

# Simulation of a plasmoid penetrating a magnetic barrier

Herbert Gunell<sup>a</sup>, Tomas Hurtig<sup>b</sup>, Hans Nilsson<sup>c</sup>, Jeffrey Walker<sup>a</sup>,  
Mark Koepke<sup>a,d</sup>, and Nils Brenning<sup>d</sup>

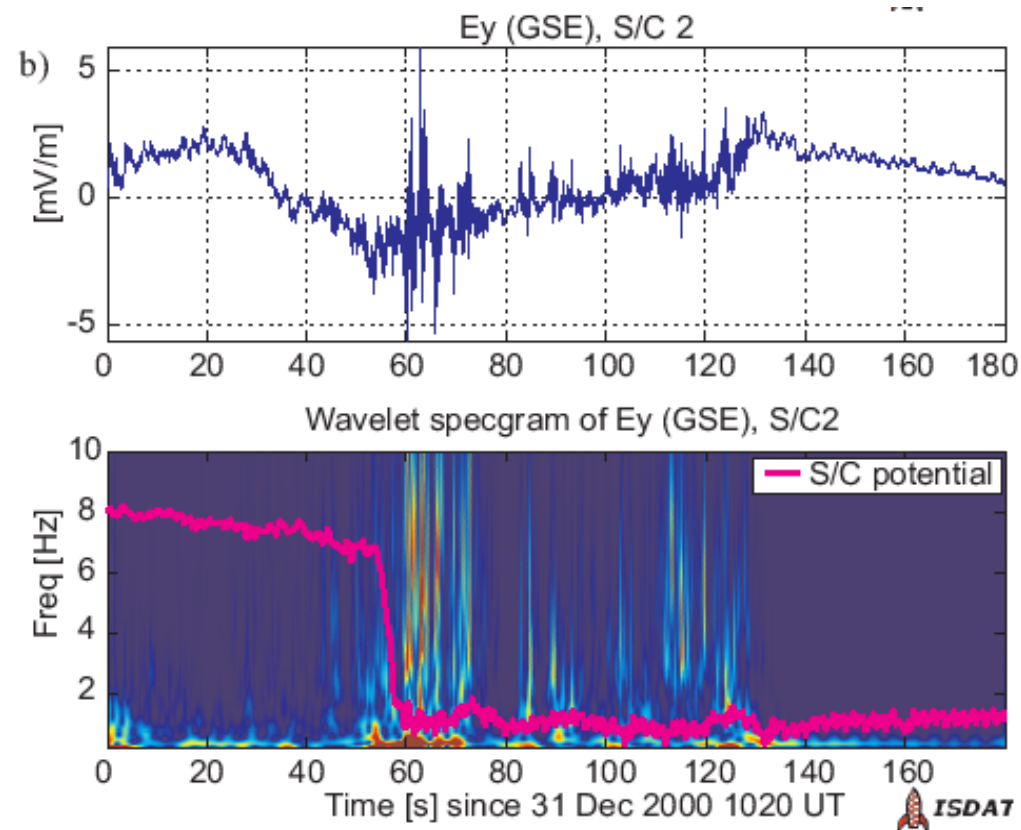
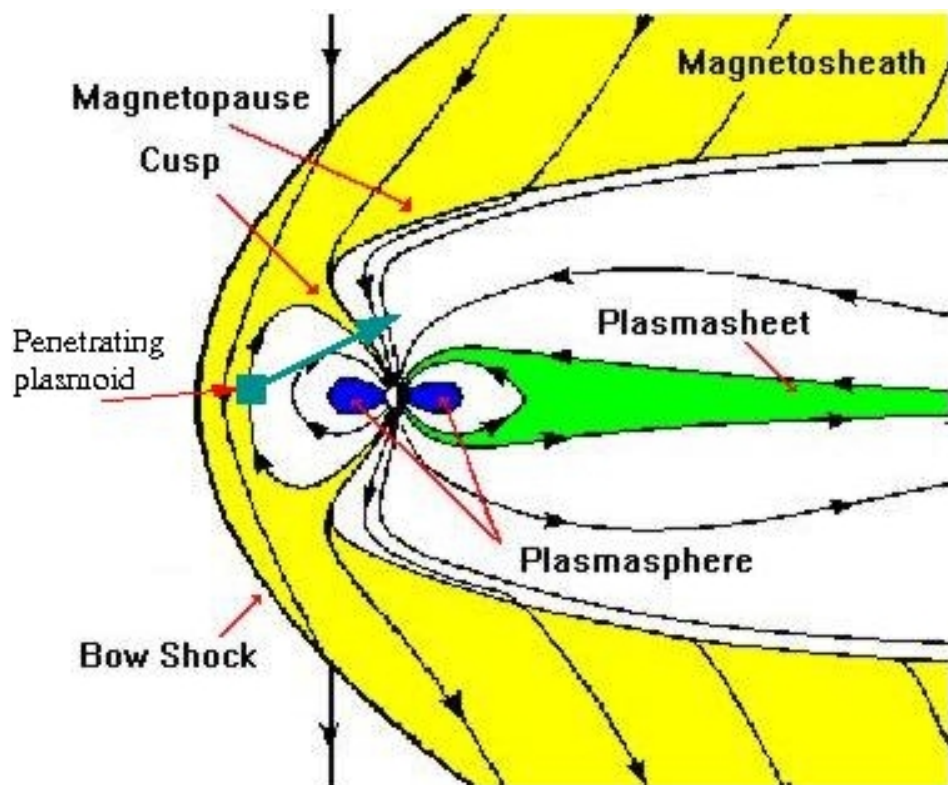
<sup>a</sup> West Virginia University, Morgantown, WV, USA

