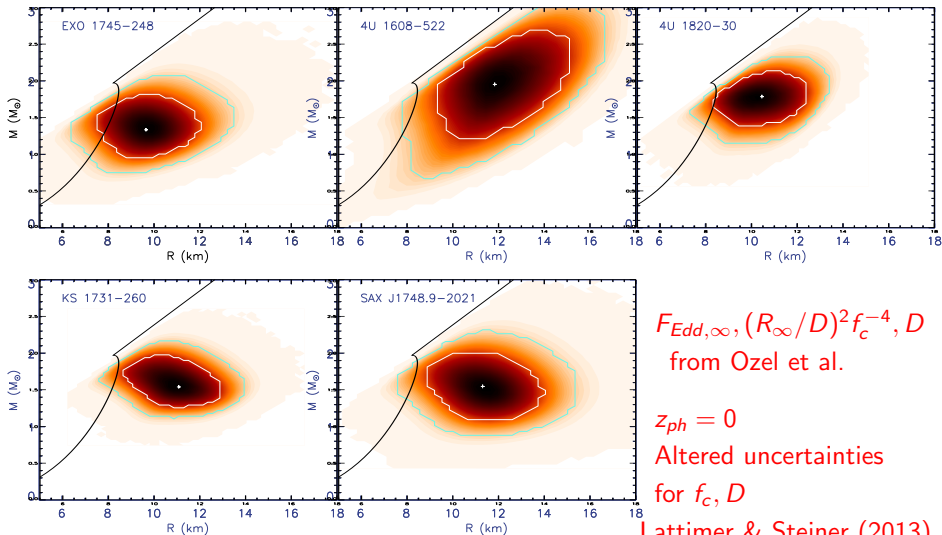
$F_{Edd,\infty}, (R_{\infty}/D)^2 f_c^{-4}, D,$   
 $f_c$  from Özel et al.

$Z_{ph} = Z$

Lattimer & Steiner (2013)



# $M - R$ PRE Burst Estimates

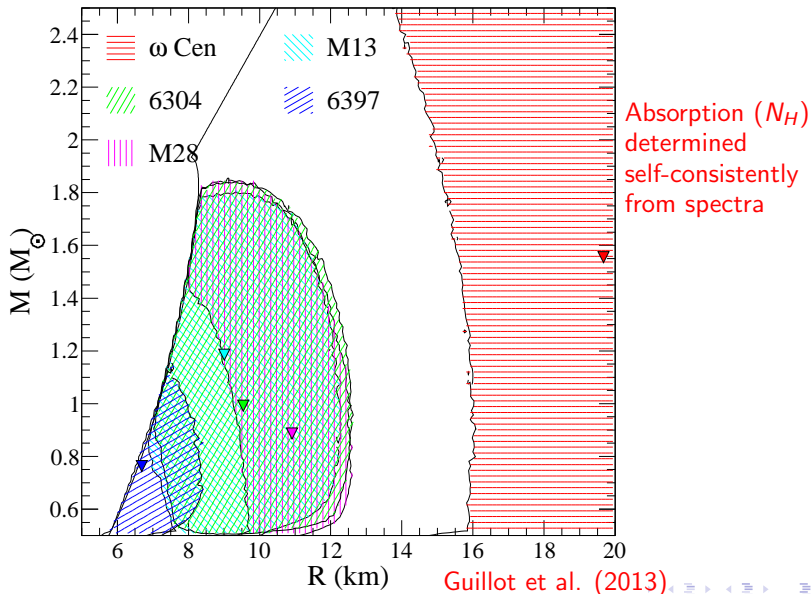


$F_{Edd,\infty}, (R_{\infty}/D)^2 f_c^{-4}, D$   
from Özel et al.

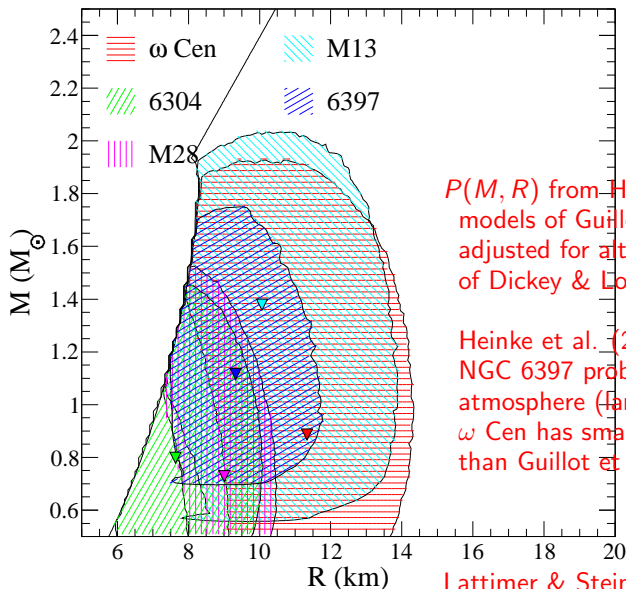
$z_{ph} = 0$   
Altered uncertainties  
for  $f_c, D$

