

Gravitational waves and dynamical mass ejection from binary neutron-star mergers

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In collaboration with Hotokezaka, Kiuchi, Kyutoku, Okawa
Sekiguchi, M. Tanaka, & Wanajo

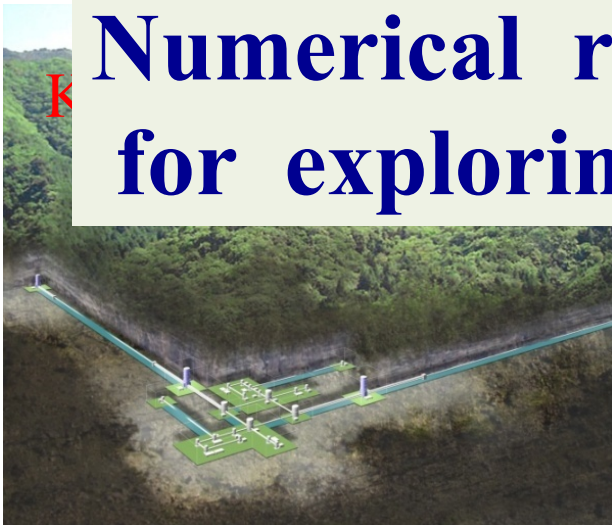
Contents

1. **One page introduction**
2. **Our latest numerical-relativity activity for NS-NS**
 - **Gravitational waves from late inspiral**
 - **Dynamical mass ejection**

Why NS-NS mergers are important ?

1. Most promising sources of gravitational waves for LIGO/VIRGO/KAGRA
2. Invaluable laboratory for studying high-density nuclear matter
3. Promising origins of short-hard GRBs
4. Sources of strong transient EM emission
5. Possible site for r-process heavy elements

Numerical relativity is the powerful tool for exploring these issues quantitatively



GOLD EXPLOSION New observations suggest that colliding neutron stars (shown in this artist's conception) produce short gamma-ray bursts. Such collisions also eject material that may be the source of the universe's gold and other heavy elements.

2A Gravitational waves (see John's talk)

Early Inspiral

$$(r_{\text{orb}} \gg R_{\text{NS}})$$

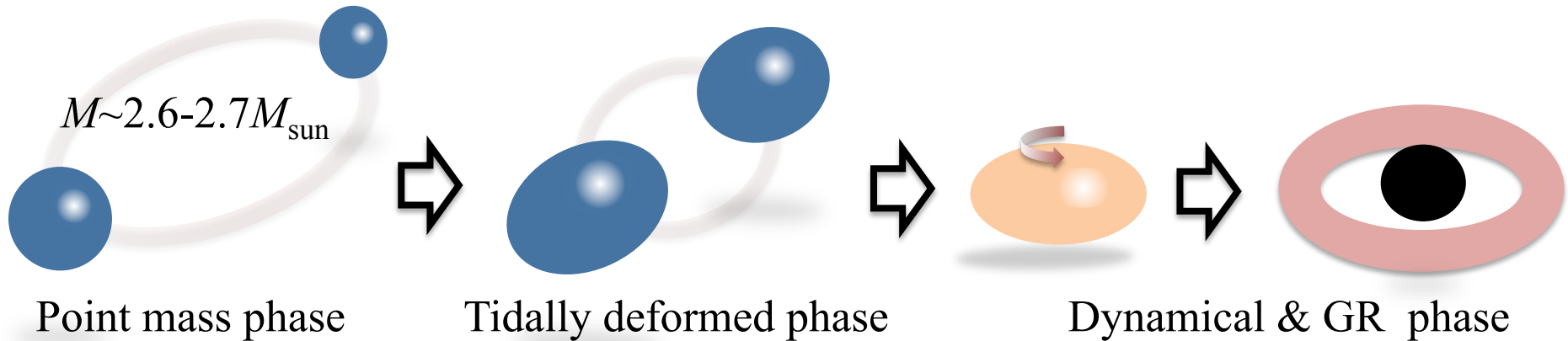
Late inspiral

$$(r_{\text{orb}} \leq 5R_{\text{NS}})$$

Merger =>

Hypermassive NS

Black hole & torus
& GRB?



Point mass phase
Adiabatic phase

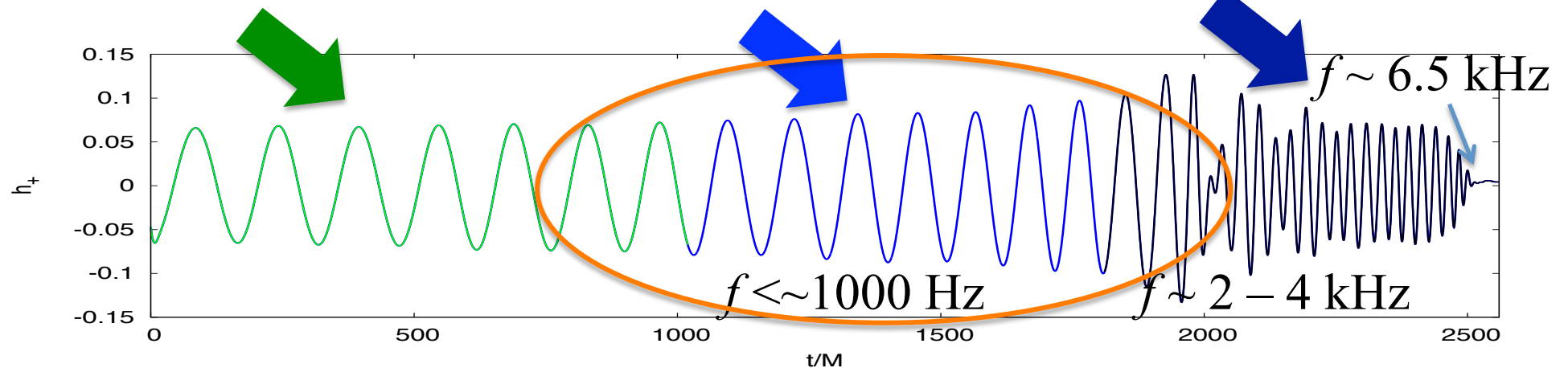
Tidally deformed phase

Dynamical & GR phase

Post-Newton

Late inspiral

Post merger

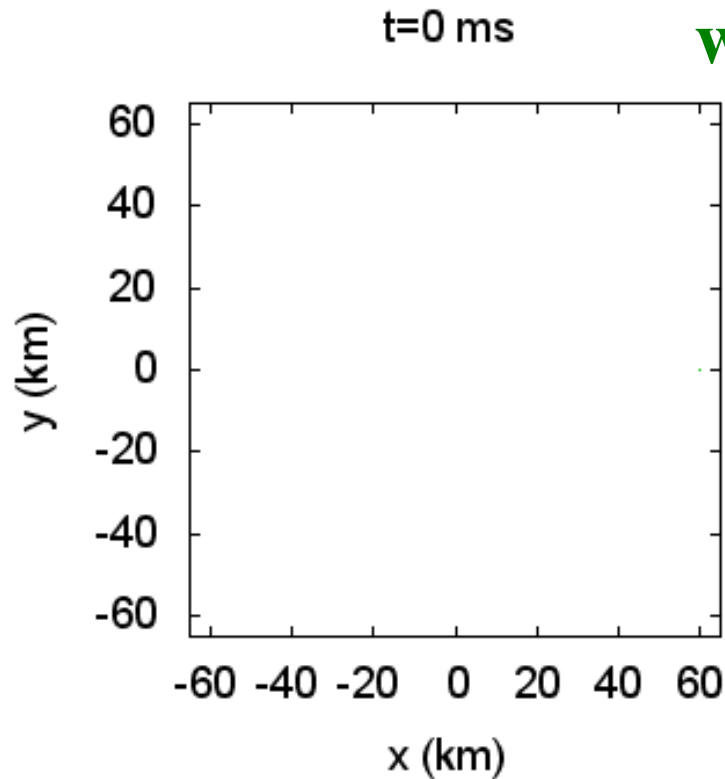


Analytic Computation (Effective One-Body)

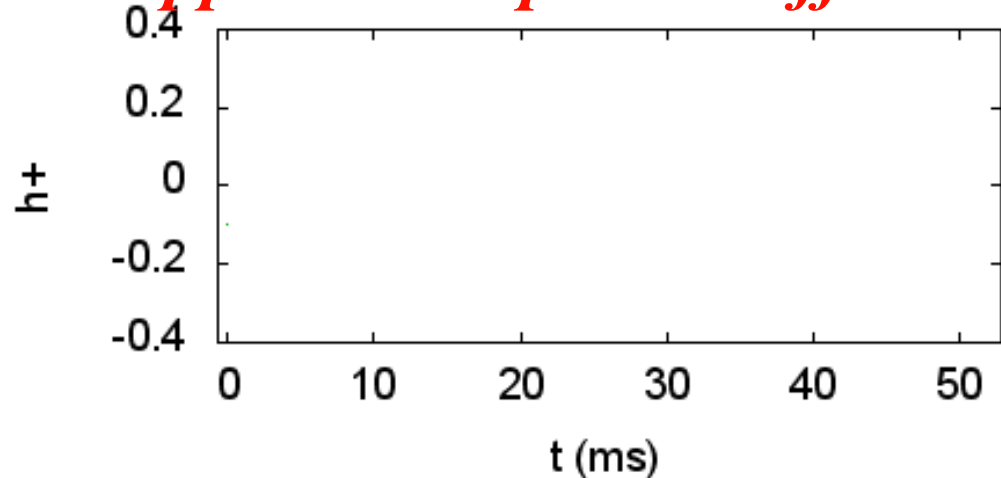
1.35-1.35 M_{sun} , EOS: MS1 ($R=14.5\text{km}$)

without tidal effects

with tidal effects



Appreciable phase difference



Calculation by Hotokezaka

EOB version (a bit old):
Damour et al., (2012)

Predicting more accurate GWs is urgent

Three key elements for deriving accurate gravitational waves in numerical rela

- ✧ Longterm simulation
- ✧ Eccentricity reduction for initial condition
- ✧ Extrapolation using high-quality data



Eccentricity reduction by Kyutoku+ 2014 ($e < \sim 0.001$)

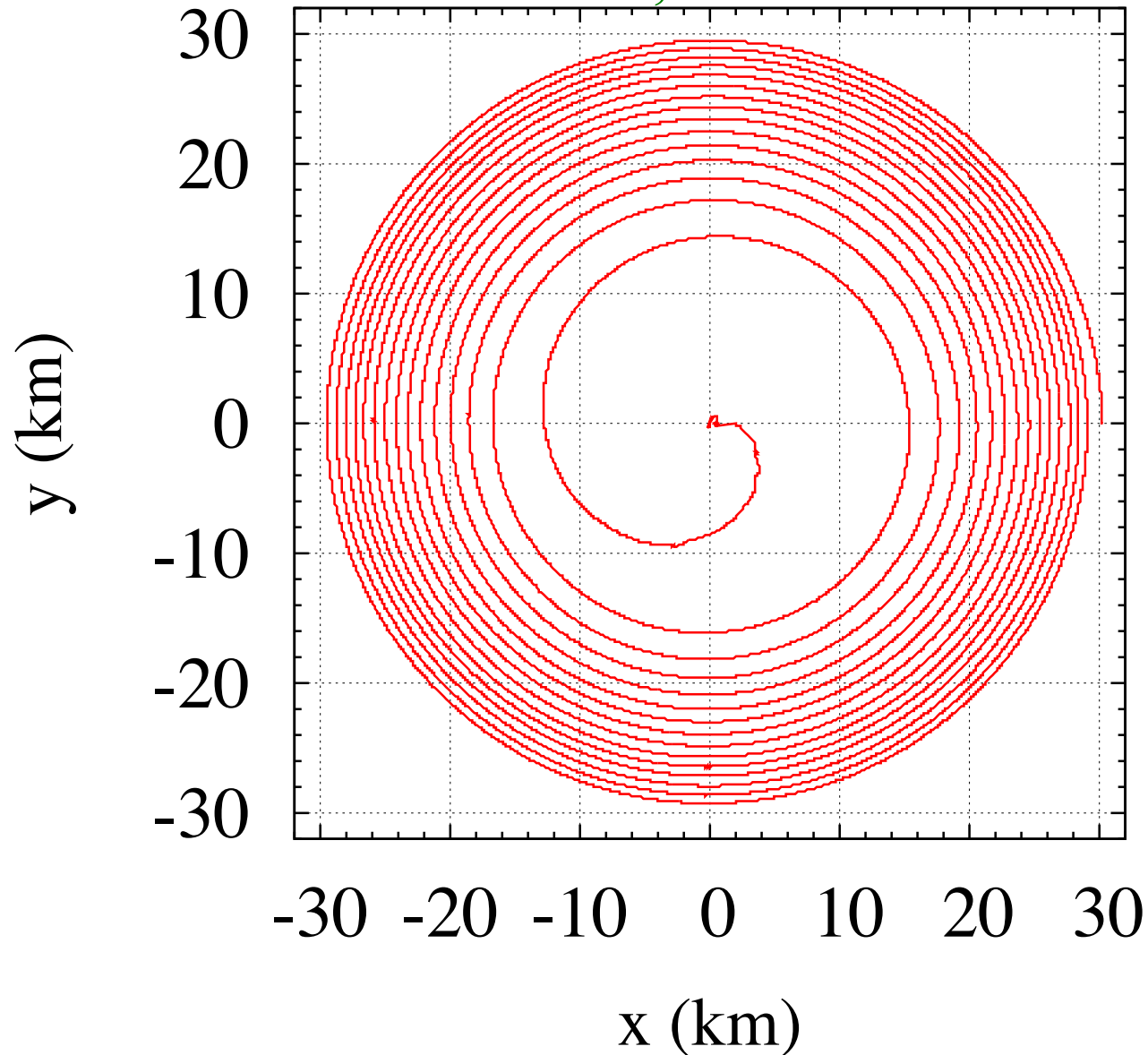
Constraint propagation by BSSN+Z4c prescription
(we locally used Z4c) (Hilditch & Bernuzzi)

→ less numerical error & good convergence

◆ See also the talk by Roland Haas

Our 15-orbits simulation with eccentricity reduction

It looks nice, however....



Hotokezaka, Kyutoku,
Okawa, Shibata,
PRD 91, 2015

H4-EOS:
R=13.6km
Mass:1.35-1.35 M_{sun}

However, at best, 3rd—4th-order convergence

We can never obtain exact numerical waveform
in hydrodynamics simulation !!

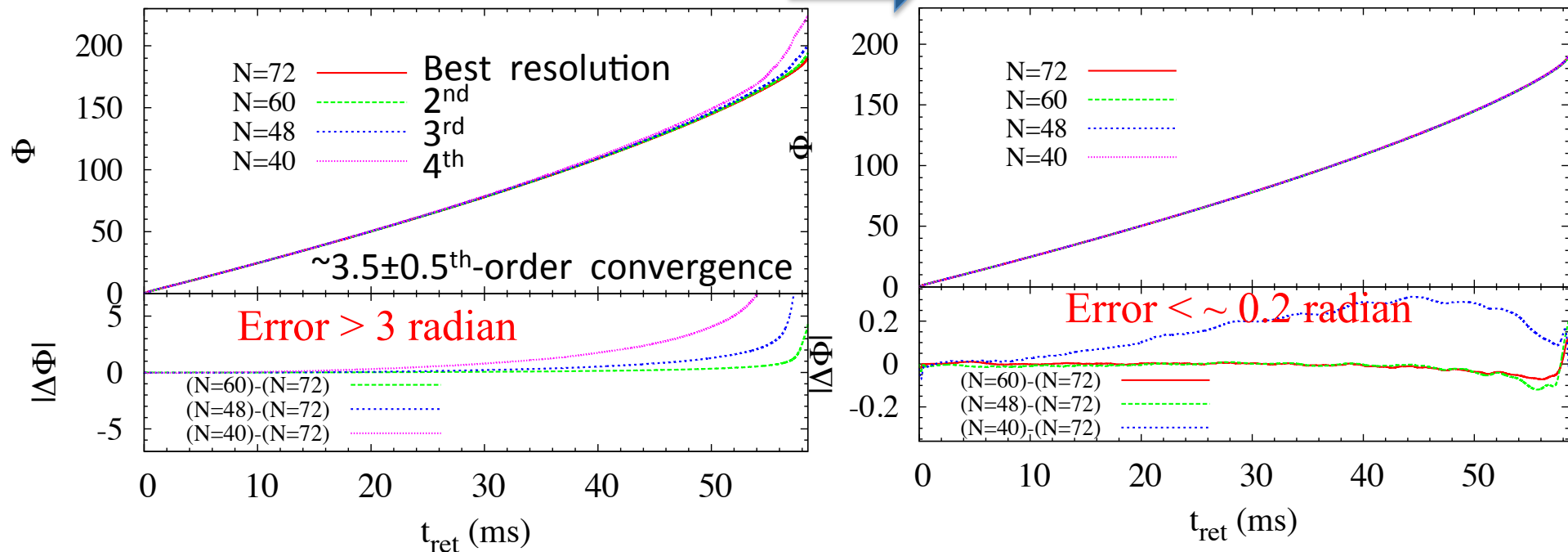
→ Extrapolation is needed for an “almost” solution

$$t \rightarrow \eta t, \quad \Phi = \int 2\pi f d(\eta t): \quad \eta = \text{const}$$