<sup>b</sup> Swedish Defence Research Agency, Stockholm, Sweden

<sup>c</sup> Swedish Institute of Space Physics, Kiruna, Sweden

<sup>d</sup> Royal Institute of Technology, Stockholm, Sweden

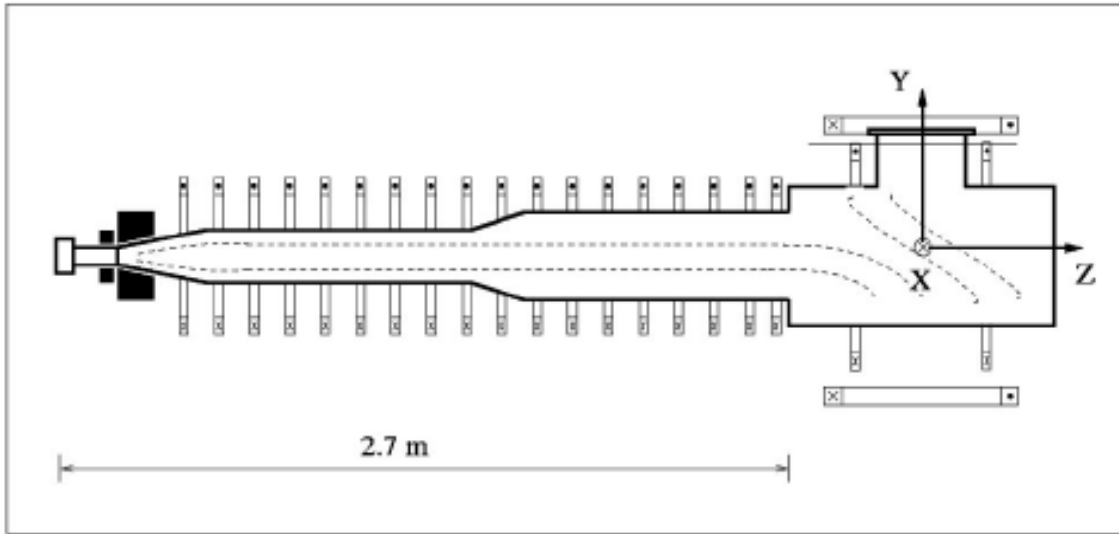
# Introduction – space



Left: Impulsive penetration at the magnetopause.

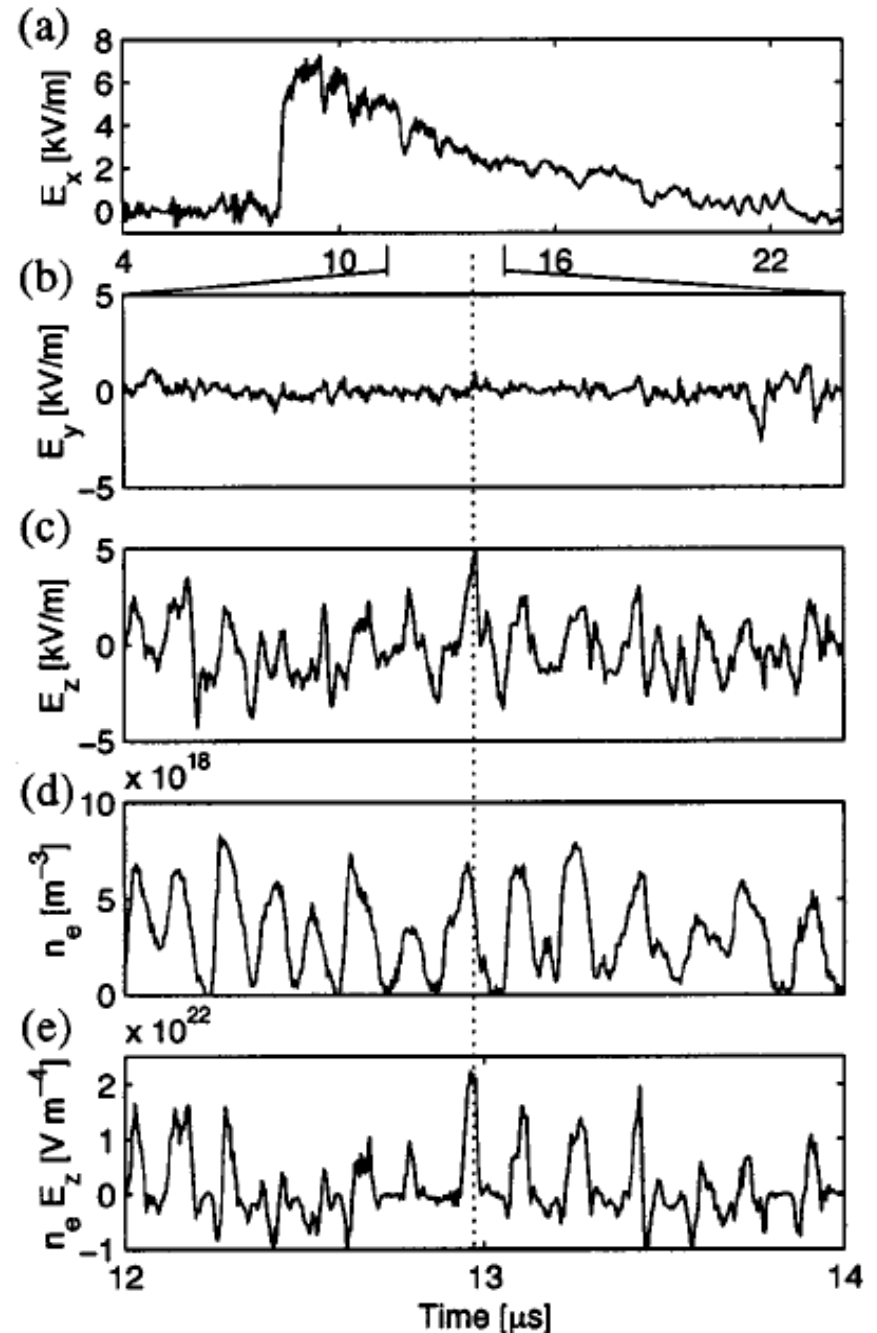
Right: Electric field measured by Cluster showing waves in the lower hybrid frequency range (*André, et al. 2001*).

