Saturation of the f-mode instability

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Workshop on binary neutron star mergers

Thessaloniki, 29.5.2015



Outline



- Fluid equations
- Classes of modes

2 The f-mode instability

- The CFS instability
- The instability window

3 Mode coupling

- Quadratic perturbation equations
- Equations of motion
- Parametric resonance instability
- Saturation conditions

4 Results and remarks

Saturation of the *f*-mode instability Coscillation modes Fluid equations

Oscillation modes

• Fluid equations (in corotating frame):

$$\dot{
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abla \cdot (
ho oldsymbol{v}) = 0$$

 $\dot{oldsymbol{v}} + (oldsymbol{v} \cdot
abla)oldsymbol{v} + 2oldsymbol{\Omega} imes oldsymbol{v} + oldsymbol{\Omega} imes (oldsymbol{\Omega} imes oldsymbol{r}) = -rac{
abla p}{
ho} -
abla \phi$
 $abla^2 \phi = 4\pi G
ho$
 $p = p(
ho, \mu)$

• *Linearly* perturbed fluid equations:

$$\begin{split} \delta\rho + \nabla \cdot (\rho \boldsymbol{\xi}) &= 0\\ \ddot{\boldsymbol{\xi}} + 2\boldsymbol{\Omega} \times \dot{\boldsymbol{\xi}} &= -\frac{\nabla \delta p}{\rho} + \frac{\nabla p}{\rho^2} \delta\rho - \nabla \delta\phi\\ \nabla^2 \delta\phi &= 4\pi G \delta\rho\\ \\ \Delta p \\ p &= \left(\frac{\partial \ln p}{\partial \ln \rho}\right)_{\mu} \frac{\Delta \rho}{\rho} + \left(\frac{\partial \ln p}{\partial \ln \mu}\right)_{\rho} \frac{\Delta \mu}{\mu} \end{split}$$

Eulerian (δ) and Langrangian (Δ) perturbations related via $\Delta f = \delta f + (\boldsymbol{\xi} \cdot \nabla) f$

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Eulerian (δ) and Langrangian (Δ) perturbations related via $\Delta f = \delta f + (\boldsymbol{\xi} \cdot \nabla) f$

Saturation of the *f*-mode instability Coscillation modes Classes of modes

Oscillation modes

$$\boldsymbol{\xi}(r,\theta,\phi) = \sum_{l=0}^{\infty} \sum_{m=-l}^{l} \left[U_l^m(r) Y_l^m(\theta,\phi) \hat{\boldsymbol{e}}_r + V_l^m(r) \nabla Y_l^m(\theta,\phi) + W_l^m(r) \hat{\boldsymbol{e}}_r \times \nabla Y_l^m(\theta,\phi) \right]$$

as $\Omega \to 0$

- Polar modes: $W_l^m = 0$
- Axial modes: $U_l^m = V_l^m = 0$

l :	degree
m:	order
n:	overtone



Mode name	Mode class	Mode type	Restoring force
<i>p</i> -mode	Polar	Sound wave $(\omega \to \infty \text{ as } n \to \infty)$	Prossure gradient
f-mode	Polar	Low- ω sound wave $\Big _{n} = 0$	
J-mode I ofai	High- ω gravity wave $\int n = 0$	Buoyaney	
$g ext{-mode}$	Polar	Gravity wave $(\omega \to 0 \text{ as } n \to \infty)$	Duoyancy
r-mode	Axial	Inertial wave	Coriolis
$Hybrid \mod$	Combination	Zero-buoyancy limit or r - and g -modes	
Only for non-zero rotation Only for non-zero buoyancy			

Saturation of the *f*-mode instability Coscillation modes Classes of modes

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The CFS instability

Are the perturbations stable?