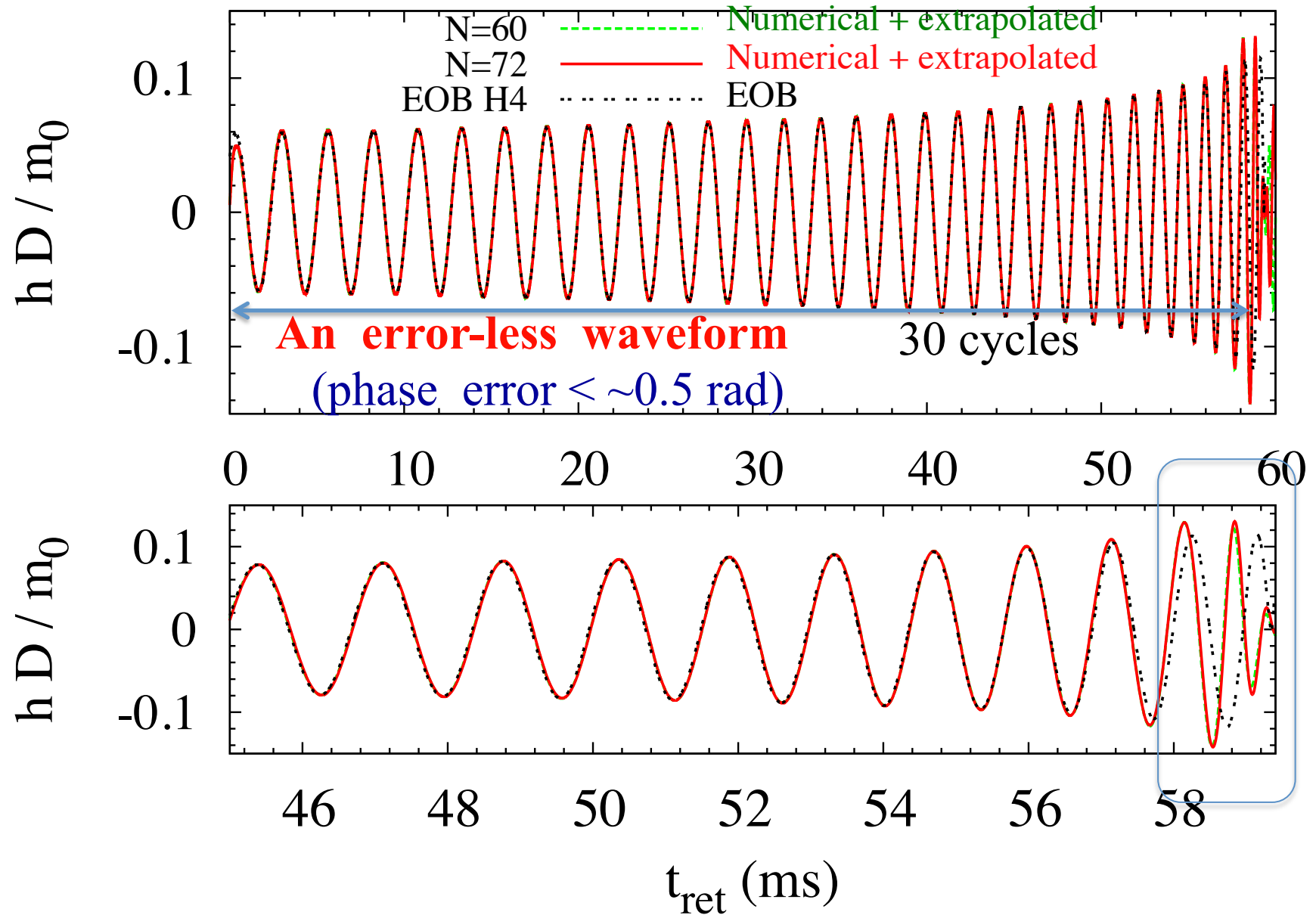
Gravitational wave phase



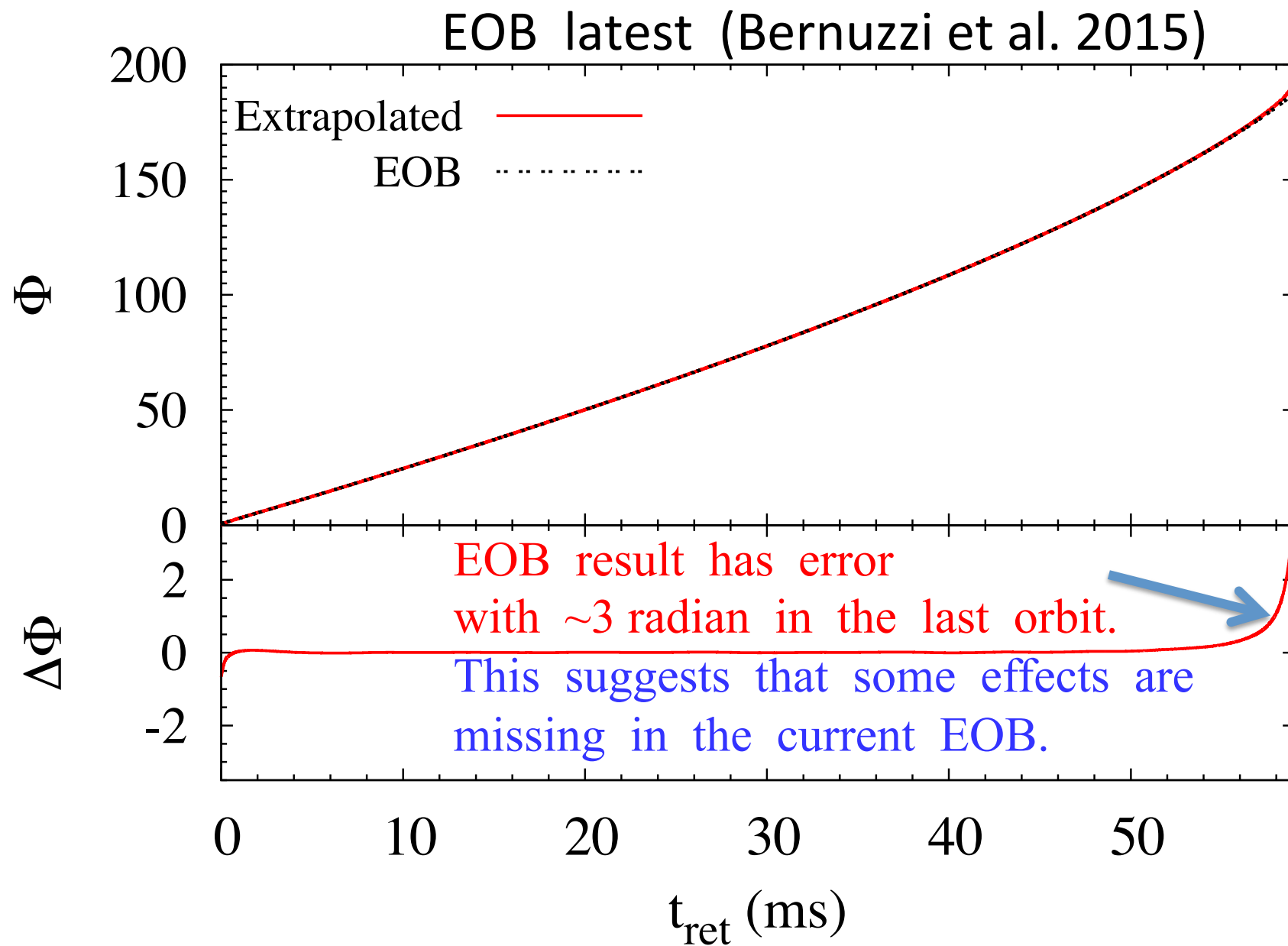
Extrapolated phase



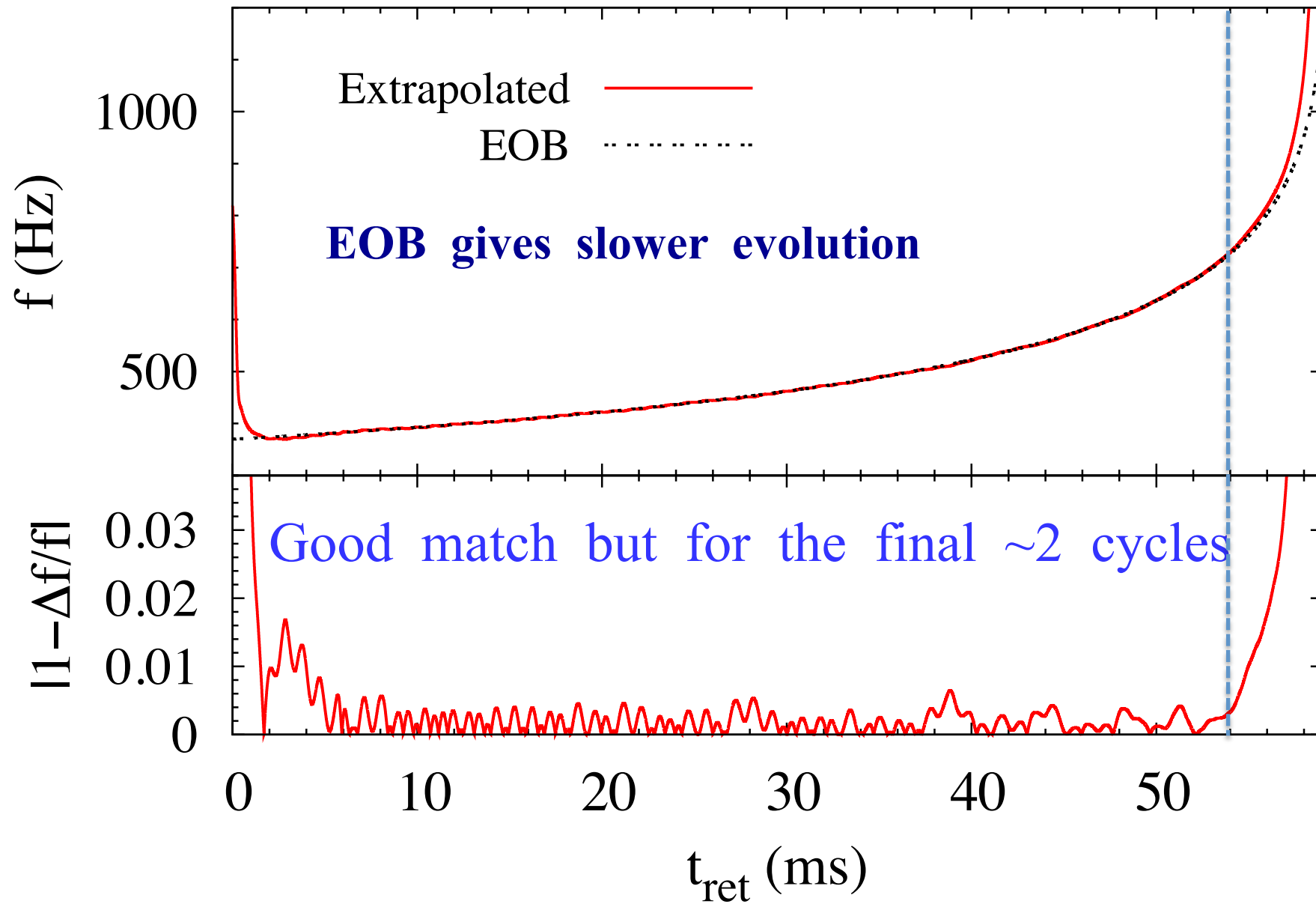
Extrapolated waveform for R=13.6 km



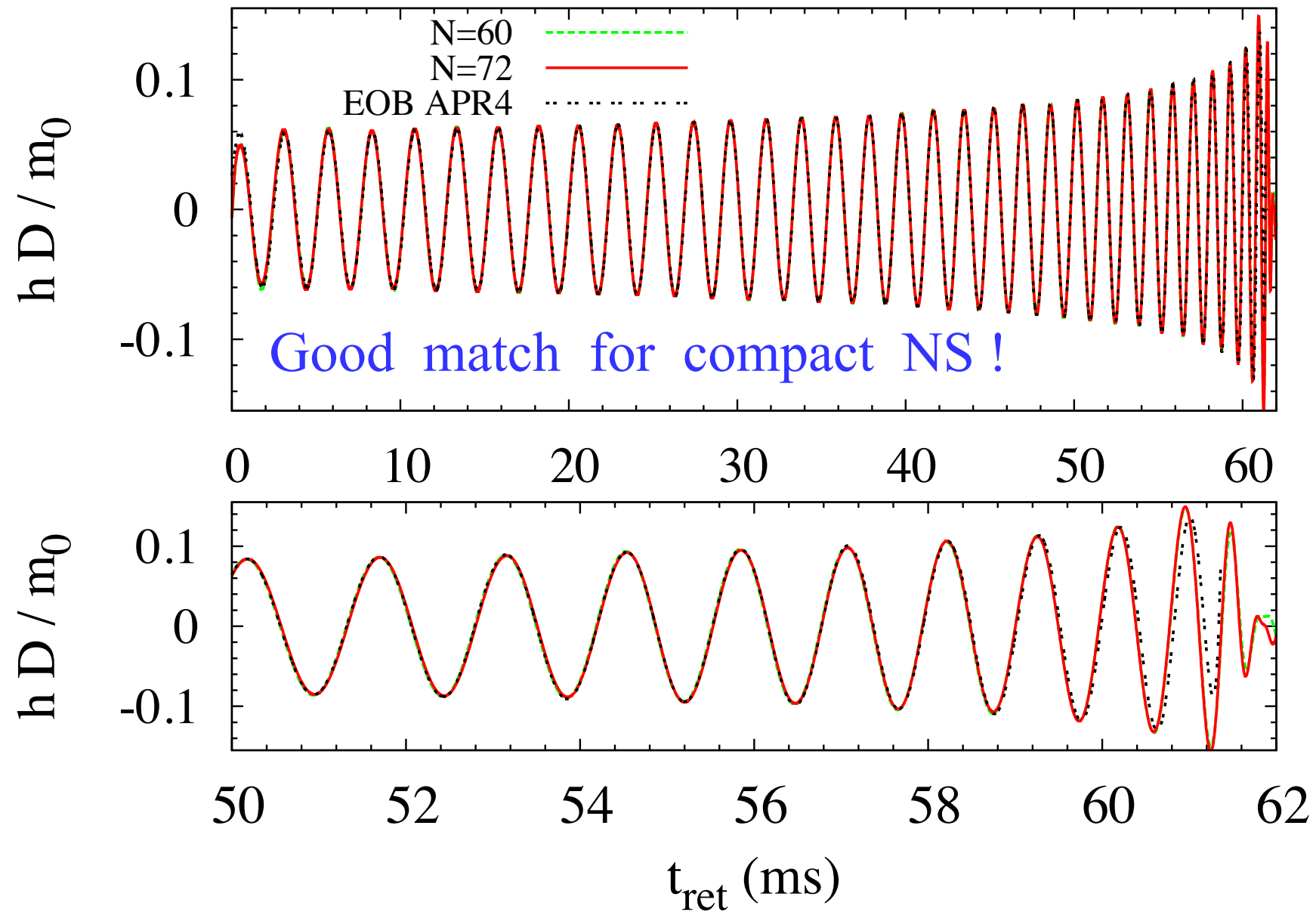
Comparison with effective-one-body approach



