

Pythagoras Program Ministry of Education Hellenic Republic

Numerical studies on the excitation of magnetosonic waves by a gravitational wave



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Collaborators

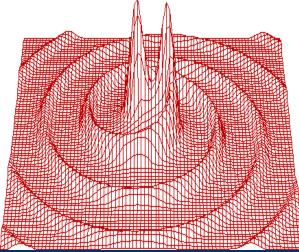
Prof. Vlahos Dr. Isliker



- Motivation
- Background
- Model of equations
- Linear limit
- Numerical results
- Conclusions
- Future Plans

Motivation

 Investigate numerically plasma wave phenomena driven by GW
Diagnostic tool for indirect detections of GW...



Basic Assumptions

- 1+3 orthonormal frame
- TT gauge
- Orthogonal frame

- COVARIANT FORM
- Strong magnetic field $B_0 = B_0 e_x$
- + polarized GW in Minkowski space-time
- constant GW amplitude

Geometry

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J. Moortgat and J. Kuijpers: Gravitational and magnetosonic waves in GRBs

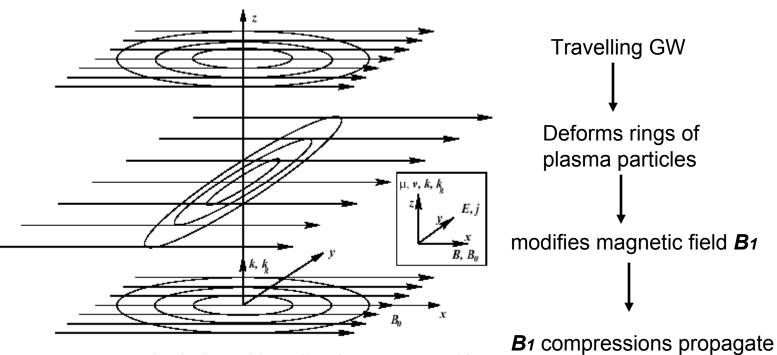


Fig. 1. A GW propagating in the positive *z*-direction across an ambient magnetic field (in the *x*-direction) excites a MSW. The orientations of the MSW components are indicated in the inset.

Model of Equations

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} - \vec{J}_B$$

$$\nabla \times \vec{B} = \frac{4\pi}{c}\vec{J}_m + \vec{J}_E + \frac{1}{c}\frac{\partial \vec{E}}{\partial t}$$

$$\vec{J}_E = -\frac{B_0}{2} \frac{\partial h}{\partial z} \hat{e}_y$$
$$\vec{J}_B = -\frac{B_0}{2c} \frac{\partial h}{\partial t} \vec{e}_x$$

Model of Equations II

$$\rho\left(\frac{\partial \vec{u}}{\partial t} + (\vec{u} \cdot \nabla)\vec{u}\right) = -\nabla P + \frac{1}{c}(\vec{J}_m \times \vec{B}) + \nu \nabla^2 \vec{u}$$

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \vec{u}) = 0$$

$$\vec{E} = -\frac{1}{c}(\vec{u} \times \vec{B}) + \eta \vec{J}_m$$

$$\nabla P = c_s^2 \nabla \rho$$
, where $c_s^2 \equiv \gamma P_0 / \rho_0$

One dimensional case

$$\rho(z,t)\left(\partial_t u_z(z,t) + \frac{1}{2}\partial_z u_z^2(z,t)\right) = -c_s^2 \partial_z \rho(z,t) - \frac{1}{c}J_{my}(z,t)B + \nu \partial_{zz} u_z(z,t)$$

$$E_y(z,t) + \frac{B}{c}u_z(z,t) = \eta J_{my}(z,t),$$

$$\begin{split} \partial_t \rho(z,t) &+ \partial_z (\rho(z,t) u_z(z,t)) = 0. \\ \partial_t E_y(z,t) &= c \partial_z B_x(z,t) - 4\pi J_{my} - c J_E \\ \partial_t B_x(z,t) &= c \partial_z E_y(z,t) - c J_B \end{split} \overset{\textbf{z} \ \texttt{k,kg}}{\underbrace{ \begin{array}{c} \forall \textbf{k,kg} \\ \textbf{y} \end{array}}} u(z,t), \textbf{J}(z,t) \\ \underbrace{ \begin{array}{c} \forall \textbf{y} \end{array}}_{\textbf{k}} B_{\textbf{b},\textbf{B}(z,t)} \\ \underbrace{ \begin{array}{c} \forall \textbf{y} \end{array}}_{\textbf{k}} B_{\textbf{k}} B_{\textbf{k}}(z,t) \\ \underbrace{ \begin{array}{c} \forall \textbf{y} \end{array}}_{\textbf{k}} B_{\textbf{k}} B_{\textbf{k$$