# Introduction – lab

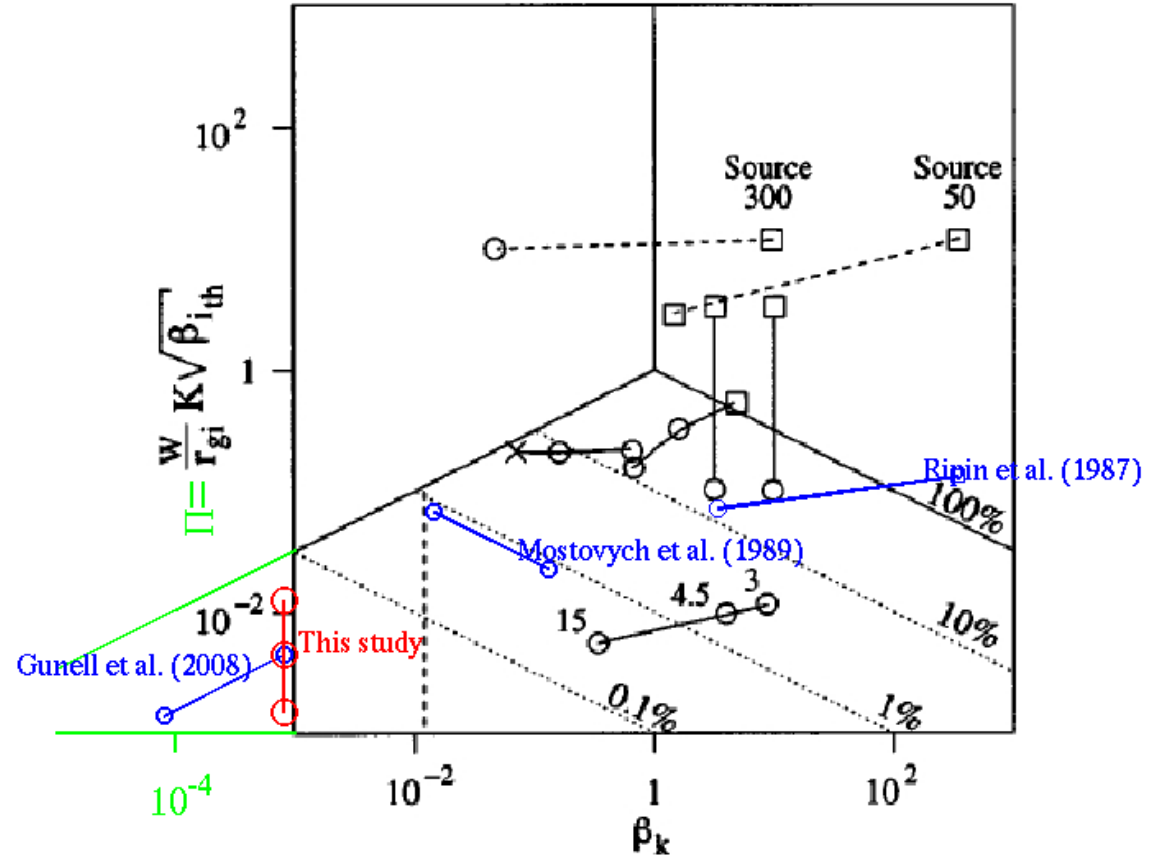
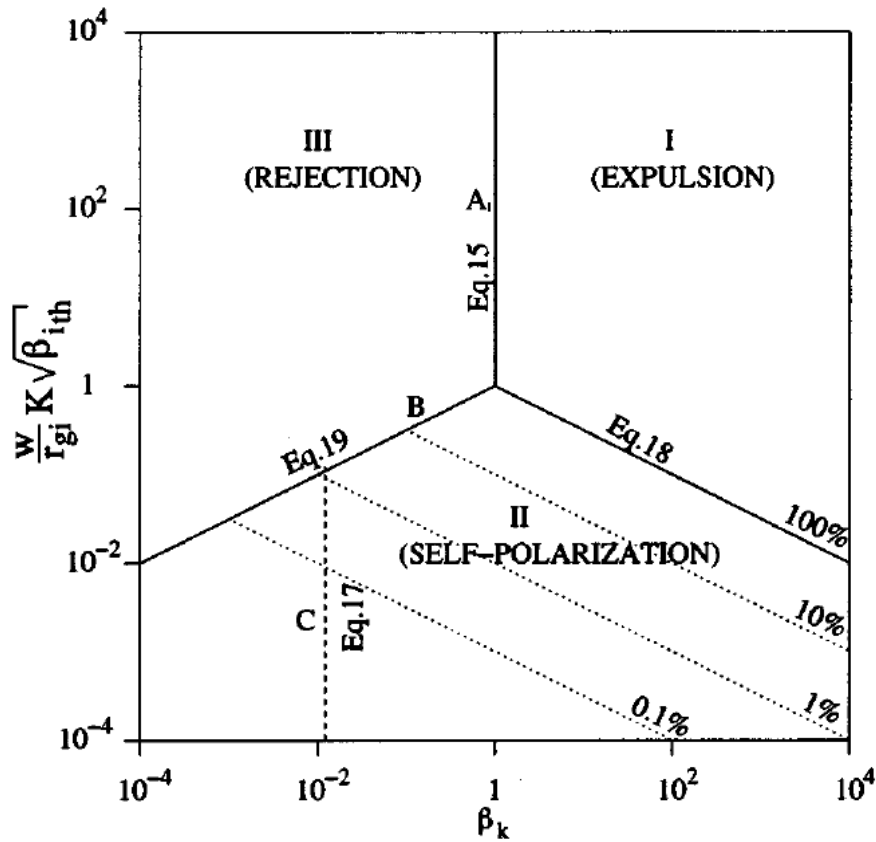


Above: A plasma cannon.

Right: a) Polarisation field  $E_x$ ; b) vertical electric field  $E_y$ ; c) E-field in flow direction  $E_z$ ; d) plasma density; e)  $n_e E_z$ .