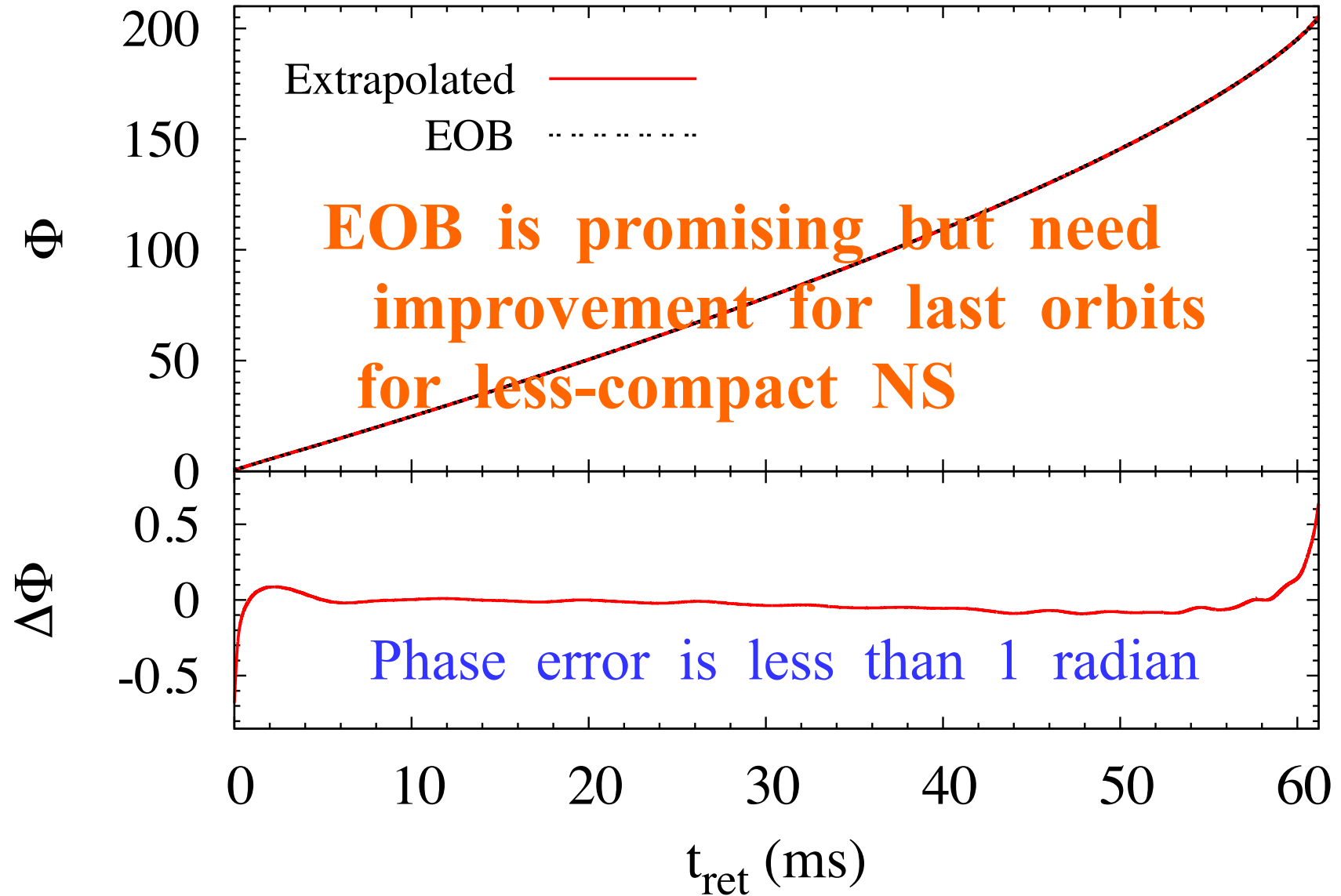
Comparison with EOB: frequency



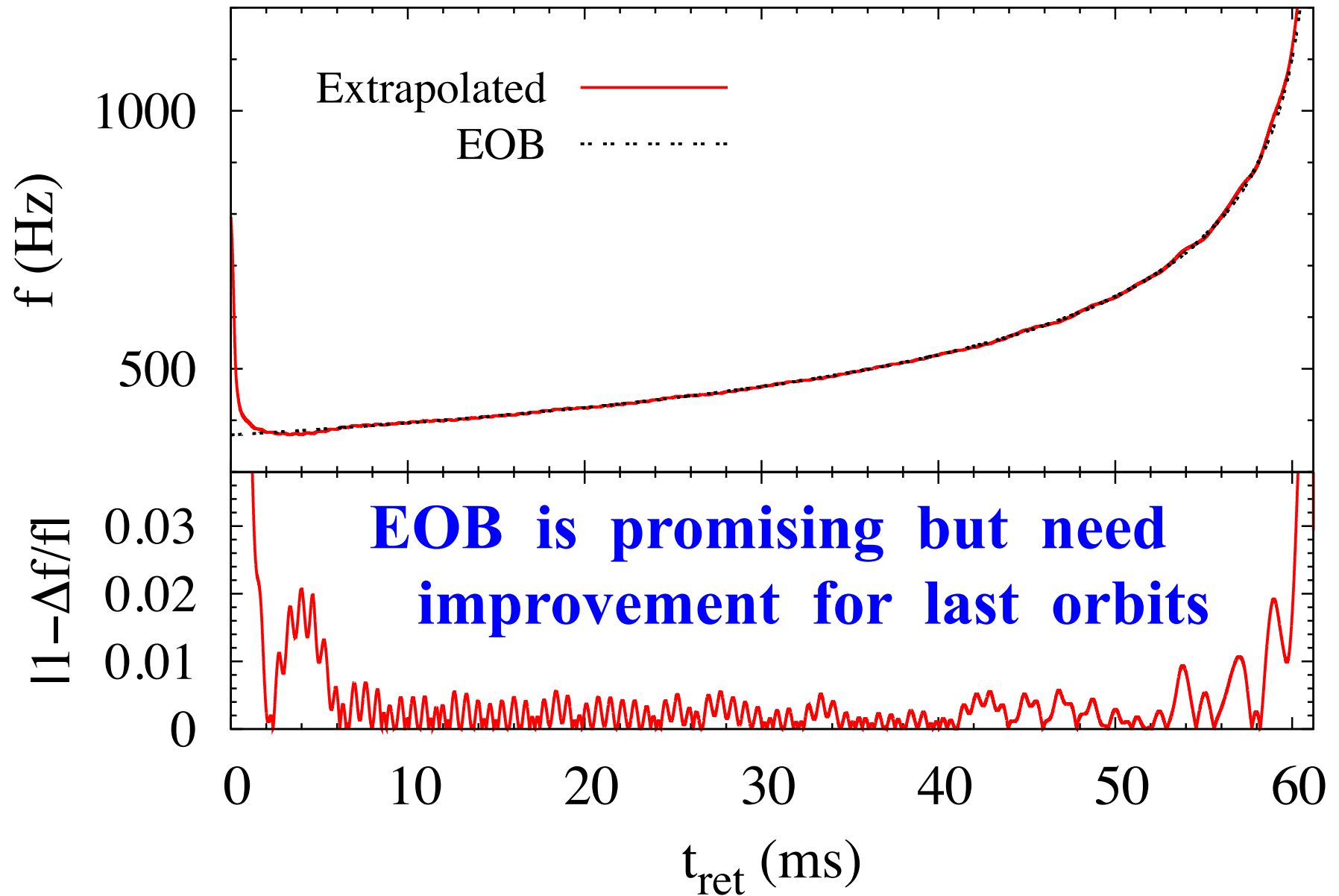
Extrapolated waveform for $R=11.1$ km



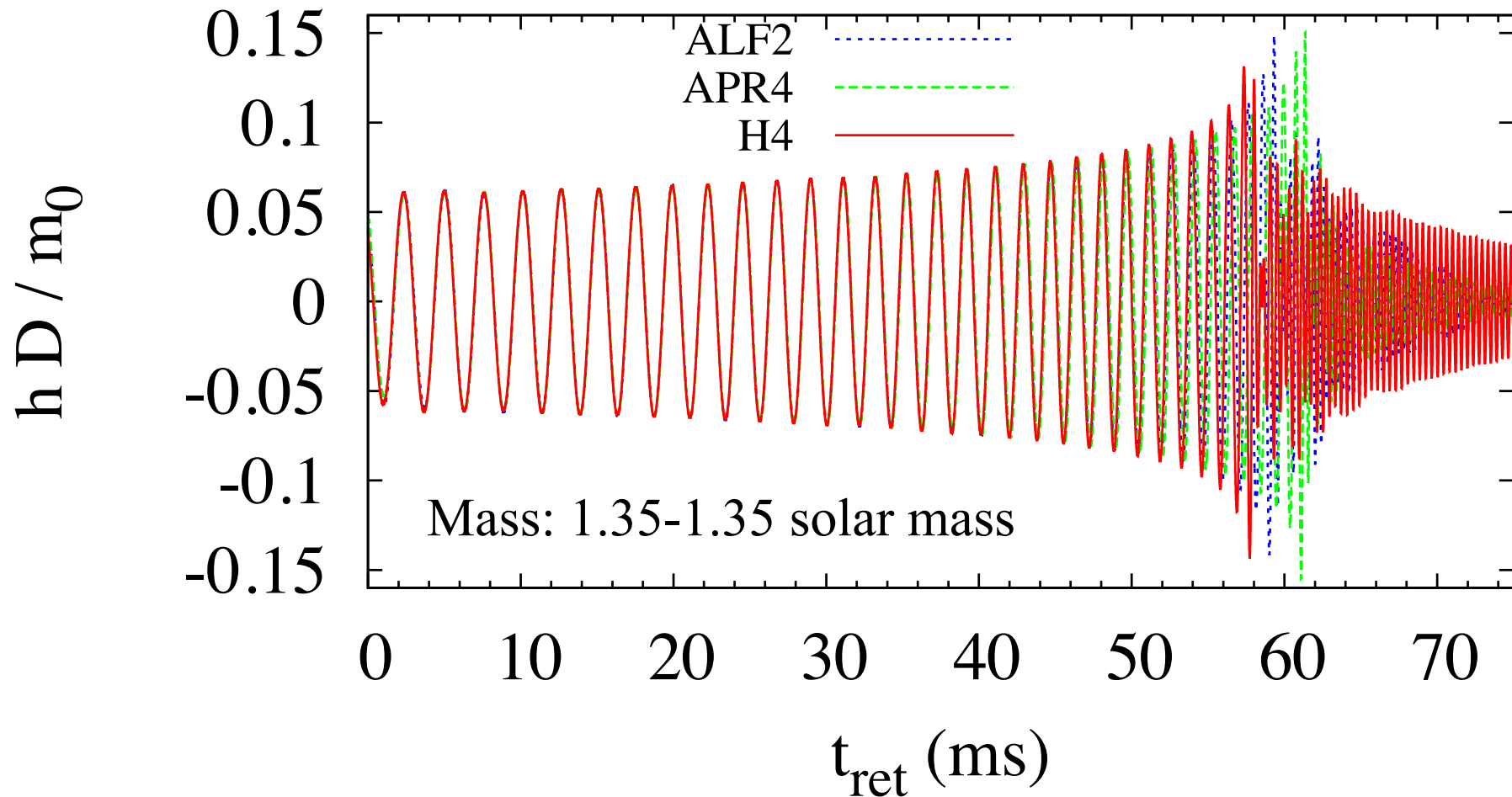
For soft EOS, current EOB is good



Good match

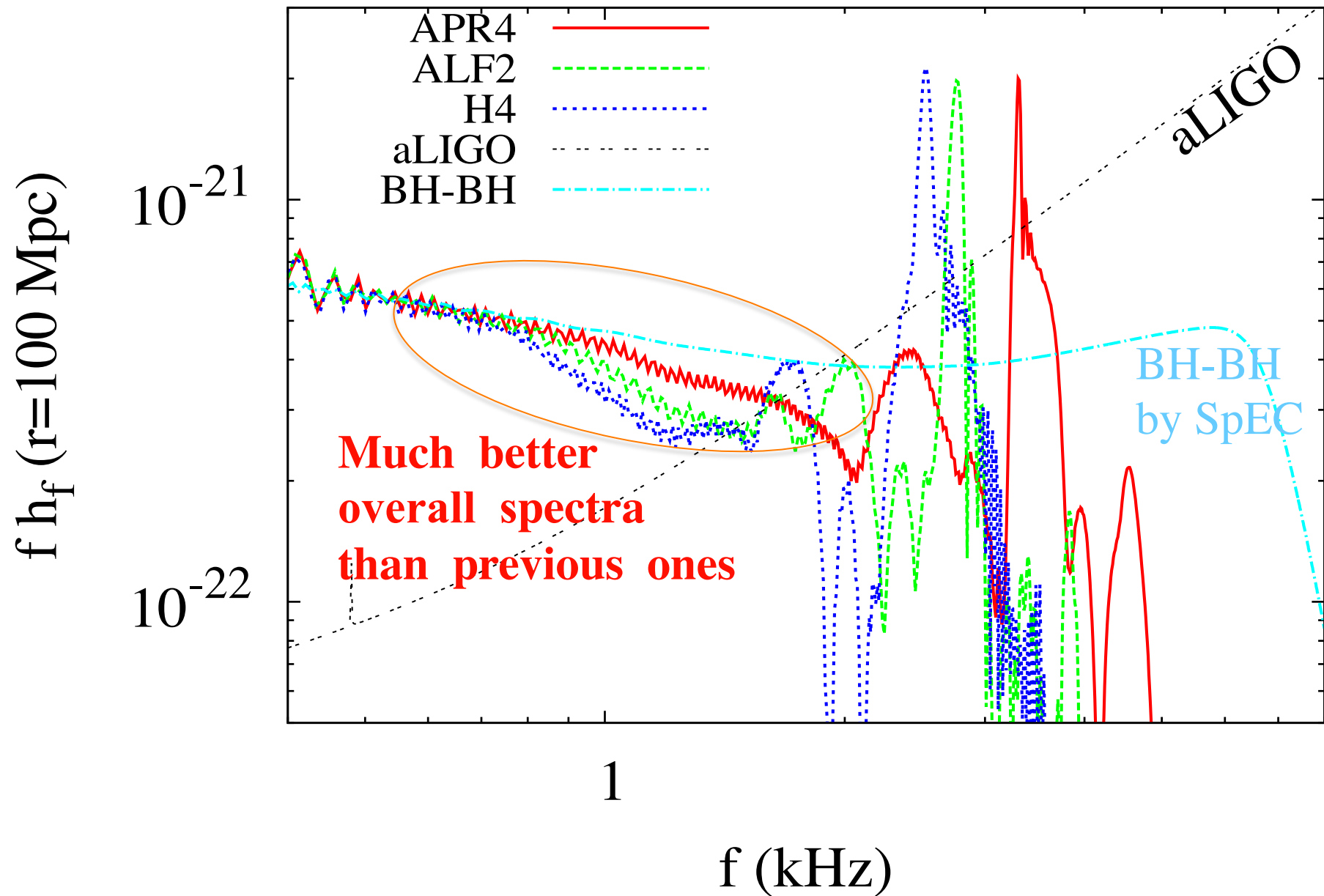


These data will be used for data analysis simulation (Talk by John)



Many simulations with a variety of tabulated EOS are ongoing.

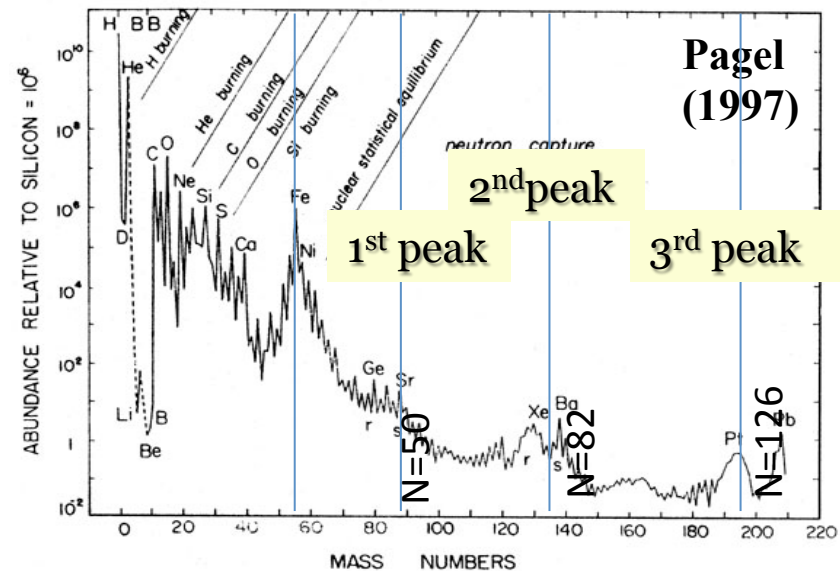
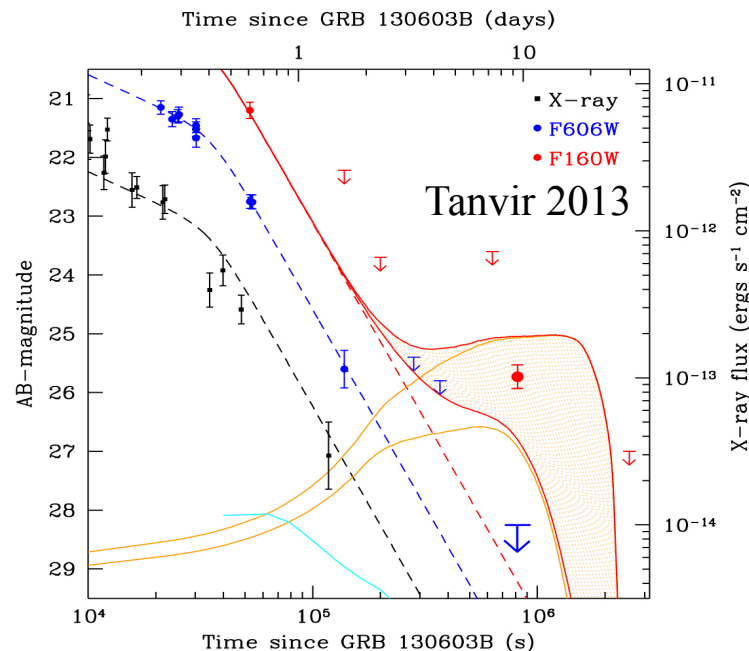
Overall spectrum



2A Dynamical mass ejection

- WHY important ?
 - That could shine and be an EM counterpart of GW source (radioactively-powered nova, radio flare, ...)
 - That could be main source of r-process elements

New topics in numerical relativity !



Dynamical mass ejection mechanism

2 major effects drive ejection

- 1) **Strong shock at the merger** → enhanced thermal pressure ejects material (like supernova)
 - 2) **Tidal torque by *non-axisymmetric merger remnant*** → Give angular momentum to the material in the envelope, subsequently ejected.
- ❖ Note that other effects like **magnetic** or **viscous** or neutrino wind could play a role (e.g., talks by Kiuchi, Just, ...)