Lattimer & Steiner (2013)

# M – R QLMXB Estimates



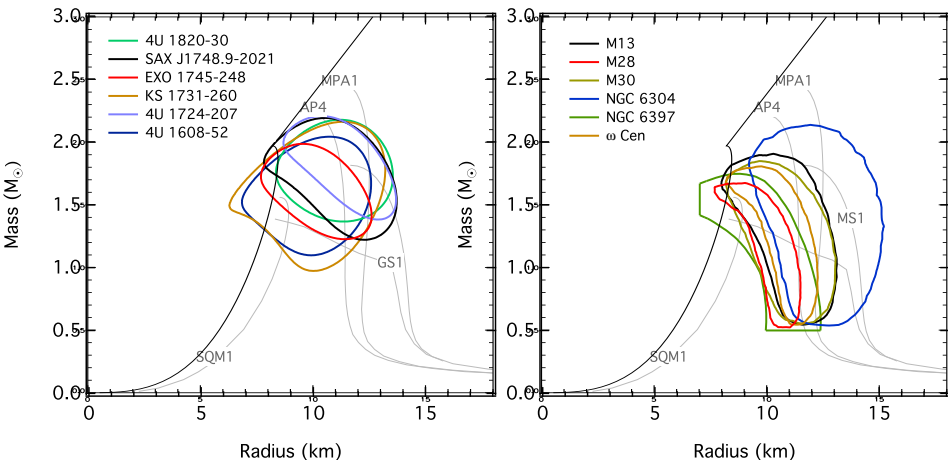
# M – R QLMXB Estimates



$P(M, R)$  from H atmosphere models of Guillot et al. (2013), adjusted for alternate  $N_H$  values of Dickey & Lockman (1990).

Heinke et al. (2014) found NGC 6397 probably has He atmosphere (larger  $R$ );  $\omega$  Cen has smaller  $N_H$  (and  $R$ ) than Guillot et al. (2013) found.

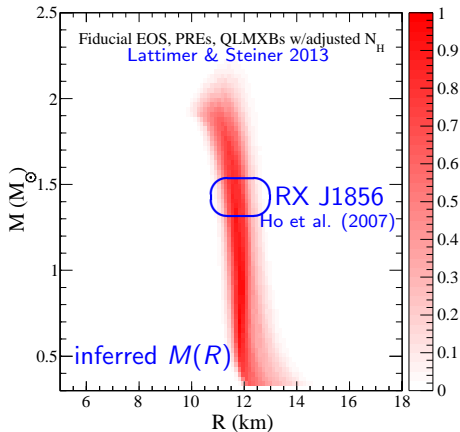
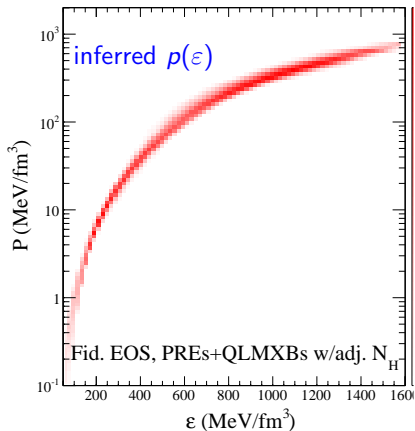
Lattimer & Steiner (2013)



Özel et al. (2015) adopts Bayesian approach following Steiner et al. (2010).  
 Mean XRB radius changes from  $9.74 \pm 0.5$  km to  $10.6 \pm 0.8$  km.  
 Neglect of causality and TOV constraints underestimate radius.  
 Steiner et al. analysis is vindicated.