Linear limit

Linearize and keep First Order perturbed quantities

 $A=A_0 + A_1 * exp(kz-\omega t)$

$$\begin{split} \left[\omega^{2}\left(1+\frac{v_{A}^{2}}{c^{2}}\right)-k^{2}(c_{s}^{2}+v_{A}^{2})+i\frac{\eta\omega c^{2}}{4\pi}\left(k^{2}-\frac{\omega^{2}}{c^{2}}\right)\left(1-\frac{c_{s}^{2}k^{2}}{\omega^{2}}+i\frac{k^{2}\nu}{\rho_{0}\omega}\right)+i\frac{k^{2}\nu\omega}{\rho_{0}}\right]E_{y} = \\ \left[-B_{0}k^{2}v_{A}^{2}h_{0}+i\eta\frac{B_{0}k^{2}h_{0}c^{2}\omega}{4\pi}\left(1-\frac{c_{s}^{2}k^{2}}{\omega^{2}}+i\frac{k^{2}\nu}{\rho_{0}\omega}\right)\right]\delta(k-k_{g})\delta(\omega-\omega_{g}) \end{split}$$

where $v_A = B_0 / \sqrt{4\pi\rho_0}$.

$$\begin{split} |E_y|^2 &= (B_0 k_g^2 u_A^2 h_0)^2 \frac{1 + \left(\eta \frac{\omega c^2}{4\pi v_A^2}\right)^2}{\left[\omega^2 - k^2 u_A^2\right]^2 + \left[\frac{\eta \omega c^2}{4\pi} \left(k^2 - \omega^2/c^2\right) + \frac{k^2 \nu \omega}{\rho_0}\right]^2 \left(1 + \frac{v_A^2}{c^2}\right)^{-2}} \end{split}$$

where $u_A^2 = v_A^2/[1 + v_A^2/c^2]$ is the relativistic Alfvén velocity.

Numerical solution

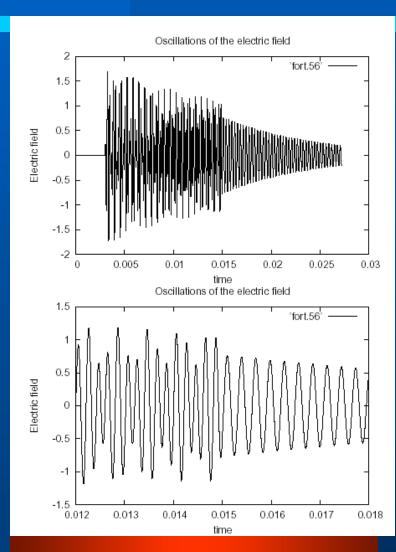
- Pseudo-spectral method
- Periodic boundary conditions
- Advance in time in fourier space (RK4)
- Non—linear part calculated in real space by using a FFT algorithm
- 256 grid points \rightarrow 128 modes \rightarrow 84 modes!

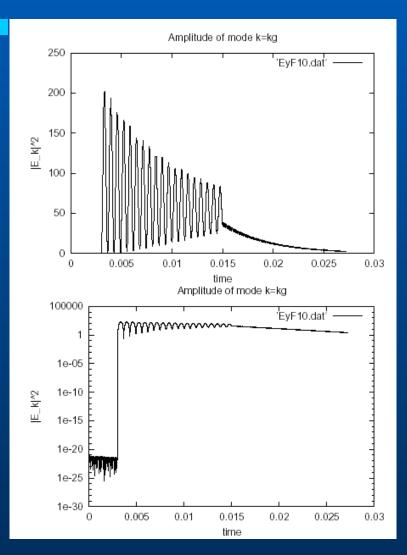
Typical values of the parameters

- k0=1.16 x 10**(-7) 1/cm
- ΔZ=5.4 x 10**7 cm
- T=1159420 K (=100 eV)
- ρ0=10**(-14) gr/cm**3
- Spitzer resistivity
- v=m uth $/\pi \lambda D^{**}2$
- kg=9k0