Mass ejection at merger

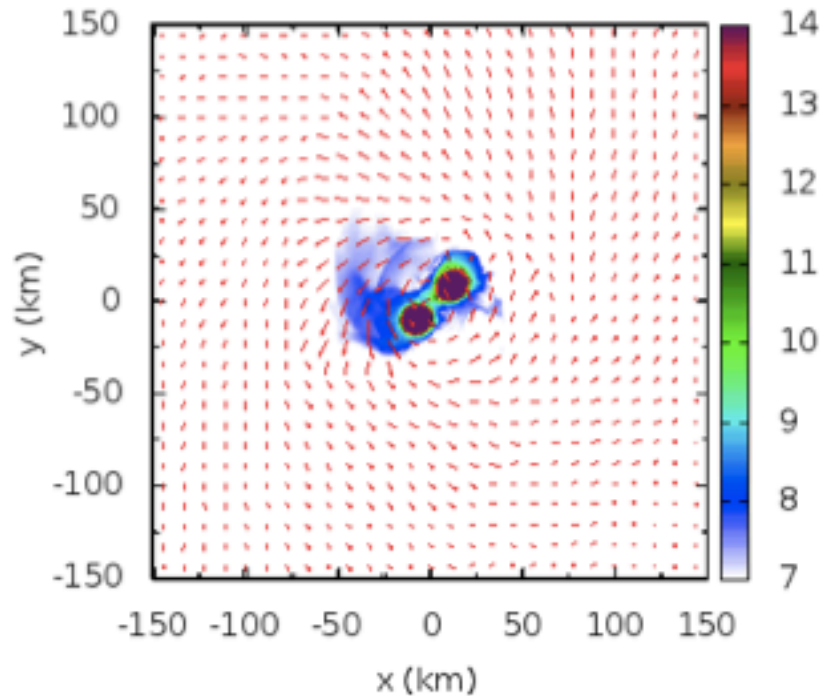
Model : $1.2M_{\text{sun}} - 1.5M_{\text{sun}}$, EOS=APR4, $R \sim 11$ km

Equatorial plane

300 km x 300 km

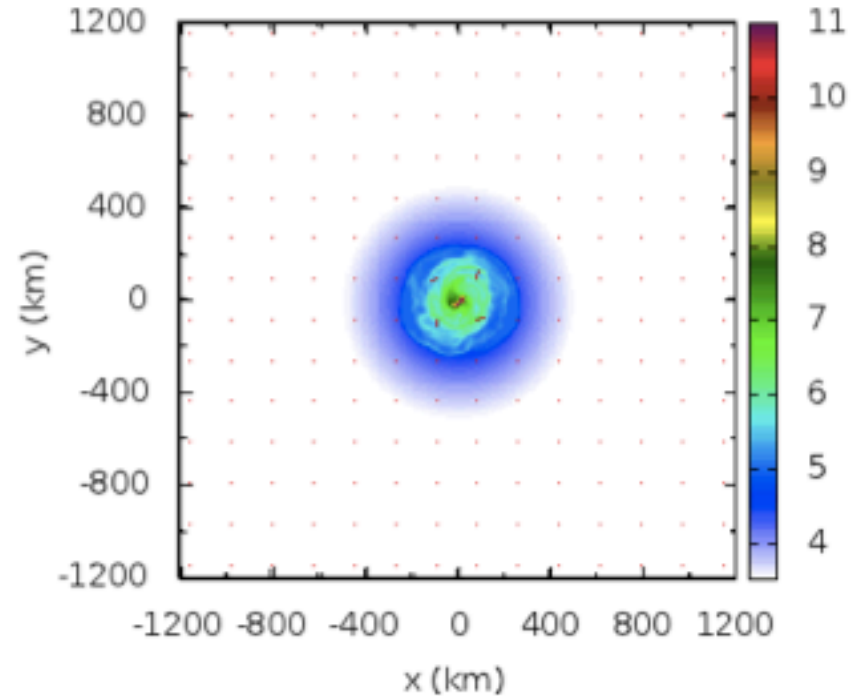
t=9.1854 ms

Log(ρ g/cc)



2400 km x 2400 km

t=9.1854 ms



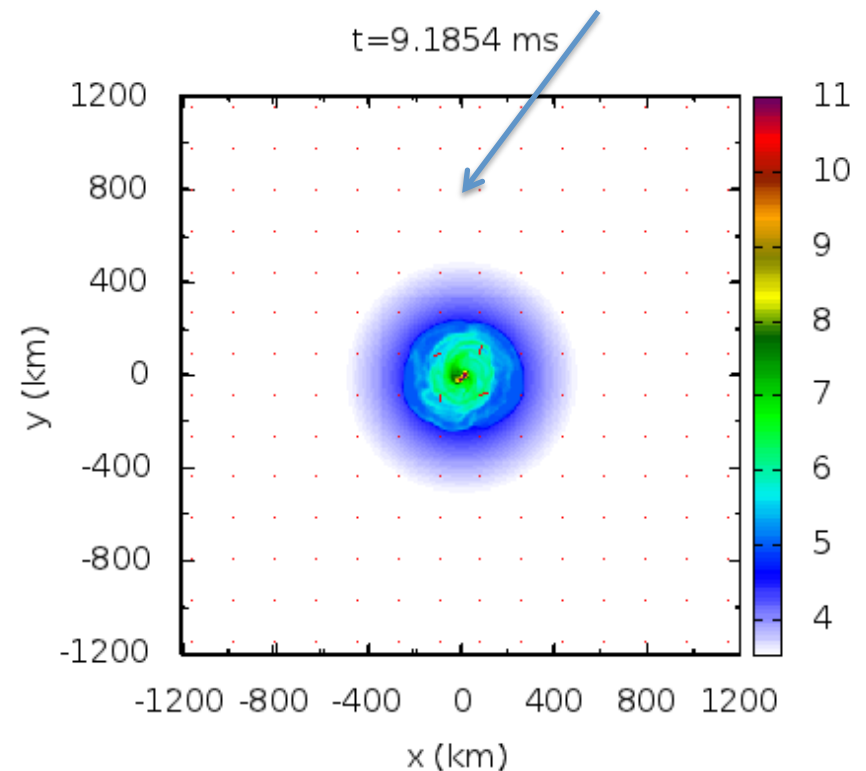
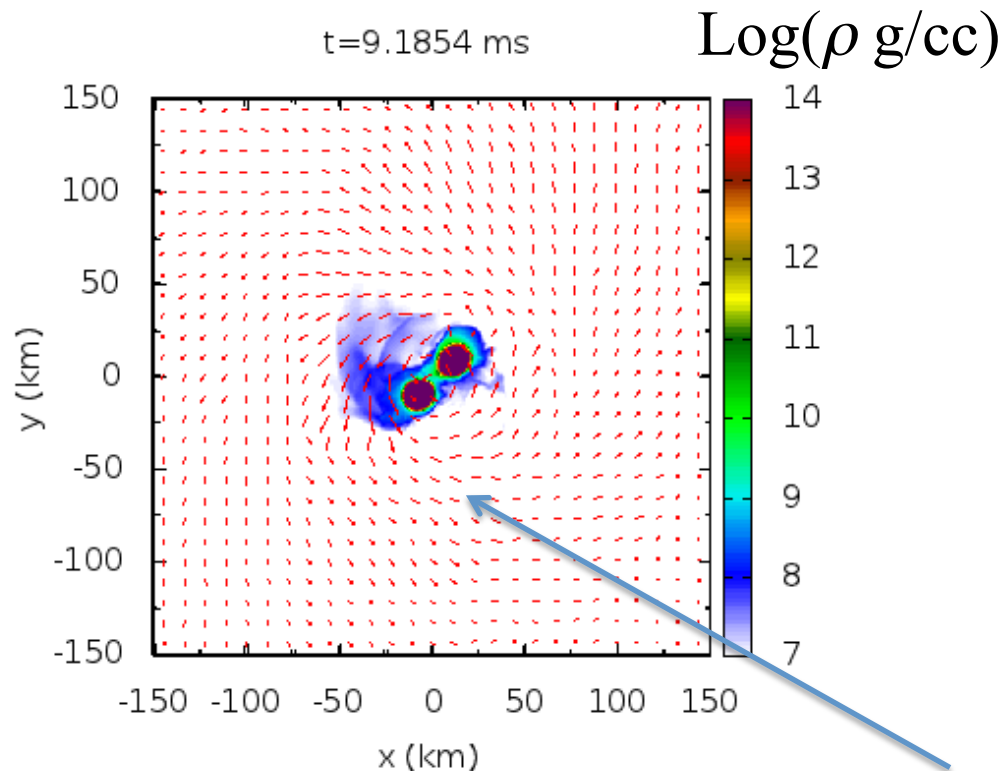
Mass ejection at merger

Model : $1.2M_{\text{sun}} - 1.5M_{\text{sun}}$, EOS=APR4, $R \sim 11$ km

Equatorial plane

$300 * 300$ km

Head speed $v \sim 0.8c$



Tidal torque plays an important role

Ejecta mass $\sim 0.01M_{\text{sun}}$, $v \sim 0.2c$ in average

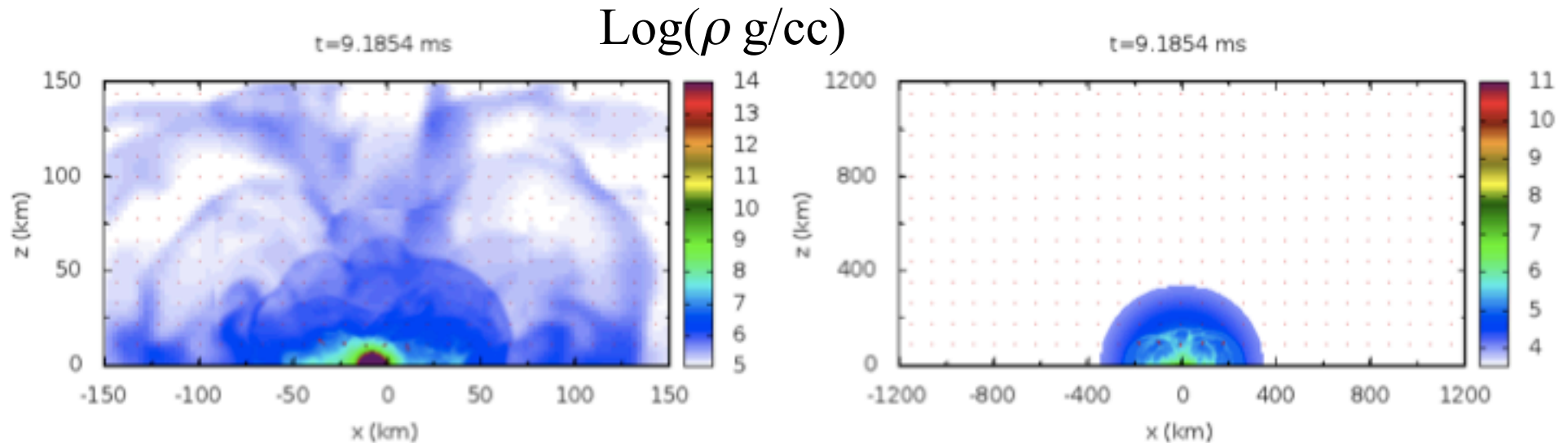
Mass ejection on the meridian plane

(x-z plane)

Model : $1.2M_{\text{sun}} - 1.5M_{\text{sun}}$, EOS=APR4, $R \sim 11$ km

300 km x 150 km

2400 km x 1200 km



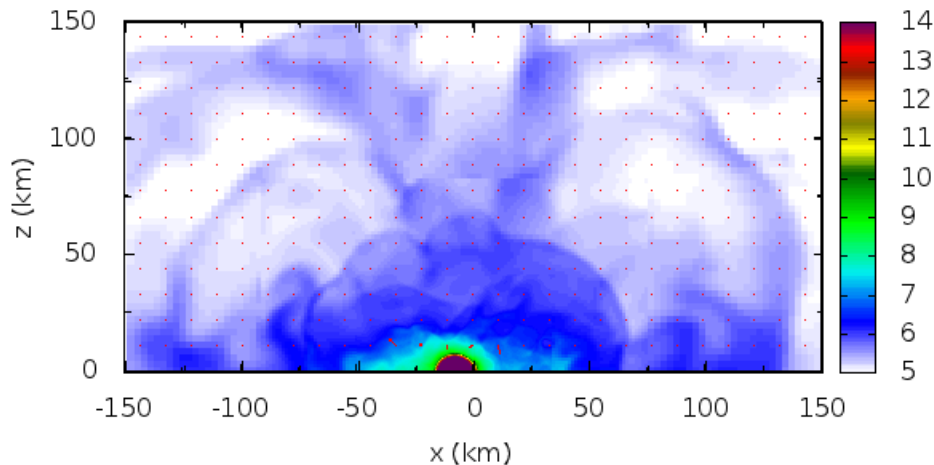
Mass ejection on the meridian plane

(x-z plane)

Model : $1.2M_{\text{sun}} - 1.5M_{\text{sun}}$, EOS=APR4, $R \sim 11$ km

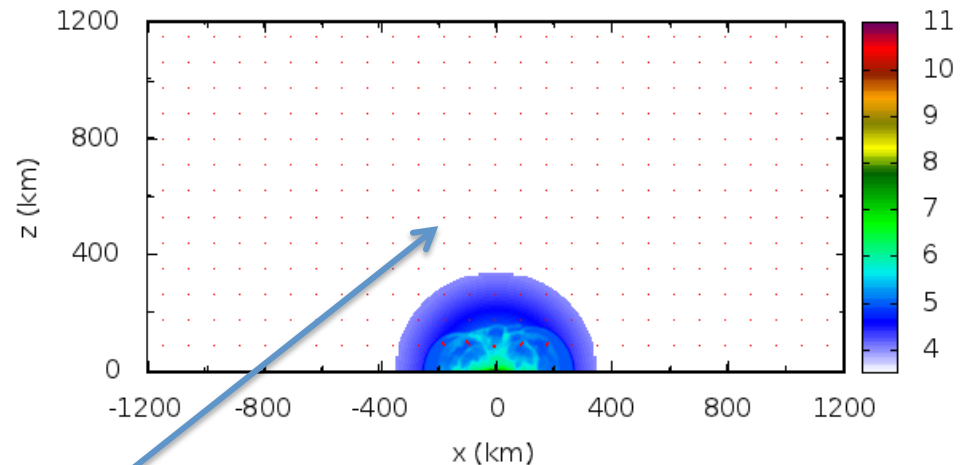
300 * 150 km

t=9.1854 ms



2400 * 1200 km

t=9.1854 ms



Ejecta is quasi-spherical:
Shock heating plays a key role.