# Bayesian TOV Inversion

- ▶  $\varepsilon < 0.5\varepsilon_0$ : Known crustal EOS
- ▶  $0.5\varepsilon_0 < \varepsilon < \varepsilon_1$ : EOS parametrized by  $K, K', S_V, \gamma$
- ▶ Polytropic EOS:  $\varepsilon_1 < \varepsilon < \varepsilon_2$ :  $n_1$ ;  $\varepsilon > \varepsilon_2$ :  $n_2$
- ▶ EOS parameters  $K, K', S_V, \gamma, \varepsilon_1, n_1, \varepsilon_2, n_2$  uniformly distributed
- ▶  $M_{\max} \geq 1.97 M_{\odot}$ , causality enforced
- ▶ All 10 stars equally weighted



# Astronomy vs. Astronomy vs. Physics

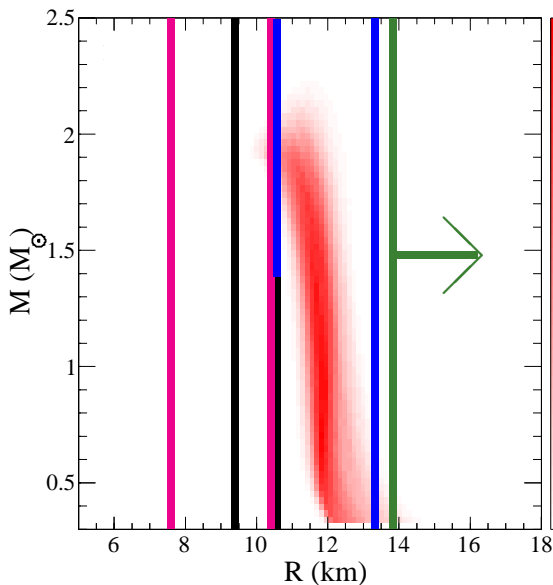
Ozel et al., XRB+QLMXB,  
 $M_{max} > 2M_{\odot}$ , crust,  $z_{ph} = z$ :  
 $R = 10.6 \pm 0.6$  km.

Suleimanov et al., long  
XRB:  $R_{1.4} \gtrsim 13.9$  km

Guillot et al. (2013),  
QLMXB, equal radii stars,  
self  $N_H$ :  $R = 9.1^{+1.3}_{-1.5}$  km.

Lattimer & Steiner (2013),  
XRB+QLMXB, TOV, crust,  
causality,  $M_{max} > 2M_{\odot}$ ,  
 $z_{ph} \neq z$ , alt  $N_H$ .

Lattimer & Lim (2013),  
nuclear experiments:  
 $29 \text{ MeV} < S_v < 33 \text{ MeV}$ ,  
 $40 \text{ MeV} < L < 65 \text{ MeV}$ :  
 $R_{1.4} = 12.0 \pm 1.4$  km.



# Additional Proposed Radius and Mass Constraints

## ▶ Pulse profiles

Hot or cold regions on rotating neutron stars alter pulse shapes: NICER and LOFT will enable timing and spectroscopy of thermal and non-thermal emissions. Light curve modeling  $\rightarrow M/R$ ; phase-resolved spectroscopy  $\rightarrow R$ .

## ▶ Moment of inertia

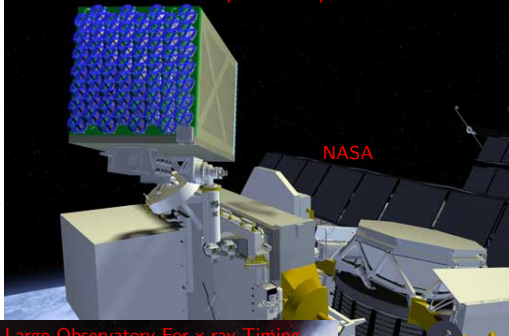
Spin-orbit coupling of ultra-relativistic binary pulsars (e.g., PSR 0737+3039) vary  $i$  and contribute to  $\dot{\omega}$ :  $I \propto MR^2$ .

## ▶ Supernova neutrinos

Millions of neutrinos detected from a Galactic supernova will measure  $BE = m_B N - M, \langle E_\nu \rangle, \tau_\nu$ .