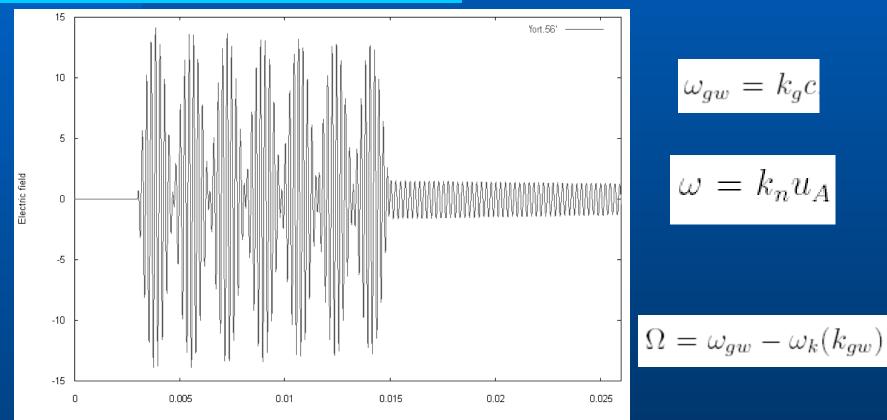
$$T_{gw}(k_g) = 0.0002 \text{ sec}$$
$$T(k_g) = 0.00029 \text{ sec}$$

Numerical Experiment



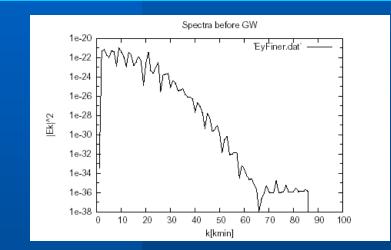


Numerical Results

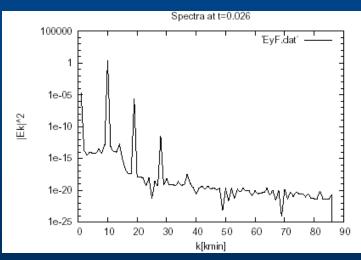


Time

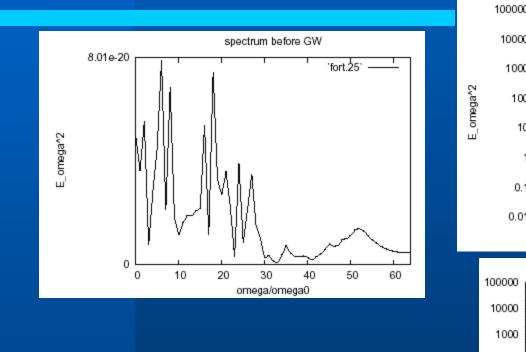
Spatial Spectrum

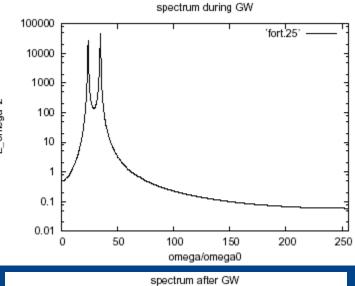


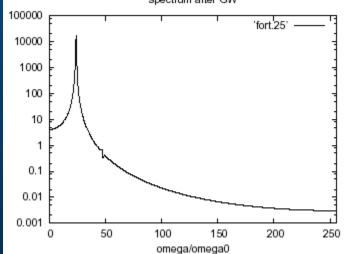
Spectra during GW 100 'EyFgrav.dat' 1 0.01 0.0001 1e-06 |Ek|^2 1e-08 1e-10 1e-12 1e-14 1e-16 1e-18 0 10 20 30 40 50 60 70 80 90 100 k[kmin]



Temporal Spectrum







Dominant harmonics during GW

• @ k=kg

 $ω = ω_{GW}$ (ALSO @ $k = ω_{GW}/U_a$) $ω = ω_{MS}$

• @ k=2kg

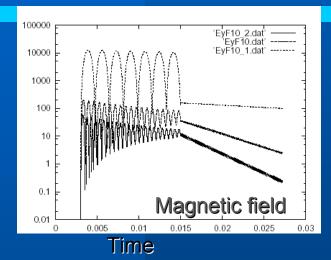
 $ω=2ω_{MS}$ $ω=2ω_{MS}+Ω$ $ω=2ω_{MS}+2Ω$