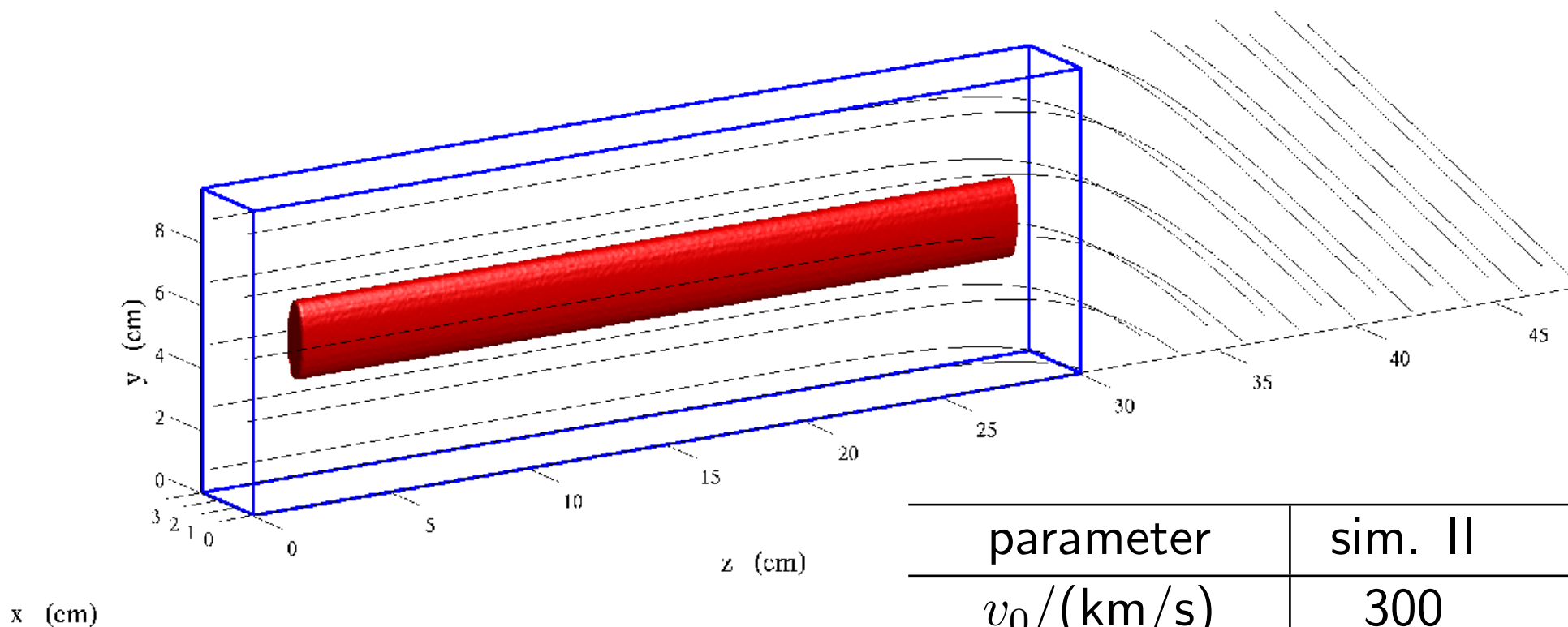


# Introduction – parameter regimes



1. **Expulsion.** A plasma structure can penetrate a magnetic barrier by expelling the magnetic field if  $\beta_k > 1$  and  $\Pi > 1/\sqrt{\beta_k}$ .
2. **Self-polarisation.** A plasma structure can penetrate a magnetic barrier by convection in a polarisation electric field if  $\Pi < \sqrt{\beta_k}$  for  $\beta_k < 1$  and  $\Pi < 1/\sqrt{\beta_k}$  for  $\beta_k > 1$ .
3. **Rejection.** The plasma cannot penetrate the magnetic barrier if  $\beta_k < 1$  and  $\Pi > \sqrt{\beta_k}$ .