Rapidly rotating stars are prone to *secular instabilities*, i.e. instabilities related to dissipation mechanisms (viscosity, gravitational radiation).



$$\left(\frac{\mathrm{d}E}{\mathrm{d}t}\right)_{\mathrm{GW}} = -\sum_{l\geq 2}^{\infty} N_l \,\omega \left(\omega - m\Omega\right)^{2l+1} \left(\left|\delta D_l^m\right|^2 + \left|\delta J_l^m\right|^2\right)$$

• Polar (axial) modes emit through the mass (current) multipoles

• If
$$\omega(\omega - m\Omega) < 0$$
, then $\left(\frac{\mathrm{d}E}{\mathrm{d}t}\right)_{\mathrm{GW}} > 0$

CFS instability

For any rotation rate Ω there is always a mode driven unstable by gravitational radiation emission [Chandrasekhar, 1970, Friedman and Schutz, 1978].

• f-modes and r-modes are the most susceptible to GW-driven instabilities

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inertial-frame frequency

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Saturation of the *f*-mode instability L The *f*-mode instability L The instability window

The instability window



Saturation of the *f*-mode instability Mode coupling Quadratic perturbation equations

Mode coupling

Do the unstable modes grow boundlessly?

Non-linear mode coupling inhibits the instability's growth

• Quadratically perturbed fluid equations:

$$\begin{split} \delta\dot{\rho} + \nabla \cdot (\rho \boldsymbol{v}) + \nabla \cdot (\delta\rho \, \boldsymbol{v}) &= 0\\ \ddot{\boldsymbol{\xi}} + \boldsymbol{\mathcal{B}}(\dot{\boldsymbol{\xi}}) + \boldsymbol{\mathcal{C}}\left(\boldsymbol{\xi}\right) + \boldsymbol{\mathcal{N}}\left(\boldsymbol{\xi}, \boldsymbol{\xi}\right) &= 0\\ \nabla^2 \delta\phi &= 4\pi G \delta\rho \end{split}$$
$$\frac{\Delta p}{p} &= \Gamma_1 \frac{\Delta \rho}{\rho} + \frac{1}{2} \left[\Gamma_1(\Gamma_1 - 1) + \left(\frac{\partial \Gamma_1}{\partial \ln \rho}\right)_{\mu} \right] \left(\frac{\Delta \rho}{\rho}\right)^2, \quad \left[\Gamma_1 = \left(\frac{\partial \ln p}{\partial \ln \rho}\right)_{\mu} \right] \end{split}$$

Eulerian (δ) and *Langrangian* (Δ) perturbations related via $\Delta f = \delta f + (\boldsymbol{\xi} \cdot \nabla) f + (\boldsymbol{\xi} \cdot \nabla) \delta f + \frac{1}{2} \boldsymbol{\xi} \cdot [\boldsymbol{\xi} \cdot \nabla (\nabla f)]$

• Perturbation decomposition:

$$\boldsymbol{\xi}(\boldsymbol{r},t) = \sum_{lpha} Q_{lpha}(t) \, \boldsymbol{\xi}_{lpha}(\boldsymbol{r}) e^{i\omega_{lpha}t}$$

Saturation of the *f*-mode instability Mode coupling Quadratic perturbation equations

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• Perturbation decomposition:

Mode coupling

 \bullet Modes couple in triplets

$$\begin{split} \dot{Q}_{\alpha} &= \gamma_{\alpha} Q_{\alpha} + i \omega_{\alpha} \mathcal{H} Q_{\beta} Q_{\gamma} e^{-i \Delta \omega t} \\ \dot{Q}_{\beta} &= \gamma_{\beta} Q_{\beta} + i \omega_{\beta} \mathcal{H} Q_{\gamma}^{*} Q_{\alpha} e^{i \Delta \omega t} \\ \dot{Q}_{\gamma} &= \gamma_{\gamma} Q_{\gamma} + i \omega_{\gamma} \mathcal{H} Q_{\alpha} Q_{\beta}^{*} e^{i \Delta \omega t} \end{split}$$