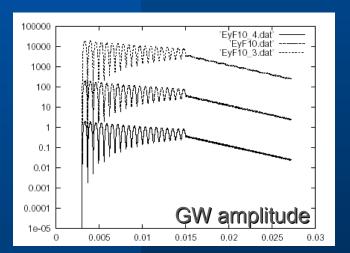
• @ k=3kg

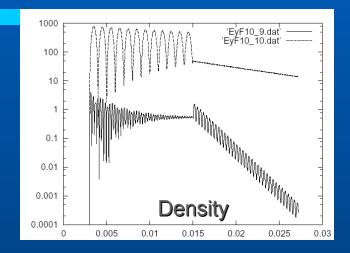
 $ω=3ω_{MS}$ $ω=3ω_{MS}+Ω$ $ω=3ω_{MS}+2Ω$

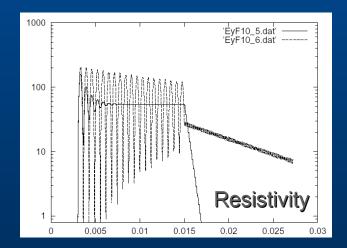
$$\Omega = \omega_{gw} - \omega_k(k_{gw})$$

Parametric Studies: electric field





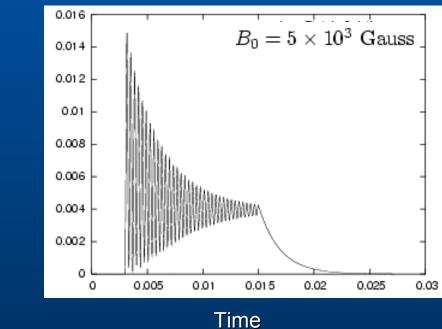


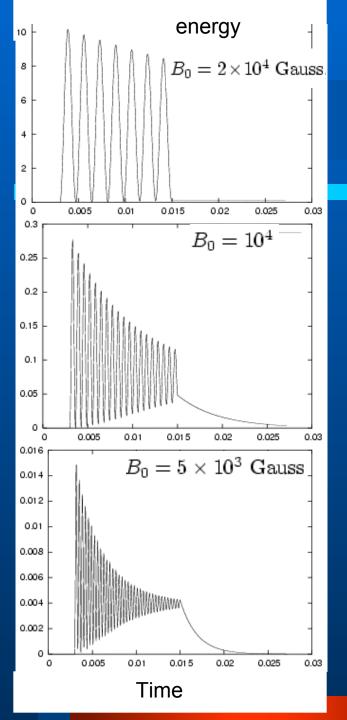


Energy deposited in the plasma

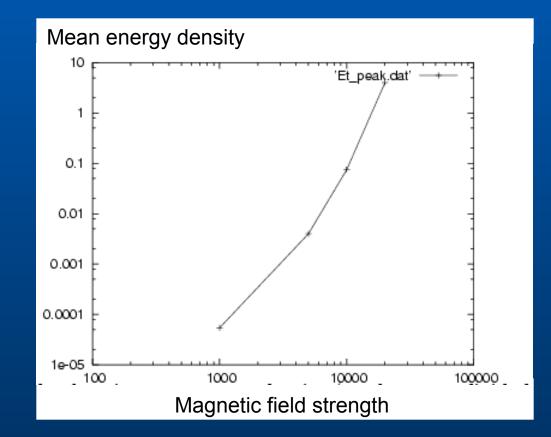
$$E_{total}(t) = \frac{1}{8\pi} \int E_y(z,t)^2 \, dz + \frac{1}{8\pi} \int (B_x(z,t)^2 - B_0^2) \, dz + \frac{1}{2} \int \rho(z,t) u_z(z,t)^2 \, dz$$

Energy density (erg/cm)

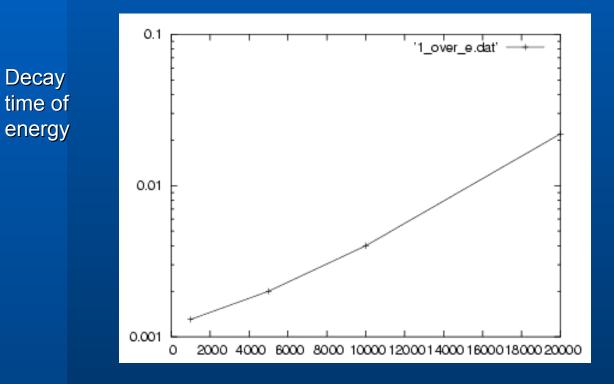




Plasma Energy







Magnetic field strength

Future studies

- Treat numerically the 2D problem
- Include both GW polarities
- Include higher order GW effects
- Include inhomogeneous plasma and magnetic field profile

THANK YOU