# The long plasmoid



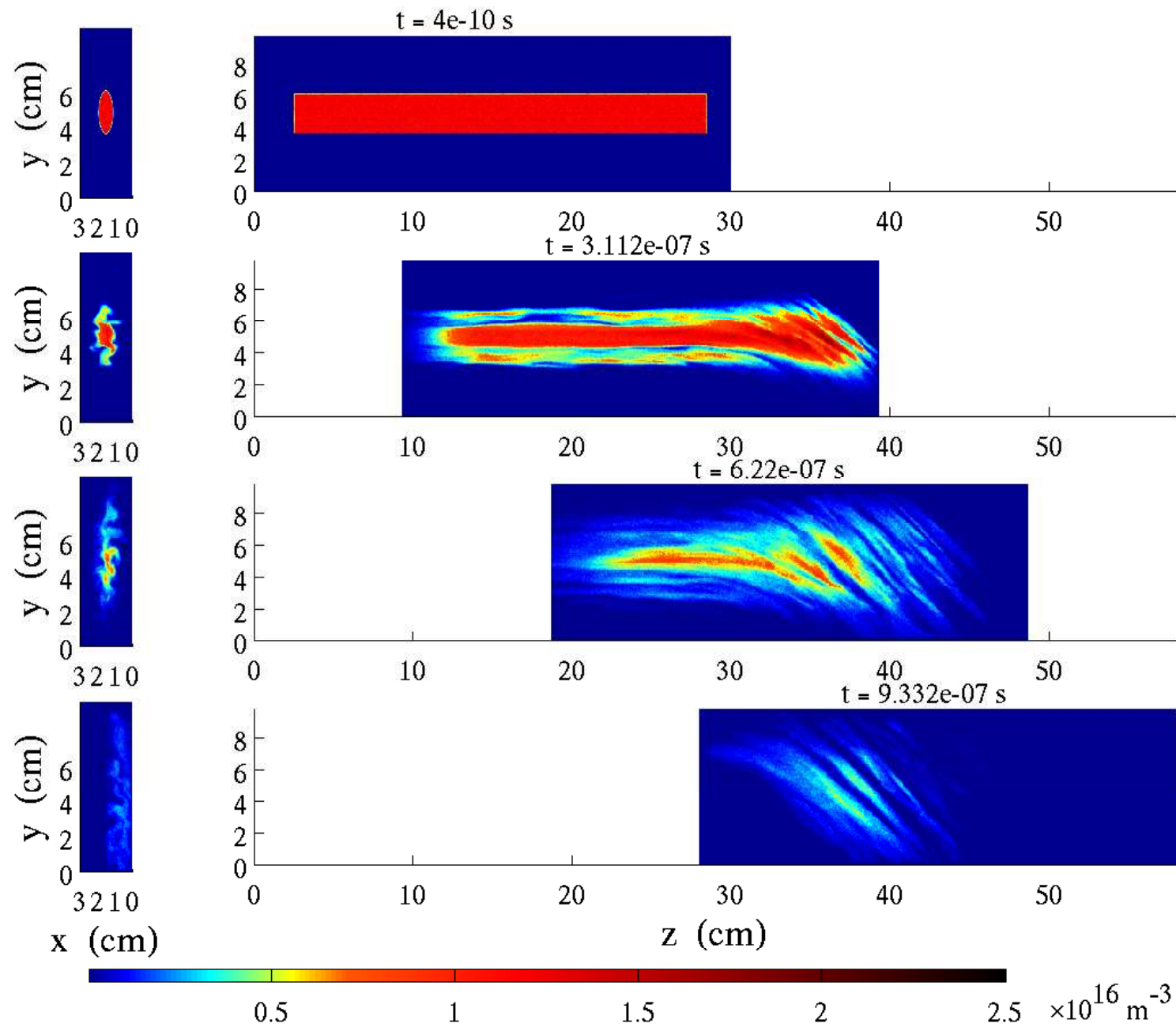
$$W_K = \frac{1}{2} n_0 m_i v_0^2$$

$$W_E = \frac{1}{2} \epsilon_0 (v_0 B_\perp)^2$$