Dynamical ejection mechanism

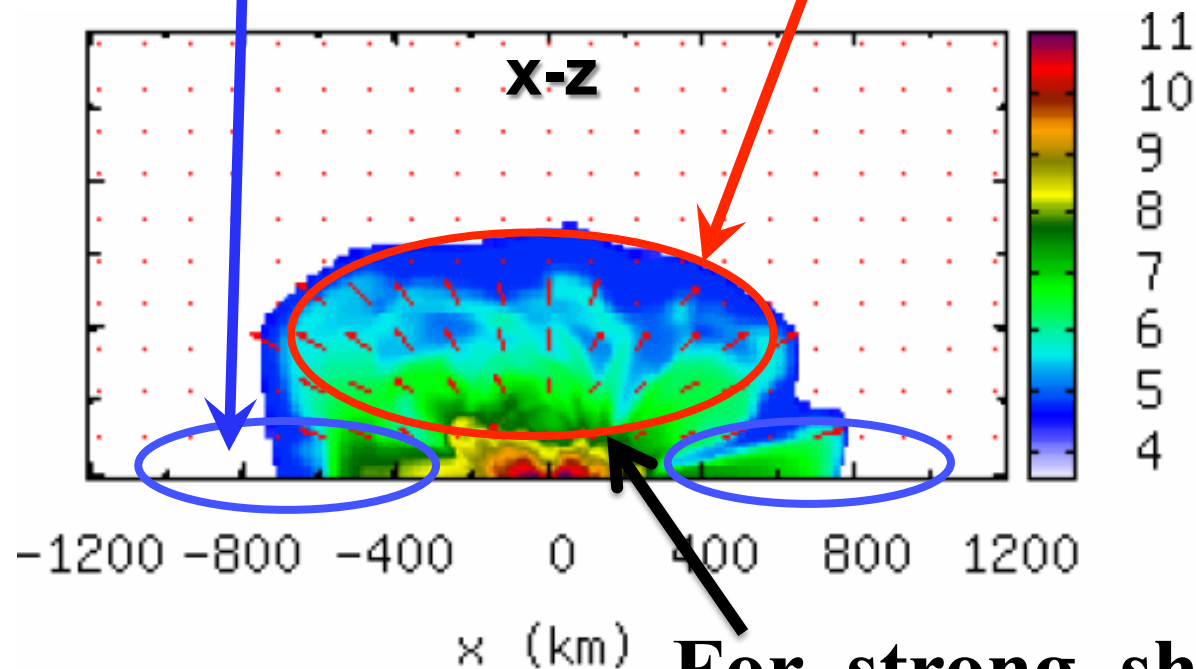
Two components

▶ Tidal component

▶ Low-temperature

▶ Shock-heated component

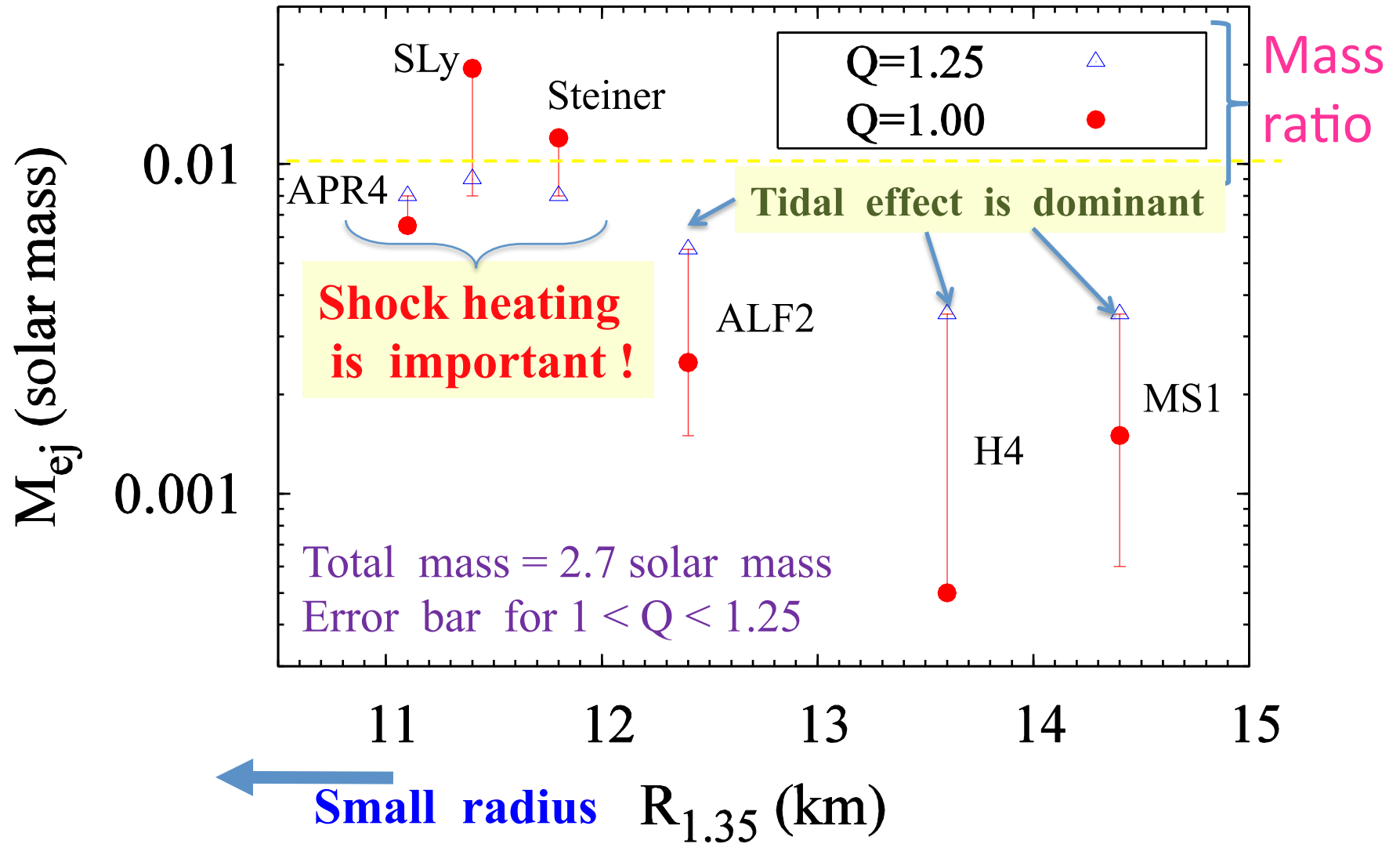
▶ High-temperature



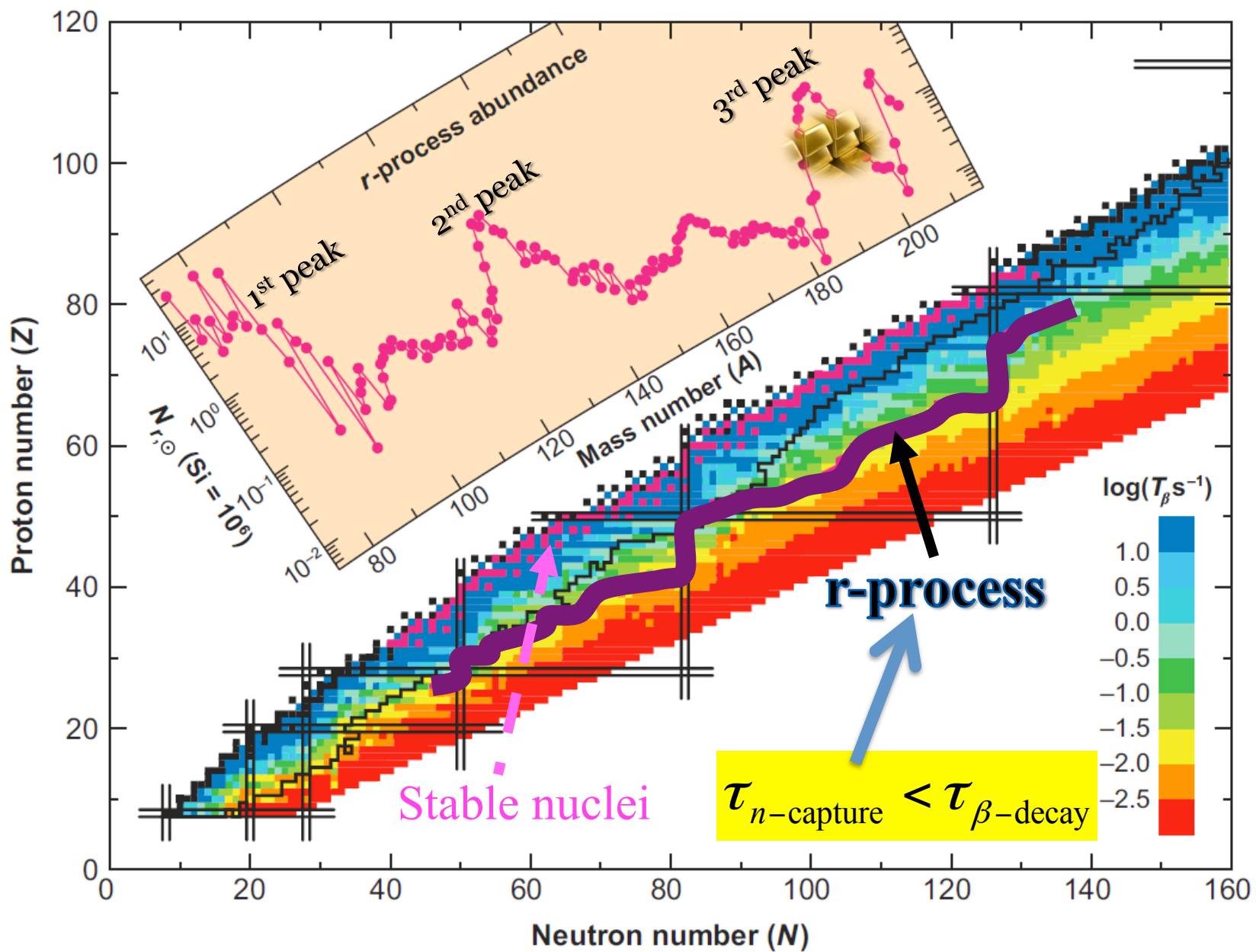
**For strong shock,
GR gravity is crucial**

Amount of ejecta depends strongly on EOS

Soft EOS \rightarrow strong gravity \rightarrow high-mass ejection



Proton number



Neutron number

Galactic r-process elements

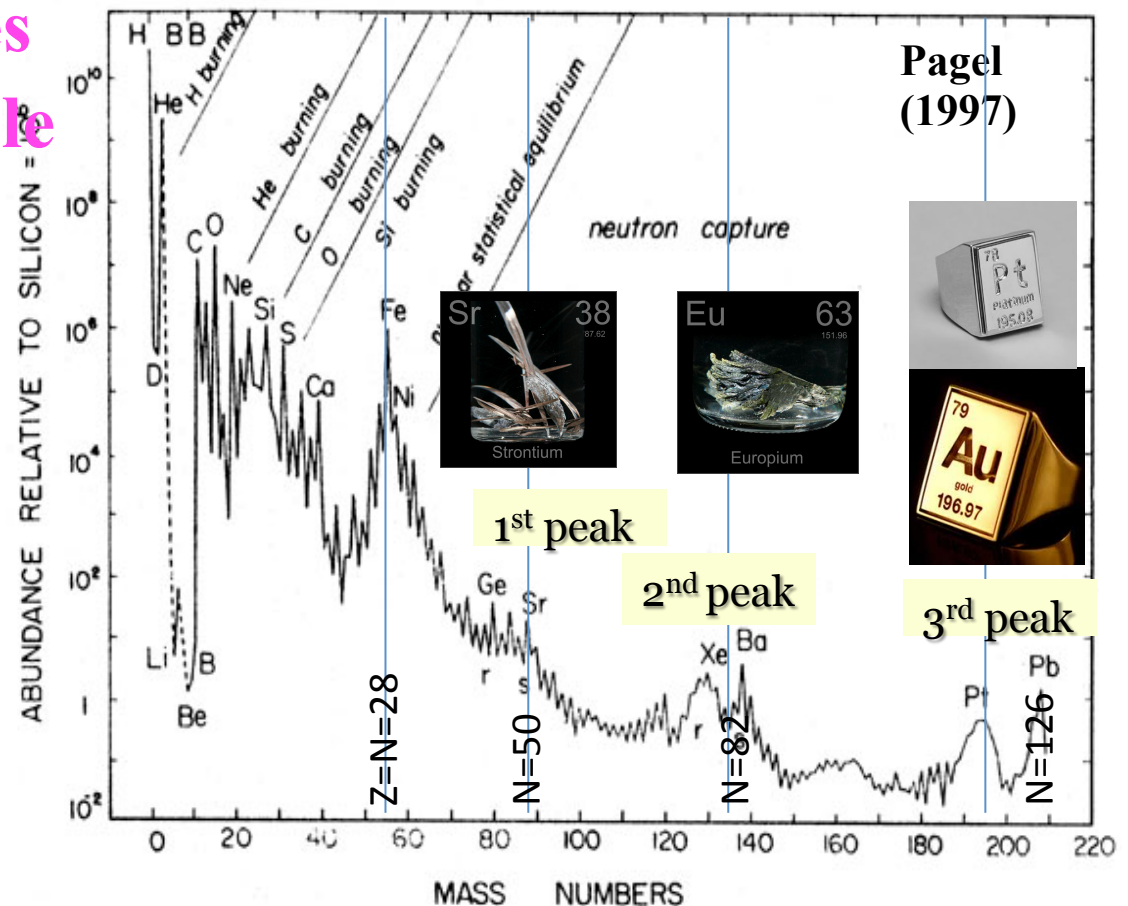
- Numerical-relativity simulations show **ejected mass per event of NS-NS could be $\sim 0.001-0.01 M_{\text{sun}}$**
- Total amount of observed r-process elements in our galaxy is **$\sim 10^4$ solar mass**
- Predicted merger rate \sim one every 10^4 yrs or less
→ total merger events $\sim 10^6$ or less in our Galaxy
- We want mass ejection per event $\sim 0.01 M_{\text{sun}}$
→ **If other contributions were absent, relatively soft EOS would be necessary**
- **IF EOS is stiff (NS has a large radius), we would need other sources or other mechanisms**

r-process nucleosynthesis study of ejecta

(By Sekiguchi & Wanajo +)

Universality of three peaks for heavy elements found in solar system & metal-poor stars

- Universality indicates the presence of single main origin
- **Question:**
Could NS-binary merger reproduce abundance pattern (all three peaks)?

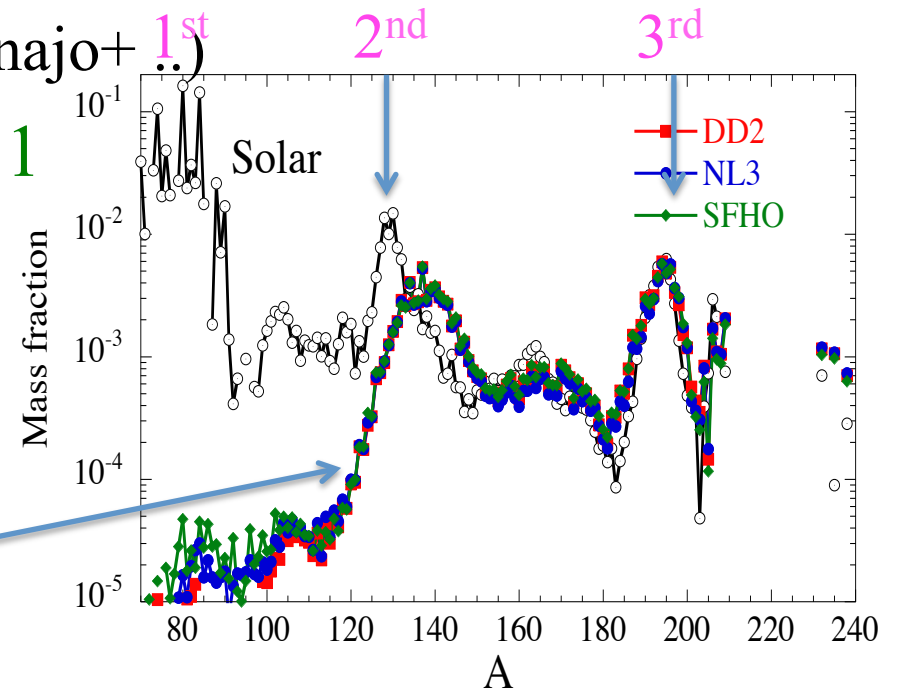


Key quantity for producing heavy elements:
electron fraction per baryon: $Y_e = [p]/([n]+[p])$

E.g., Korobkin+ (2012)

- ◆ If high $Y_e > 0.45$ = neutron less-rich
→ 3rd peak is not well reproduced
(e.g., CCSN, Roberts, Janka+, Wanajo+ (1st))

- ◆ If too neutron rich, $Y_e \sim 0.1$
e.g., BH-NS or NS-NS in
Newtonian simulation
→ 3rd peak dominant;
no/weak 1st & 2nd peaks
e.g., Goriely et al. 2011



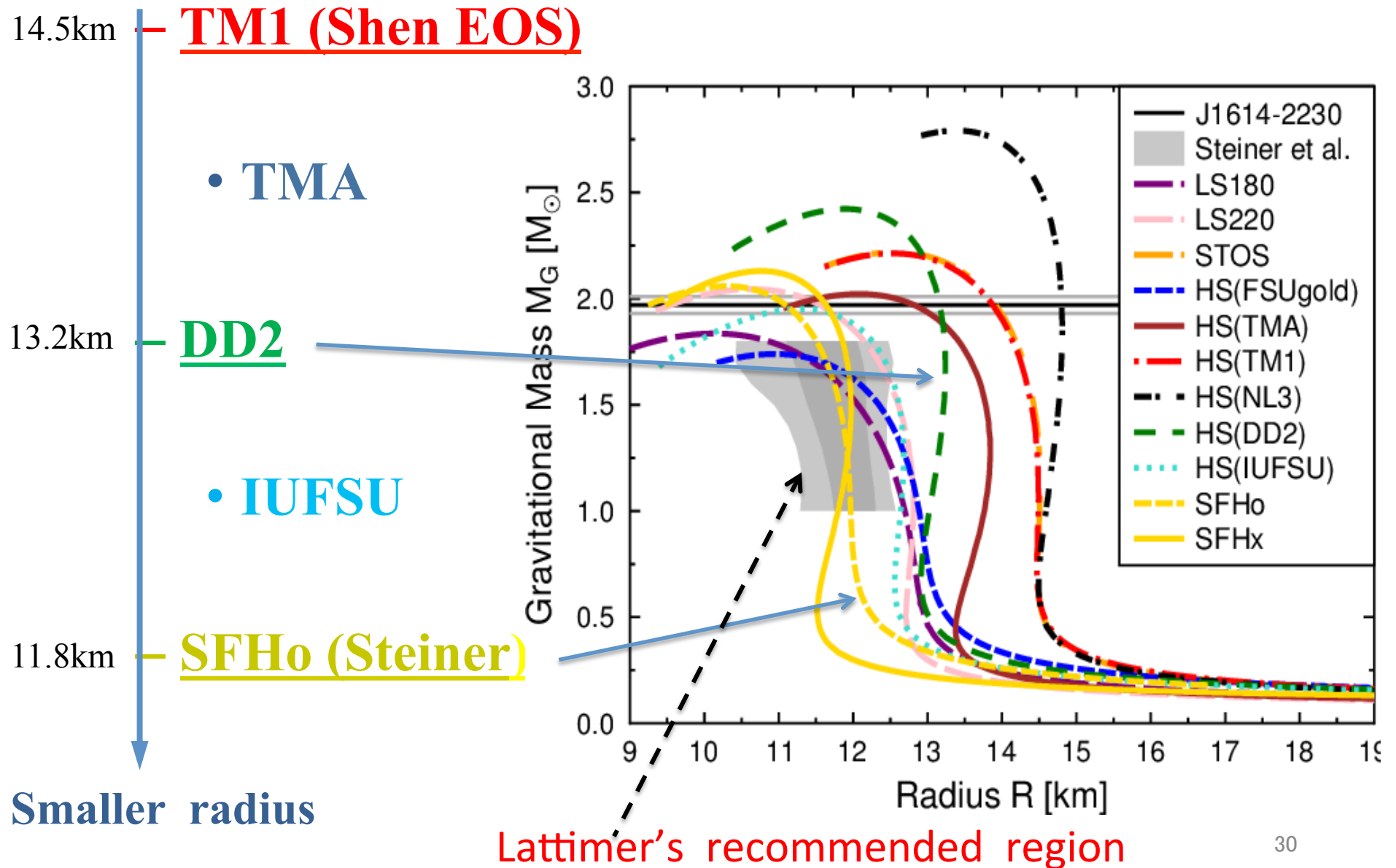
- **Appropriate blending of Y_e is needed: HOW ?**
→ **Perform numerical relativity simulation !**