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• Detuning
$$\Delta \omega \equiv \omega_{\alpha} - \omega_{\beta} - \omega_{\gamma} \approx 0$$

resonance condition

The system exhibits internal resonances

Mode coupling

• Modes couple in *triplets*

$$\begin{split} \dot{Q}_{\alpha} &= \gamma_{\alpha} Q_{\alpha} + i \omega_{\alpha} \mathcal{H} \, Q_{\beta} Q_{\gamma} \, e^{-i \Delta \omega t} \\ \dot{Q}_{\beta} &= \gamma_{\beta} Q_{\beta} + i \omega_{\beta} \mathcal{H} \, Q_{\gamma}^{*} Q_{\alpha} \, e^{i \Delta \omega t} \\ \dot{Q}_{\gamma} &= \gamma_{\gamma} Q_{\gamma} + i \omega_{\gamma} \mathcal{H} \, Q_{\alpha} Q_{\beta}^{*} \, e^{i \Delta \omega t} \end{split}$$

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$$\Delta \omega \equiv \omega_{\alpha} - \omega_{\beta} - \omega_{\gamma} \approx 0$$

resonance condition

• Coupling coefficient $\mathcal{H} \neq 0$ if

$$\begin{array}{c} m_{\alpha} = m_{\beta} + m_{\gamma} \\ l_{\alpha} + l_{\beta} + l_{\gamma} = \text{even number} \\ |l_{\beta} - l_{\gamma}| \leq l_{\alpha} \leq l_{\beta} + l_{\gamma} \end{array} \right\} \hline \text{coupling selection rules}$$

Mode coupling

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$$\begin{split} \dot{Q}_{\alpha} &= \gamma_{\alpha} Q_{\alpha} + i \omega_{\alpha} \mathcal{H} Q_{\beta} Q_{\gamma} e^{-i \Delta \omega t} \\ \dot{Q}_{\beta} &= \gamma_{\beta} Q_{\beta} + i \omega_{\beta} \mathcal{H} Q_{\gamma}^{*} Q_{\alpha} e^{i \Delta \omega t} \\ \dot{Q}_{\gamma} &= \gamma_{\gamma} Q_{\gamma} + i \omega_{\gamma} \mathcal{H} Q_{\alpha} Q_{\beta}^{*} e^{i \Delta \omega t} \end{split}$$

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 coupling selection rules

•

• Growth/damping rates $\gamma_i = \frac{1}{2E_i} \frac{\mathrm{d}E_i}{\mathrm{d}t} \gtrless 0$

$$\frac{\mathrm{d}E}{\mathrm{d}t} = \left(\frac{\mathrm{d}E}{\mathrm{d}t}\right)_{\mathrm{GW}} + \left(\frac{\mathrm{d}E}{\mathrm{d}t}\right)_{\mathrm{BV}} + \left(\frac{\mathrm{d}E}{\mathrm{d}t}\right)_{\mathrm{SV}} \gtrless 0$$

Saturation of the *f*-mode instability Mode coupling Parametric resonance instability

Parametric resonance instability

$$\begin{array}{c|c} \dot{Q}_{\alpha} = \gamma_{\alpha} Q_{\alpha} + i\omega_{\alpha} \mathcal{H} Q_{\beta} Q_{\gamma} e^{-i\Delta\omega t} & \text{Detuning } \Delta\omega \\ \dot{Q}_{\beta} = \gamma_{\beta} Q_{\beta} + i\omega_{\beta} \mathcal{H} Q_{\gamma}^{*} Q_{\alpha} e^{i\Delta\omega t} & \text{Coupling coefficient } \mathcal{H} \\ \dot{Q}_{\gamma} = \gamma_{\gamma} Q_{\gamma} + i\omega_{\gamma} \mathcal{H} Q_{\alpha} Q_{\beta}^{*} e^{i\Delta\omega t} & \text{Growth/damping rates } \gamma_{i} \end{array}$$

- Parent mode: unstable f-mode $(\gamma_{\alpha} > 0)$
- Daughter modes: other (stable) polar modes $(\gamma_{\beta,\gamma} < 0)$

Parametric resonance instability

- Parent feeds daughters and makes them grow
- Parametric instability threshold: daughters grow when

$$|Q_{\alpha}|^{2} > |Q_{\rm PIT}|^{2} \equiv \frac{\gamma_{\beta}\gamma_{\gamma}}{\omega_{\beta}\omega_{\gamma}\mathcal{H}^{2}} \left[1 + \left(\frac{\Delta\omega}{\gamma_{\beta} + \gamma_{\gamma}}\right)^{2}\right]$$