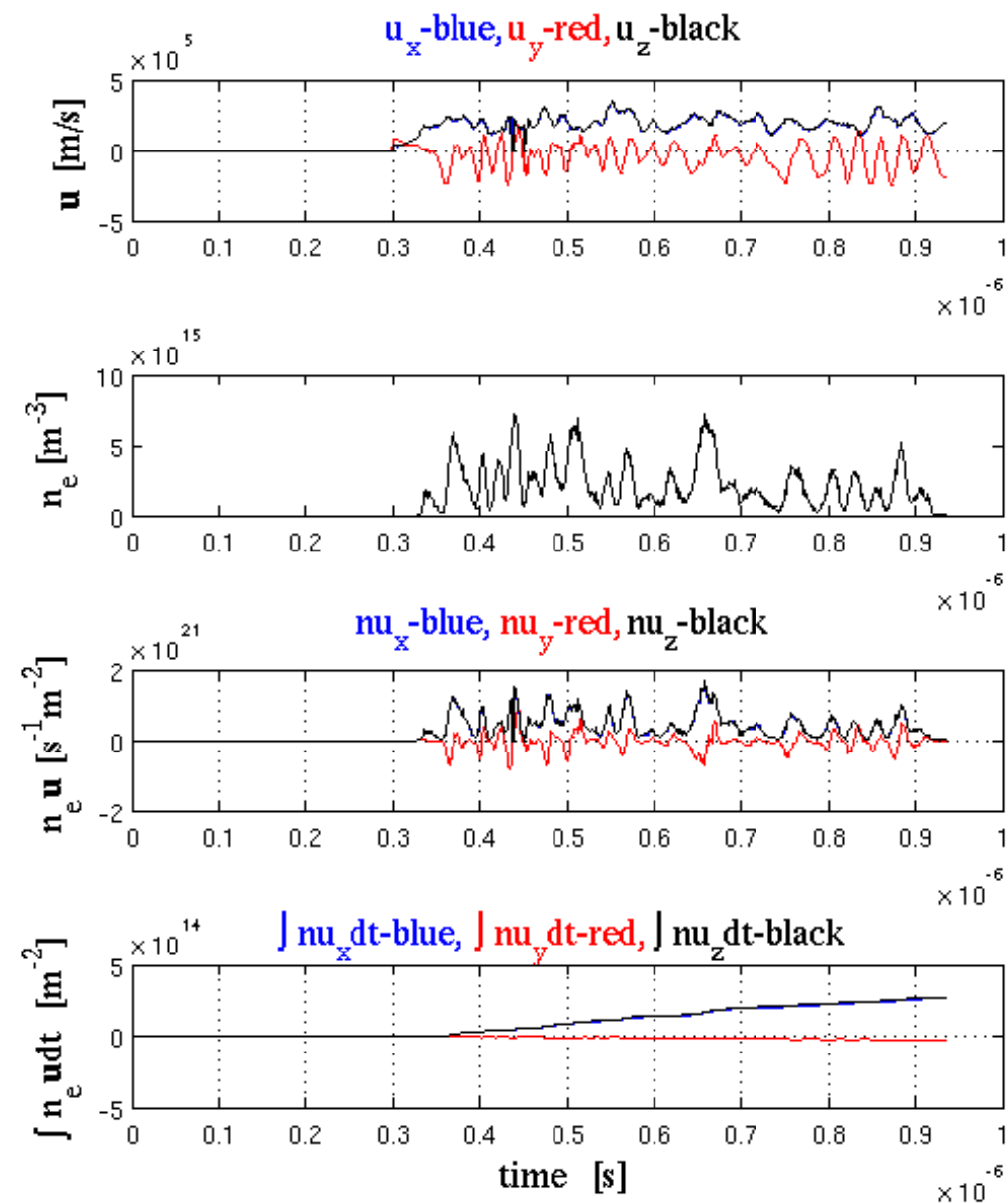
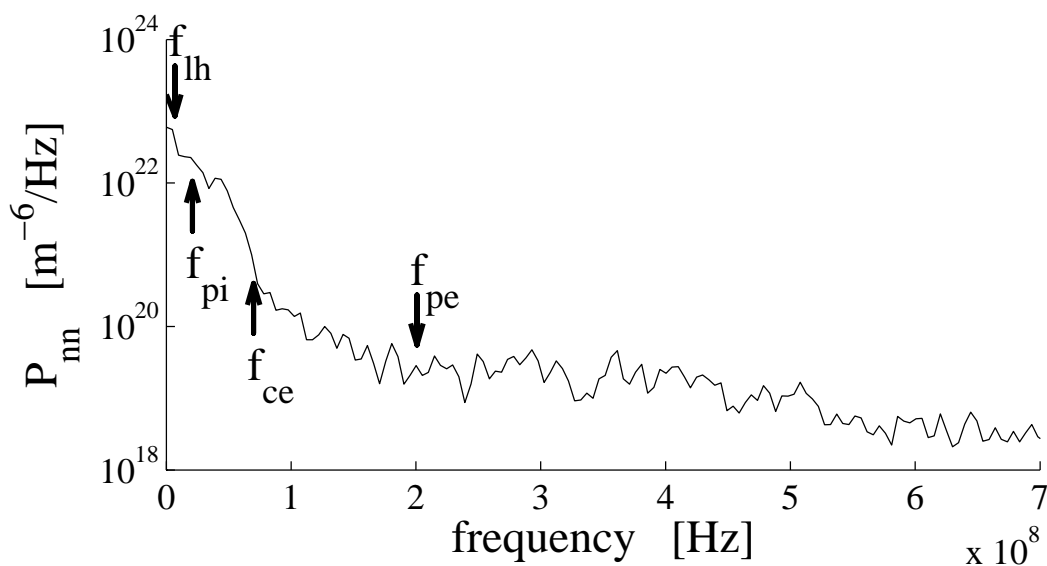
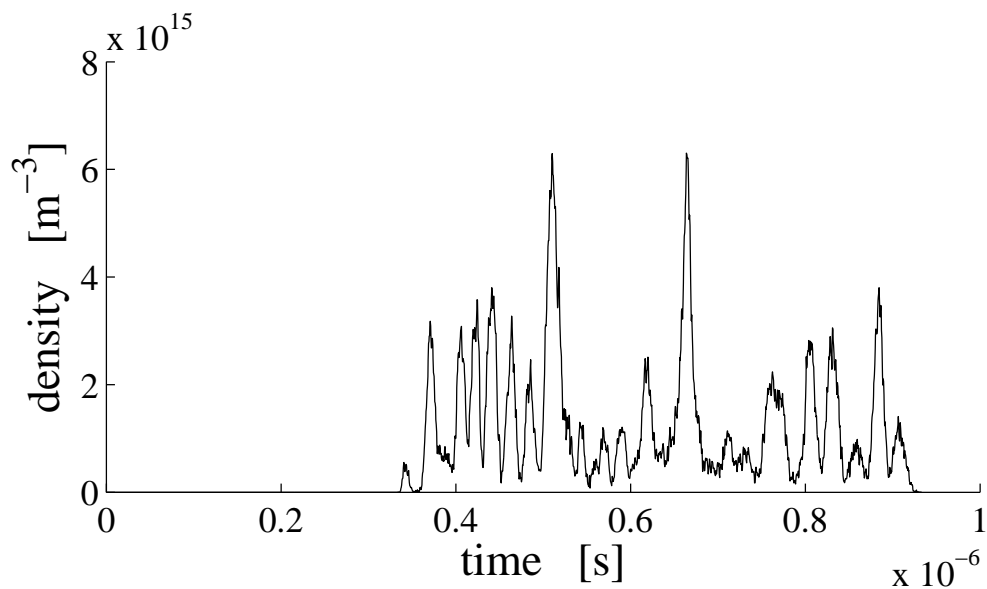
$$W_B = \frac{B_\perp^2}{2\mu_0}$$