GR neutrino-radiation hydrodynamics

(Sekiguchi's GR radiation hydro code)

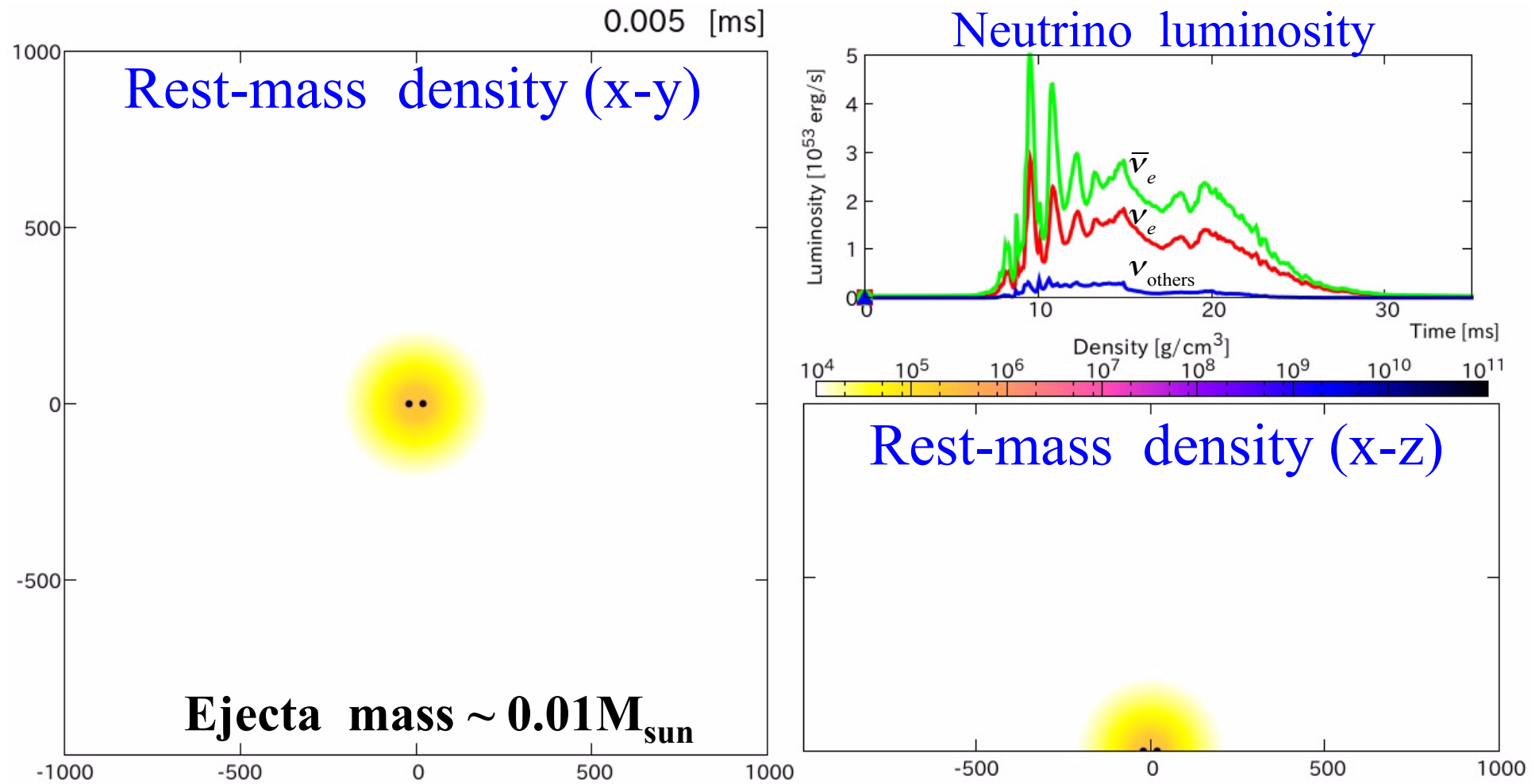
- Einstein's eq: BSSN + puncture (+ local Z4c)
- Radiation: Leakage + fully covariant truncated moment scheme with M1 closure (gray) for heating
pure M1 scheme (gray) works but expensive
- EOS: SFHo, IUFSU, DD2, TMA, TM1
- Grid size: 580*580*290*9 level (fixed mesh refinement) with $\Delta x=150\text{--}160$ m for the finest domain
- CPU time: 500-700k node-hours by K-computer with ~7000 cores (864 nodes)
- Binary mass: 1.30-1.30, 1.35-1.35, 1.30-1.40, 1.25-1.45, 1.40-1.40 (ongoing)

Variety of EOS table (we appreciate Hempel)



SFHo ($R \sim 11.9$ km): $1.35-1.35 M_{\text{sun}}$

ρ

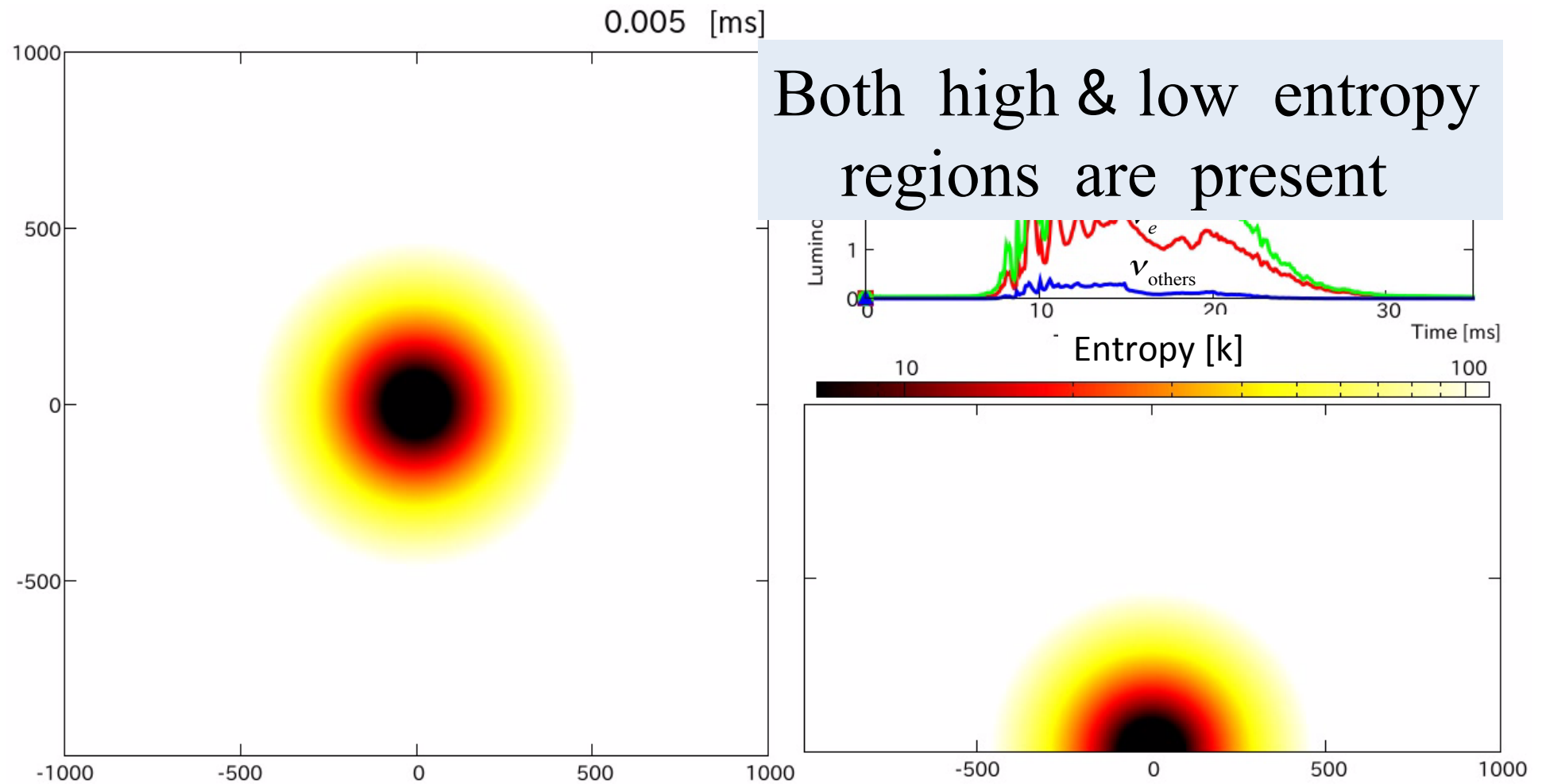


Sekiguchi et al. (2015)

SFHo ($R \sim 11.9$ km): $1.35 - 1.35 M_{\text{sun}}$

S

Specific entropy

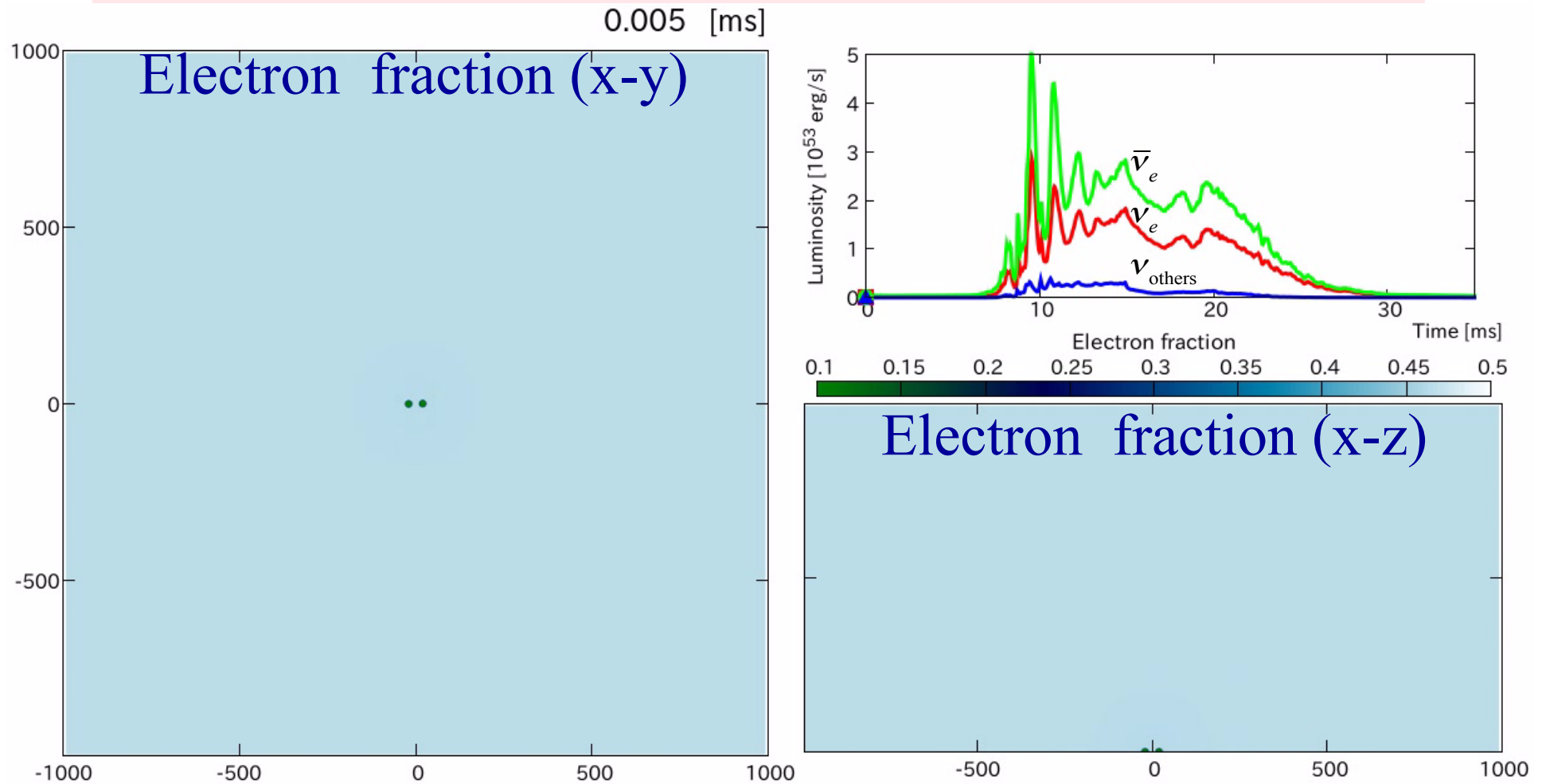


Sekiguchi et al. (2015)

SFHo (R~11.9 km): 1.35-1.35 M_{sun}

High temperature $\Rightarrow \gamma\gamma \rightarrow e^- + e^+$, $n + e^+ \rightarrow p + \bar{\nu}_e$