## ▶ QPOs from accreting sources ISCO and crustal oscillations

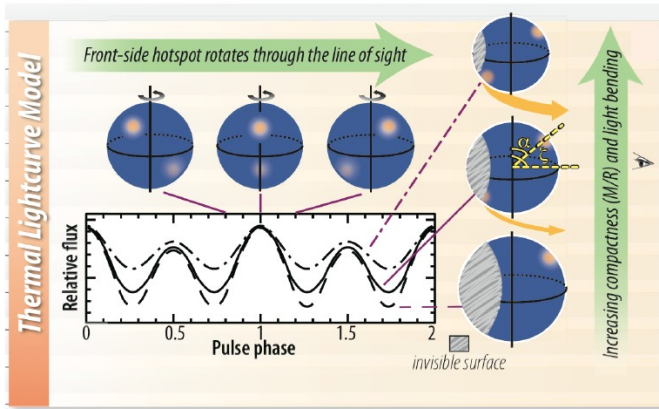
Neutron star Interior Composition Explorer



Large Observatory For x-ray Timing

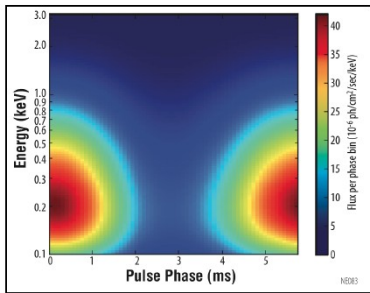


Reveal stellar structure through lightcurve modeling, long-term timing, and pulsation searches

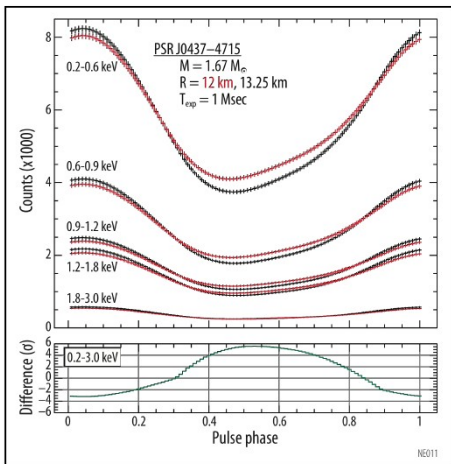


**Lightcurve modeling** constrains the compactness ( $M/R$ ) and viewing geometry of a non-accreting millisecond pulsar through the depth of modulation and harmonic content of emission from rotating hot-spots, thanks to **gravitational light-bending**...

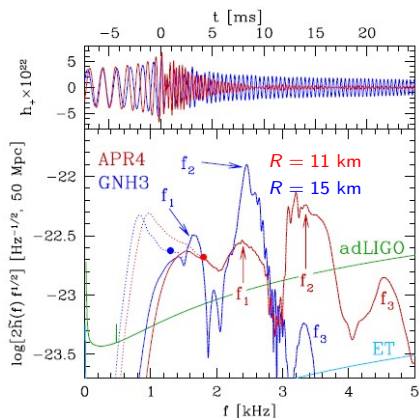




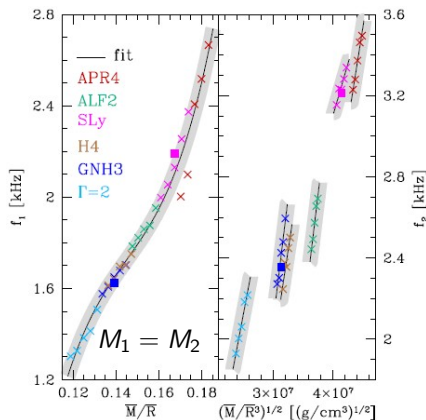
... while phase-resolved spectroscopy promises a direct constraint of radius  $R$ .



# Constraints from Observations of Gravitational Radiation



Takami, Rezzolla and Baiotti (2014)



- ▶ Chirp mass and tidal deformability measurable during inspiral.
- ▶ Frequency peaks are tightly correlated with compactness.
- ▶ Mass determinations from prompt and delayed black hole formation.
- ▶ In neutron star-black hole mergers, disc mass depends on  $a/M_{BH}$  and on  $M_{NS}M_{BH}/R^2$ .
- ▶ R-mode instabilities in rotating neutron stars.

# Conclusions

- ▶ Measured neutron star masses imply lower limits to radii of typical neutron stars.
- ▶ Symmetry energy determines typical neutron star radii.
- ▶ Nuclear experiments set reasonably tight constraints on symmetry energy parameters.
- ▶ Theoretical calculations of pure neutron matter predict very similar symmetry constraints.
- ▶ These constraints predict neutron star radii  $R_{1.4} = 12.0 \pm 1.4$  km.
- ▶ Combined astronomical observations of photospheric radius expansion X-ray bursts and quiescent sources in globular clusters suggest  $R_{1.4} \sim 12.1 \pm 0.6$  km.
- ▶ The properties of a high-density phase, such as quark matter, are tightly constrained by current mass measurements.
- ▶ A mass measurement above  $2.4M_{\odot}$  may be incompatible with other constraints, assuming GR is correct.