parameter	sim. II	exp.
$v_0 / (\text{km/s})$	300	300
$B_\perp / \text{T}$	0.05	0.015
$n_0 / \text{m}^{-3}$	$10^{16}$	$10^{18}$
$m_i / m_e$	92	1836
$W_K / W_E$	756	$8 \cdot 10^5$
$\frac{1}{10} \frac{W_K}{W_E} \sqrt{\frac{m_e}{m_i}}$	7.9	$2 \cdot 10^3$
$\beta_k = W_K / W_B$	$8 \cdot 10^{-4}$	0.8

# The long plasmoid – density

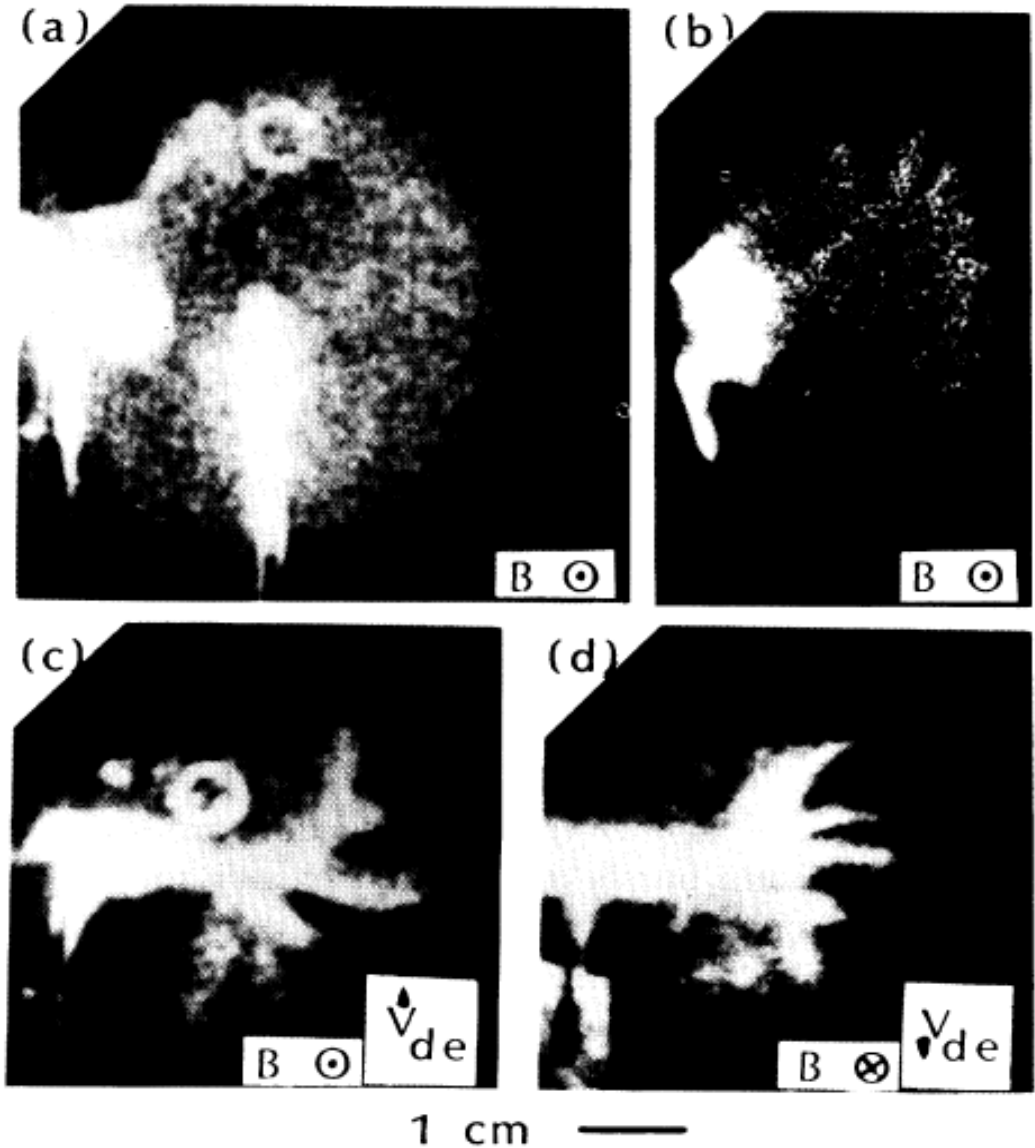
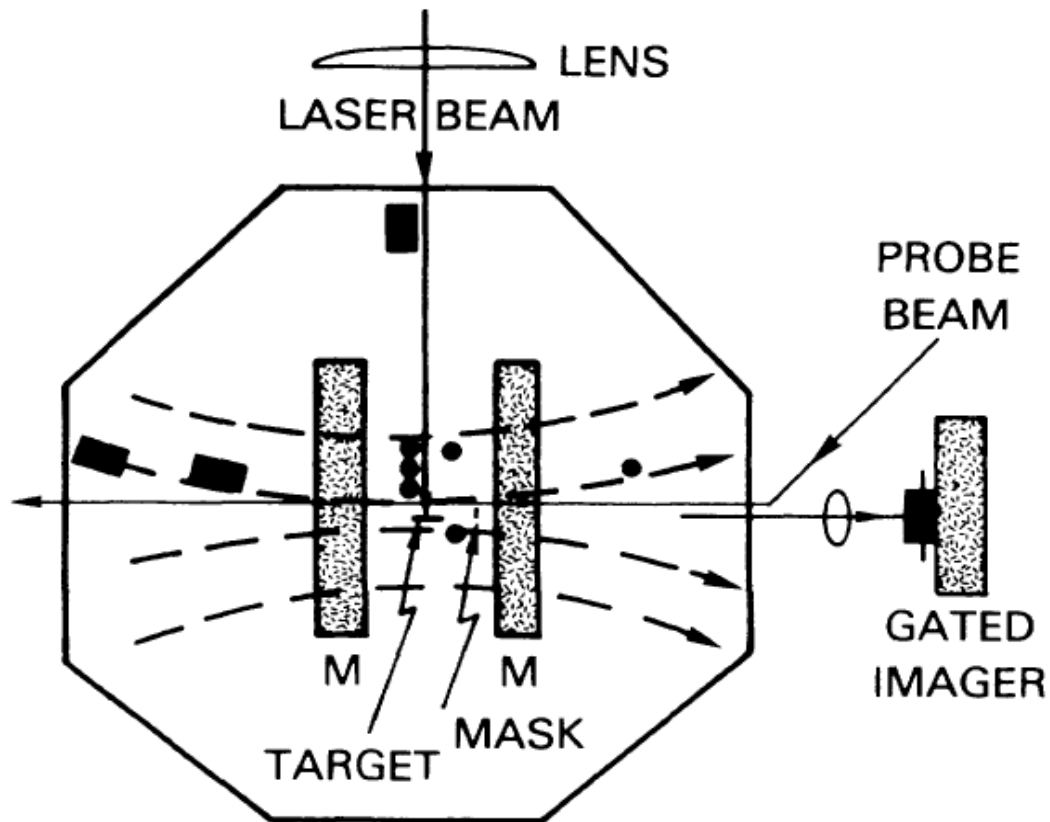


# The long plasmoid – waves



Density probe at  $z = 0.39$  m.