Y
e



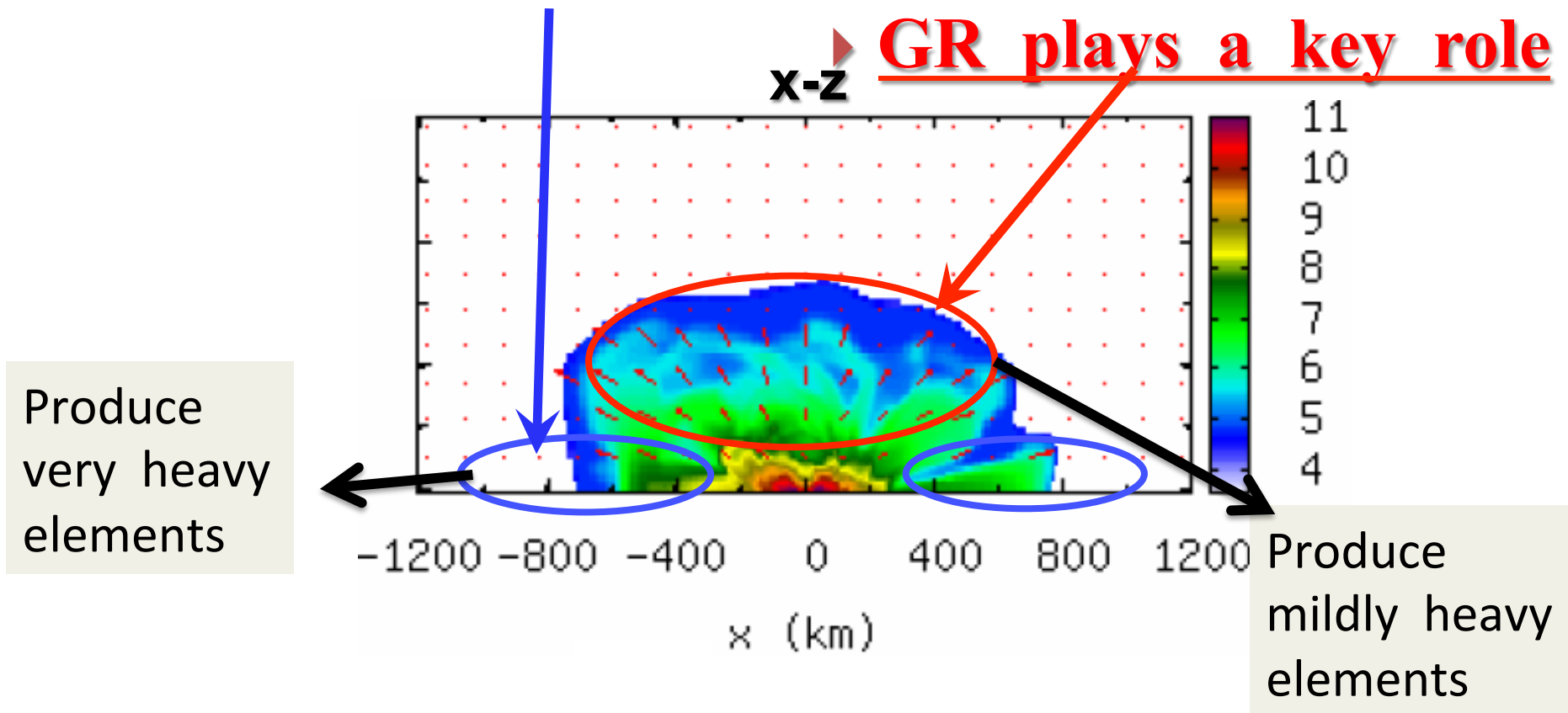
Thermodynamical properties of ejecta

Mass ejection from BNS merger : two components

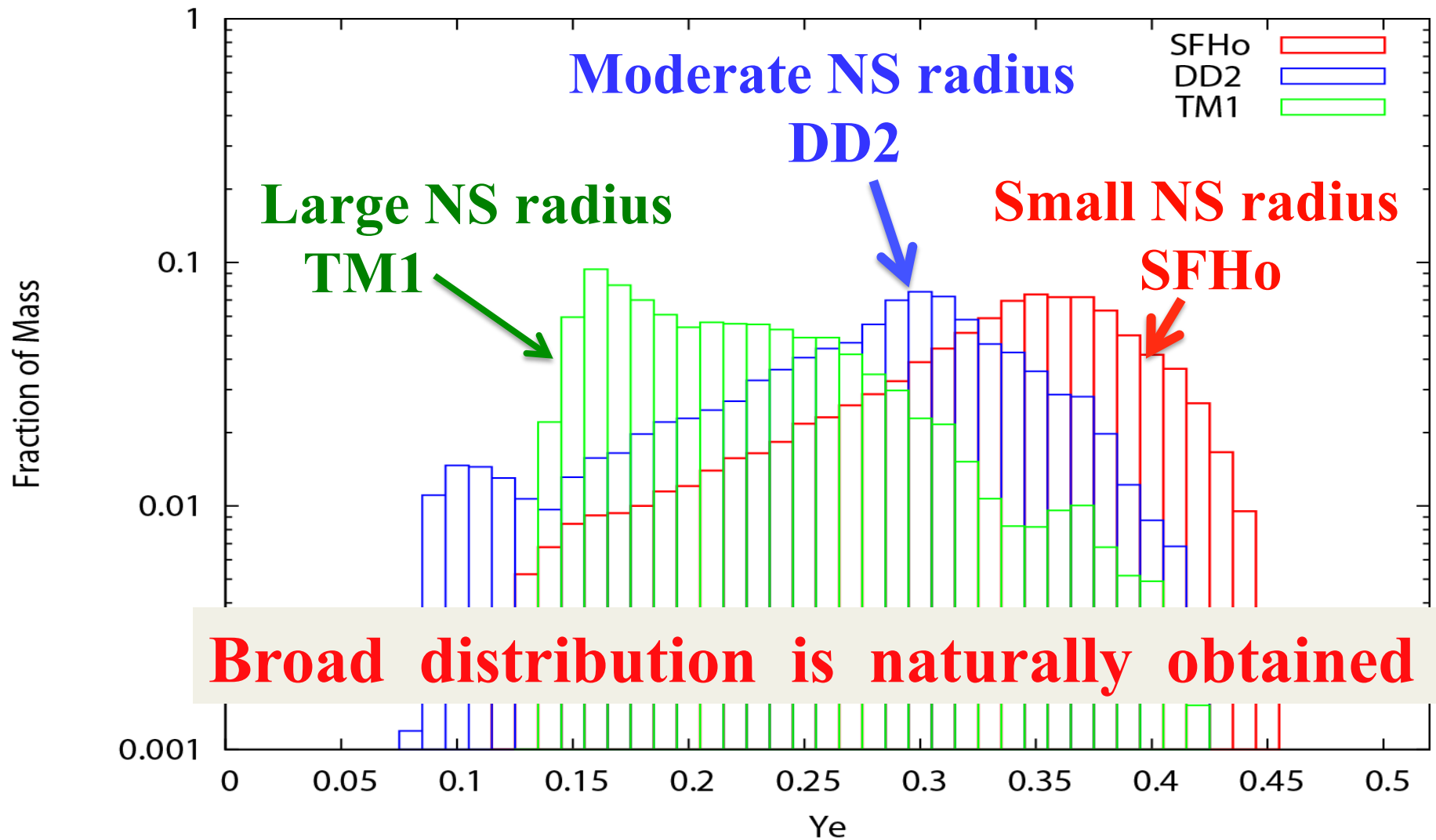
- ▶ Tidal component
- ▶ Low-temp, low Y_e

- ▶ Shock-heated component
- ▶ High-temp, high Y_e

▶ GR plays a key role

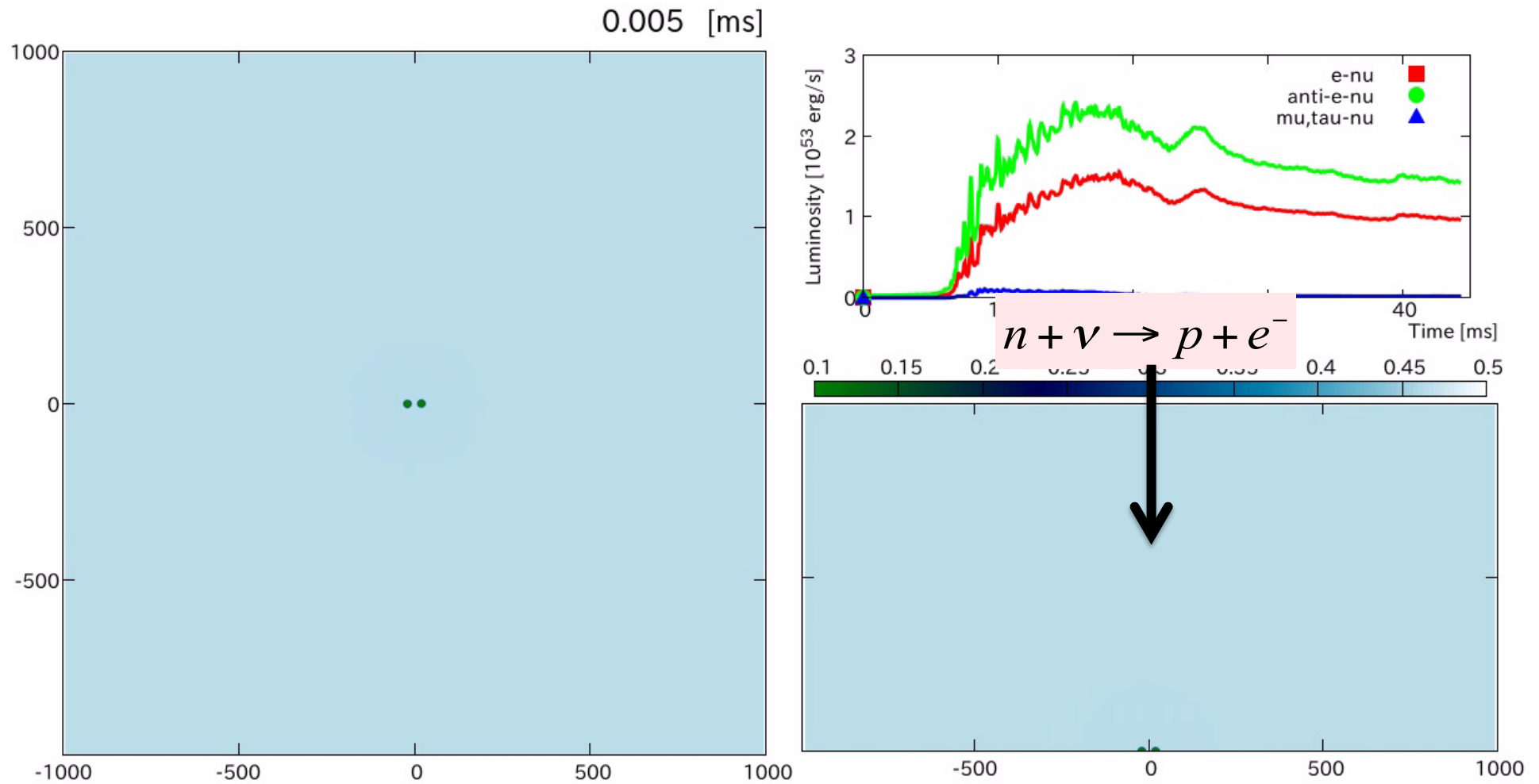


Fraction of mass as a function of Y_e

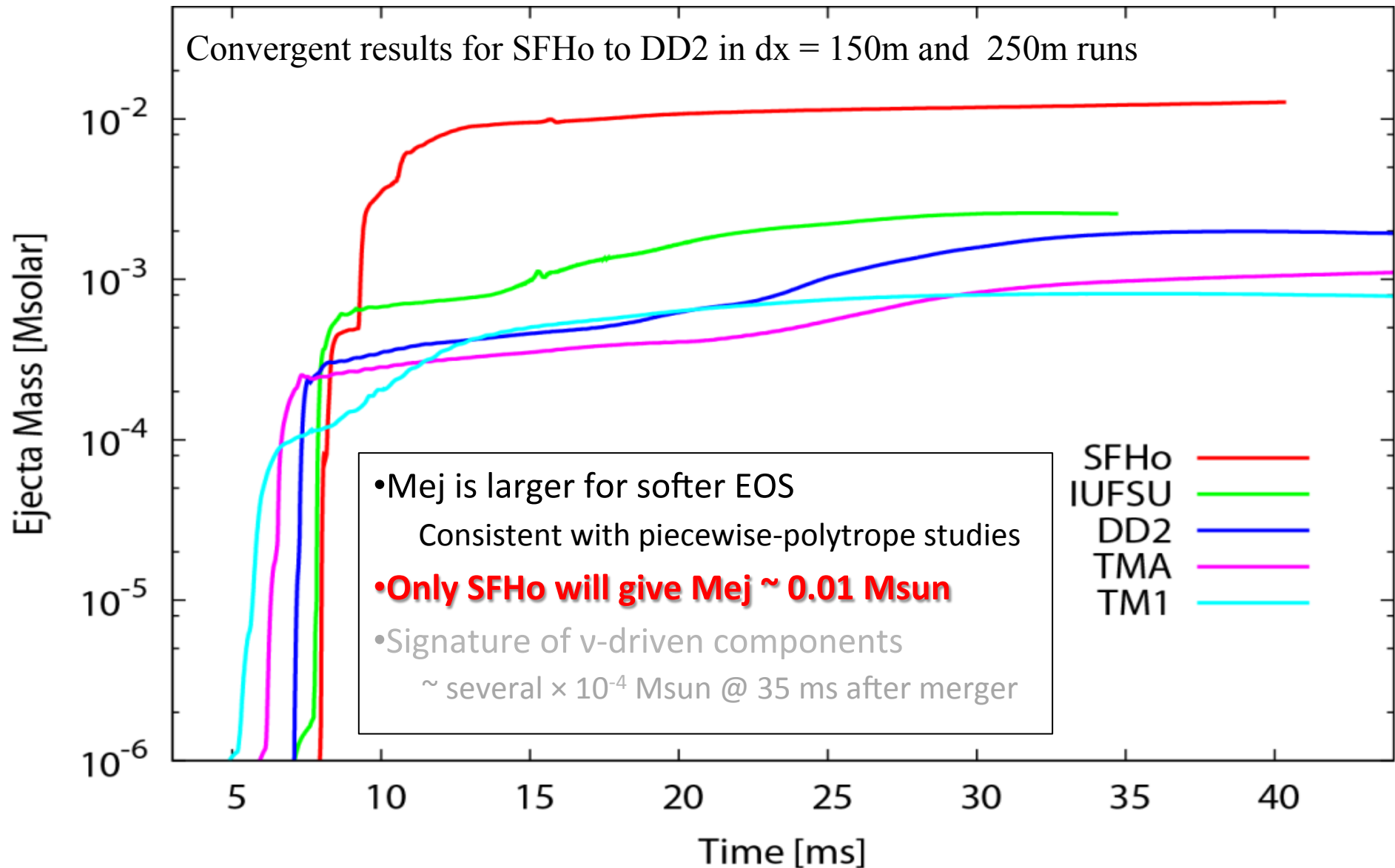


Sekiguchi et al. (2015)

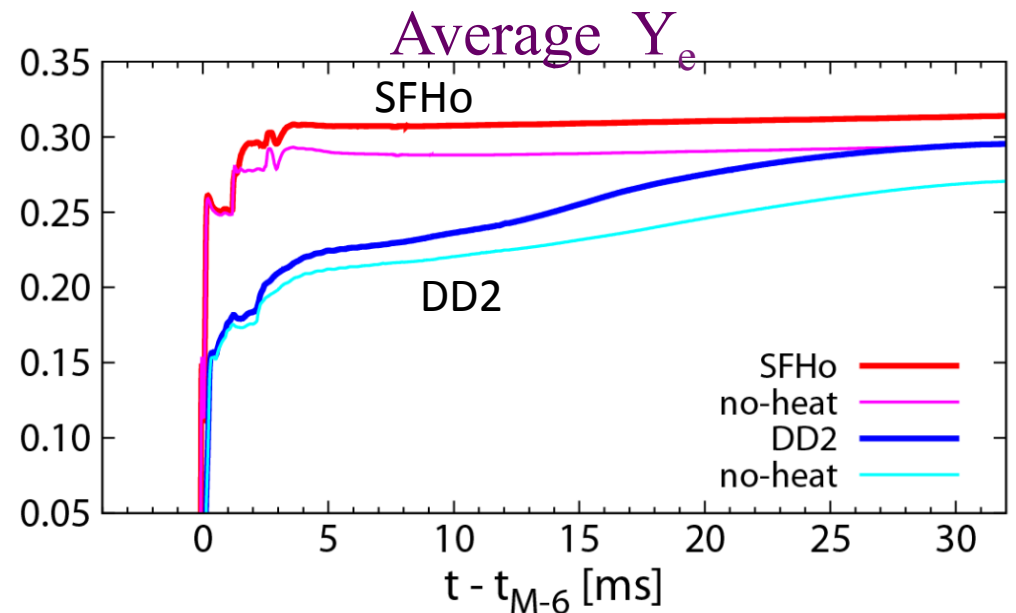
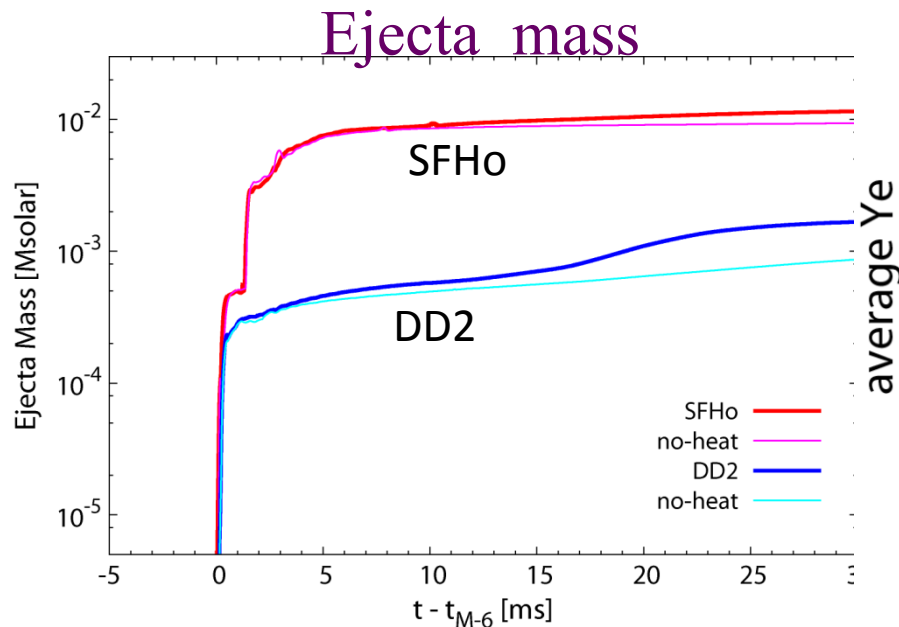
DD2 ($R \sim 13.2$ km): $1.35-1.35 M_{\text{sun}}$ **Y**
 Neutrino heating from remnant MNS is important! **e**



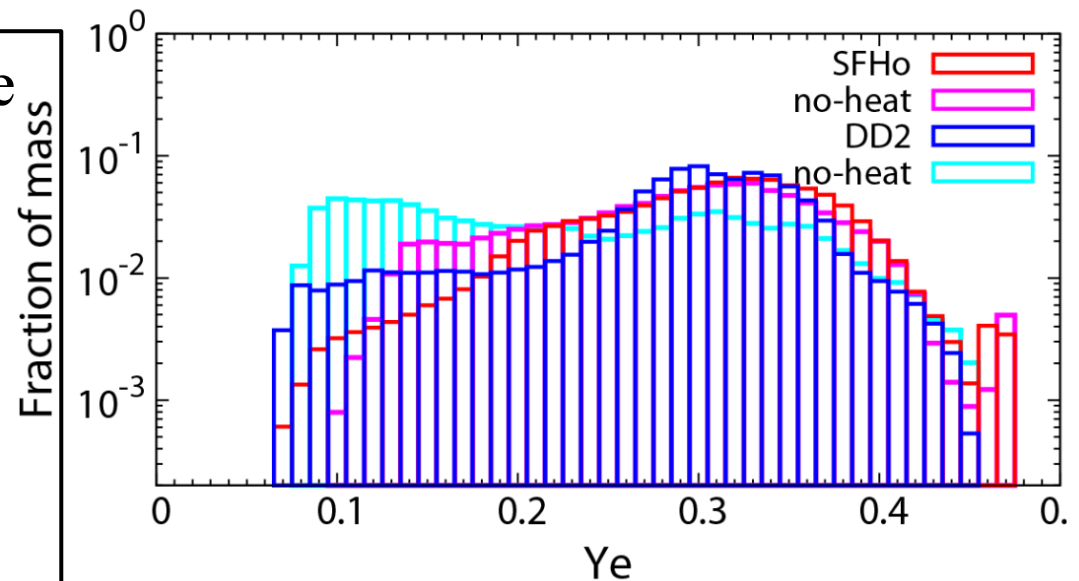
However, for stiff EOS, ejecta mass is small