• Saturation amplitude: parent saturates at

$$|Q_{\rm sat}|^2 \equiv \frac{\gamma_\beta \gamma_\gamma}{\omega_\beta \omega_\gamma \mathcal{H}^2} \left[1 + \left(\frac{\Delta \omega}{\gamma_\alpha + \gamma_\beta + \gamma_\gamma} \right)^2 \right] \approx |Q_{\rm PIT}|^2$$

Saturation of the *f*-mode instability Mode coupling Parametric resonance instability

Parametric resonance instability

$$\begin{array}{l} \dot{Q}_{\alpha} = \gamma_{\alpha} Q_{\alpha} + i \omega_{\alpha} \mathcal{H} Q_{\beta} Q_{\gamma} e^{-i\Delta\omega t} \\ \dot{Q}_{\beta} = \gamma_{\beta} Q_{\beta} + i \omega_{\beta} \mathcal{H} Q_{\gamma}^{*} Q_{\alpha} e^{i\Delta\omega t} \\ \dot{Q}_{\gamma} = \gamma_{\gamma} Q_{\gamma} + i \omega_{\gamma} \mathcal{H} Q_{\alpha} Q_{\beta}^{*} e^{i\Delta\omega t} \\ \end{array} \right| \left| Q_{\mathrm{PIT}} \right|^{2} \equiv \frac{\gamma_{\beta} \gamma_{\gamma}}{\omega_{\beta} \omega_{\gamma} \mathcal{H}^{2}} \left[1 + \left(\frac{\Delta\omega}{\gamma_{\beta} + \gamma_{\gamma}} \right)^{2} \right] \\ \end{array} \right| \begin{array}{l} \mathrm{Detuning} \ \Delta\omega \\ \mathrm{Coupling \ coefficient} \ \mathcal{H} \\ \mathrm{Growth/damping \ rates} \ \gamma_{i} \end{array}$$



Saturation of the *f*-mode instability Mode coupling Saturation conditions

Saturation conditions

$$\begin{array}{l} \dot{Q}_{\alpha} = \gamma_{\alpha} Q_{\alpha} + i\omega_{\alpha} \mathcal{H} Q_{\beta} Q_{\gamma} e^{-i\Delta\omega t} \\ \dot{Q}_{\beta} = \gamma_{\beta} Q_{\beta} + i\omega_{\beta} \mathcal{H} Q_{\gamma}^{*} Q_{\alpha} e^{i\Delta\omega t} \\ \dot{Q}_{\gamma} = \gamma_{\gamma} Q_{\gamma} + i\omega_{\gamma} \mathcal{H} Q_{\alpha} Q_{\beta}^{*} e^{i\Delta\omega t} \end{array} \right| \begin{array}{l} \gamma_{\alpha} > 0, \ \gamma_{\beta,\gamma} < 0 \\ |Q_{\mathrm{PIT}}|^{2} \equiv \frac{\gamma_{\beta} \gamma_{\gamma}}{\omega_{\beta} \omega_{\gamma} \mathcal{H}^{2}} \left[1 + \left(\frac{\Delta\omega}{\gamma_{\beta} + \gamma_{\gamma}} \right)^{2} \right] \end{array} \begin{array}{l} \begin{array}{l} \text{Detuning } \Delta\omega \\ \text{Coupling coefficient } \mathcal{H} \\ \text{Growth/damping rates } \gamma_{i} \end{array}$$

• Saturation successful if:

 $|\gamma_{\beta} + \gamma_{\gamma}| \gtrsim \gamma_{\alpha} \quad \text{and} \quad \Delta \omega \gtrsim |\gamma_{\alpha} + \gamma_{\beta} + \gamma_{\gamma}|$



Results and remarks



Figure: Model: $M = 1.4 M_{\odot}, R = 10 \text{ km}, p \propto \rho^3, \Gamma_1 = 3.1$

- Neutron star equation of state probing: GW asteroseismology
- Post-merger remnants: high angular velocities, large growth rates
- Competing mechanisms: r-mode instability, magnetic field
- r-mode saturation amplitude $\sim 10^{-6} 10^{-5}$ [Brink et al., 2004]