Effects of neutrino heating

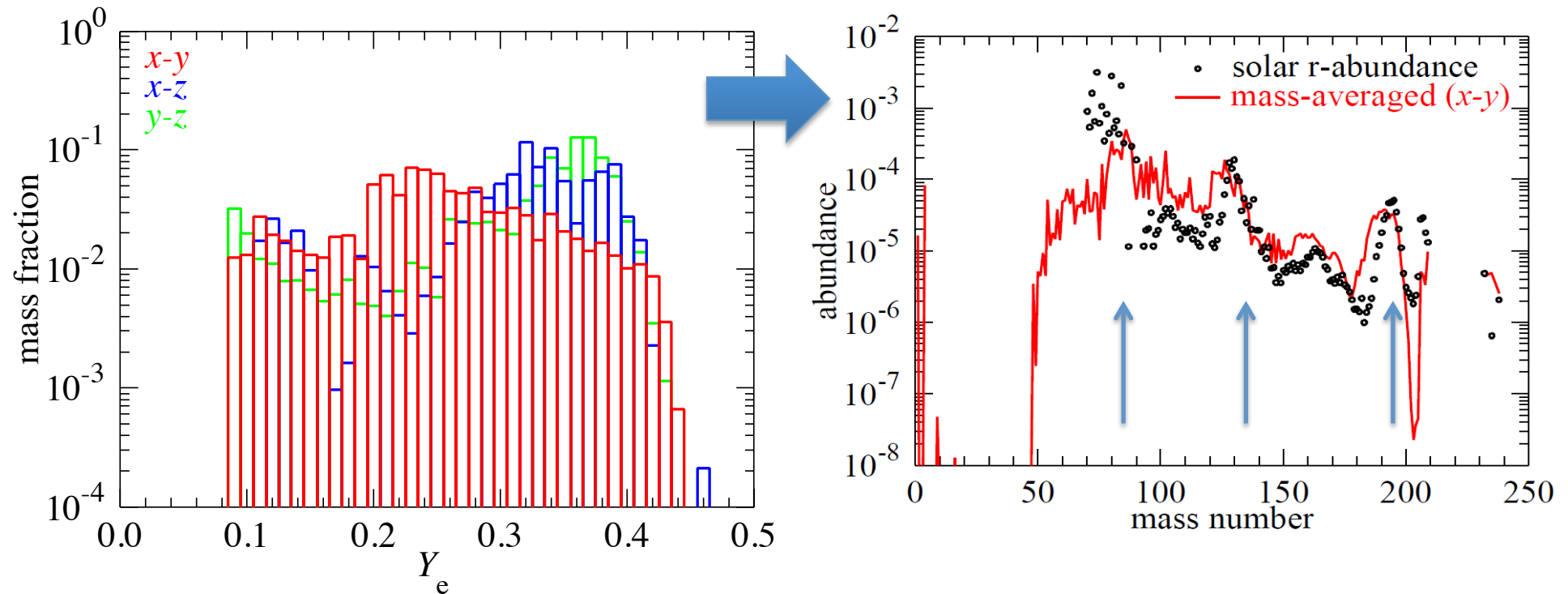


- ▶ Amount of ejecta mass can be increased by $\sim 10^{-3} M_{\text{sun}}$
- ▶ Average Y_e can change by 0.02~0.03
- ▶ For DD2 & TM1, ejecta mass is $O(0.001 \text{ solar})$
- ▶ Viscous heating ??



Our first result

(Wanajo et al. ApJ 2014)



Broad distribution for Y_e could be suitable for reproducing wide abundance pattern

Project is ongoing by Wanajo, Nishimura, Sekiguchi+

Summary

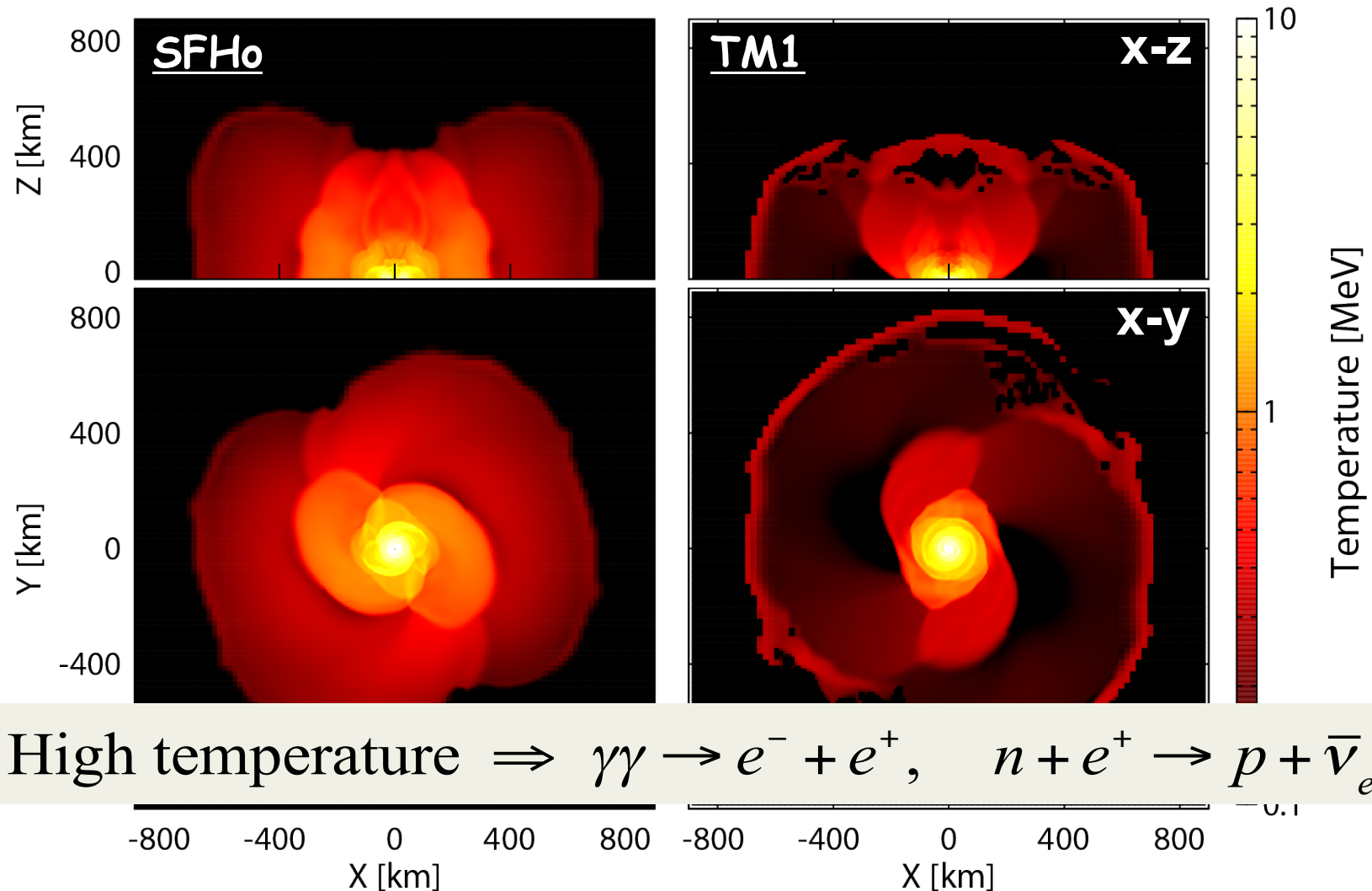
- Gravitational waves from late inspiraling phase of NS-NS is a valuable site for exploring NS EOS
→ high-resolution numerical-relativity simulations are ongoing for constructing templates
(also by Bernuzzi-Nagar +, Haas +,)
- Mass ejected in NS-NS merger is $\sim 0.001-0.01$ solar mass → EM counterparts (tomorrow's talks)
- NS-NS could be r-process nucleosynthesis site:
Three peaks could be well reproduced by shock + neutrino heating (Sekiguchi + 2015; works ongoing)
- Next issue: Adding viscous effects to remnant NS and/or BH+torus

Announcement from Yukawa Institute, Kyoto University

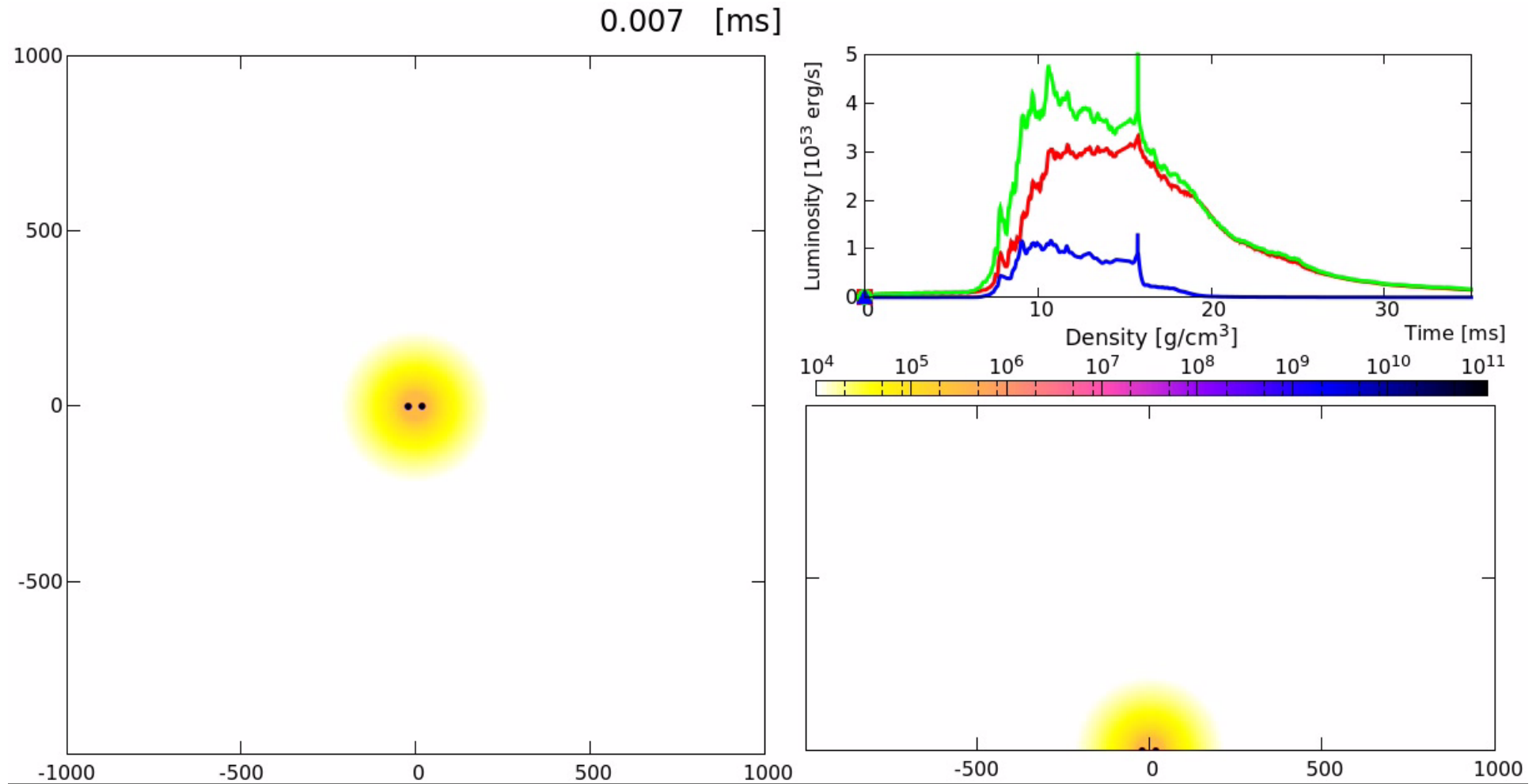
- **Longterm workshop on “Nuclear Physics and Compact Stars 2016 (NPACS 2016)”**
Oct.17 (Mon.), 2016 -- Nov.18 (Fri.), 2016.
- **In the third week, conference on “Birth, Life, and Death of Neutron Stars and Nuclei (YKIS 2016)” will be held**
Oct.31 (Mon.), 2016 -- Nov.4 (Fri.), 2016

Ejecta temperature: Depends on EOS

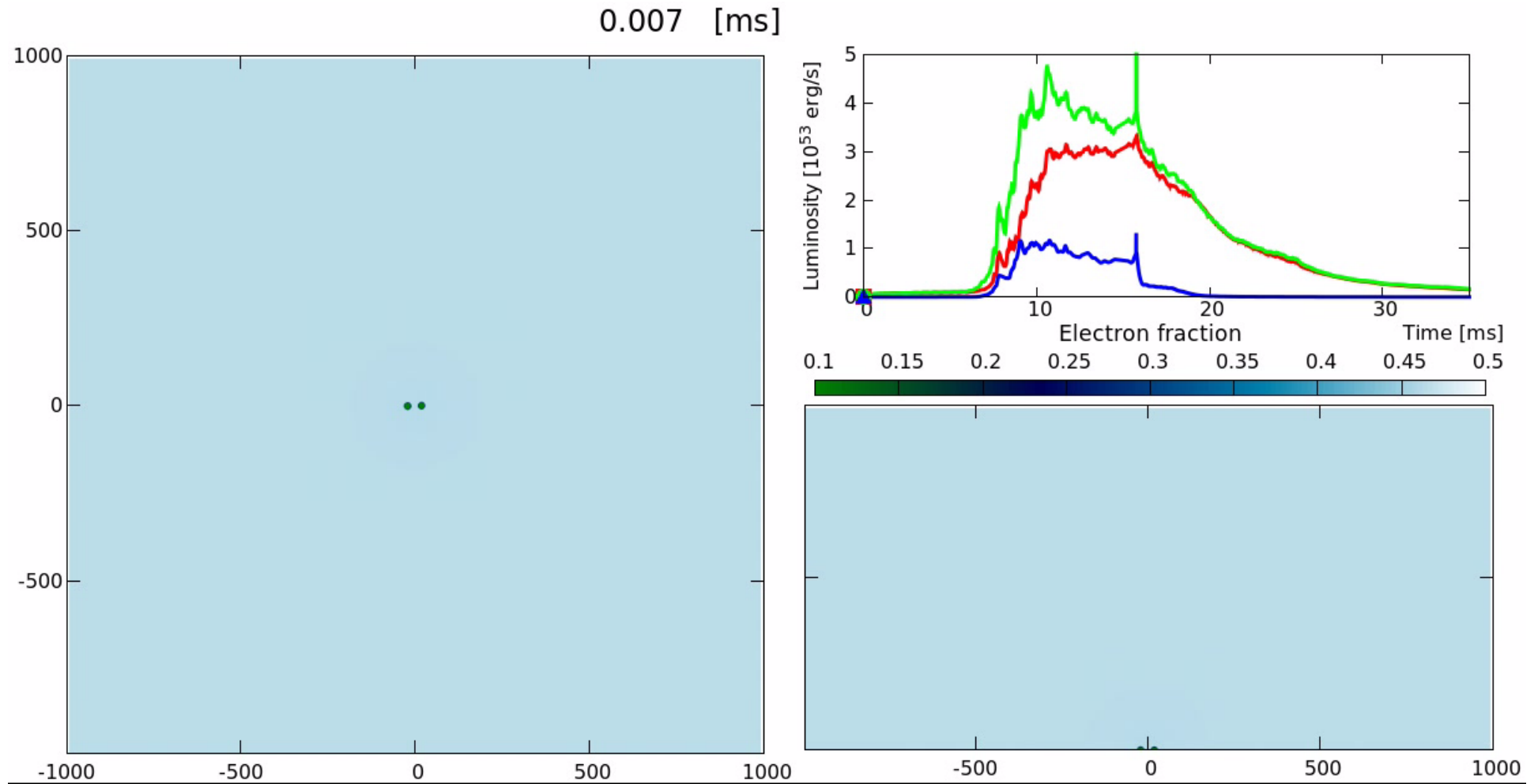
- SFHo EOS: NS=Small radius
- High temperature
- ▶ TM1 EOS: NS=Large radius
- ▶ Low temperature \rightarrow n rich



Unequal mass NS-NS system: SFHo1.25-1.45



Unequal mass NS-NS system: SFHo1.25-1.45



Unequal mass NS-NS system: SFHo1.25-1.45

- Orbital plane : Tidal effects play an important role, ejecta is neutron rich
- Meridian plane : shock + neutrinos play roles, ejecta less neutron rich