# Wide plasmoids



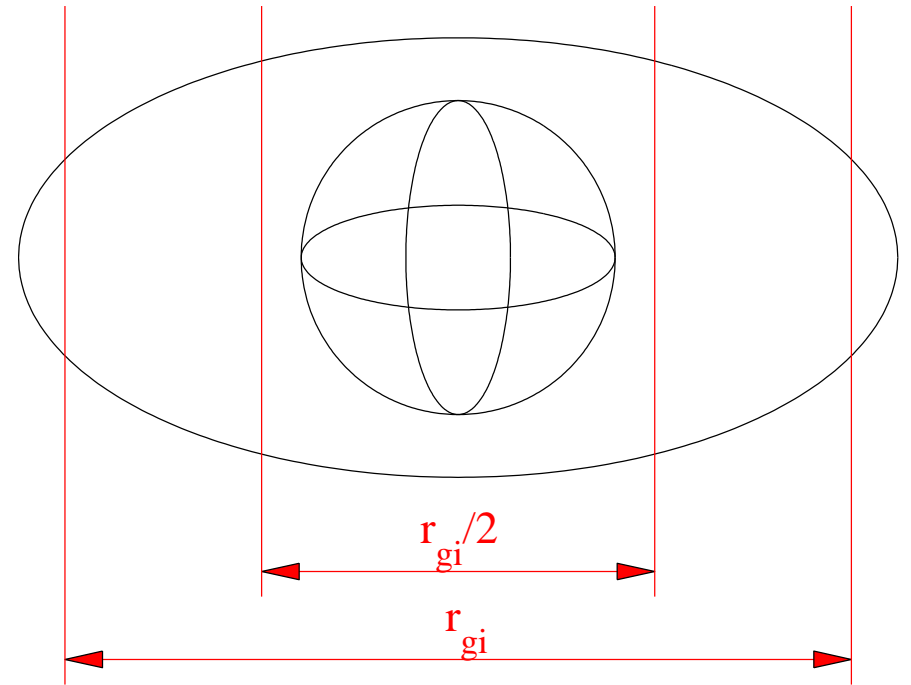
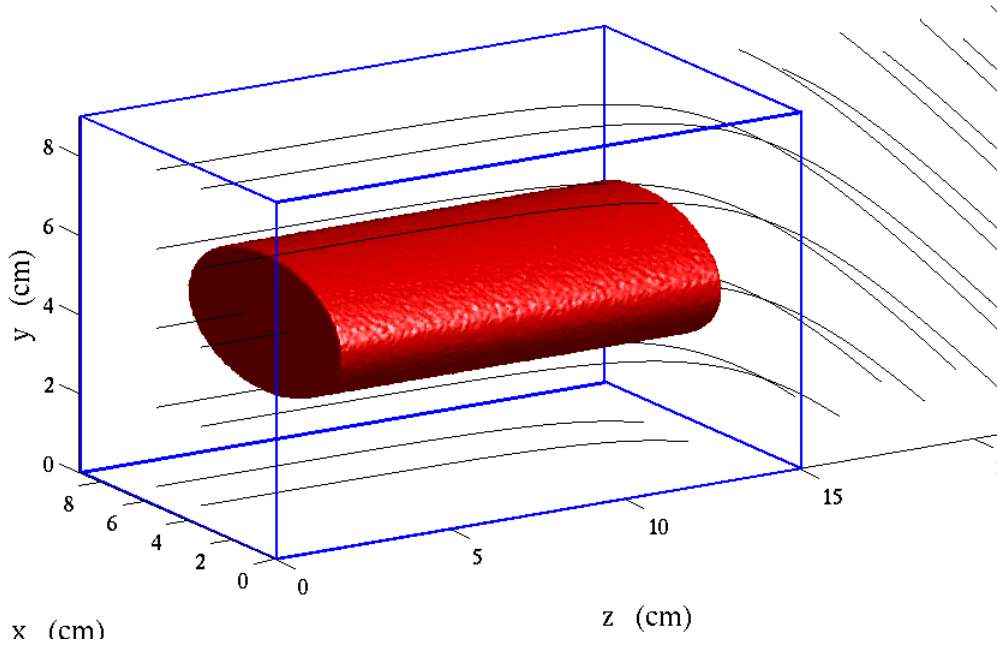
Pictures from *Ripin et al. (1987)*.



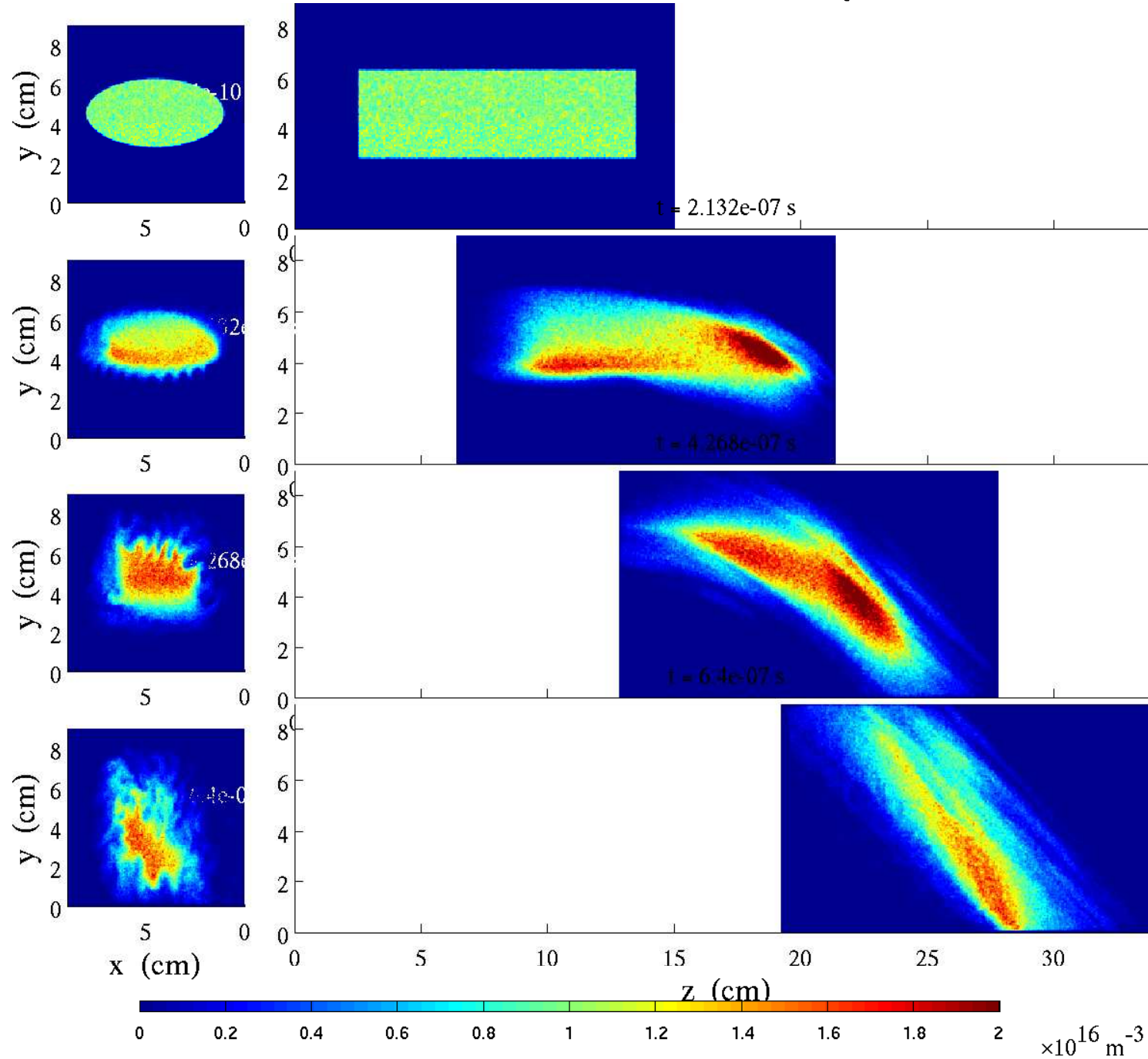
# The wide plasmoid

The limit estimated by *Lindberg (1978)*:

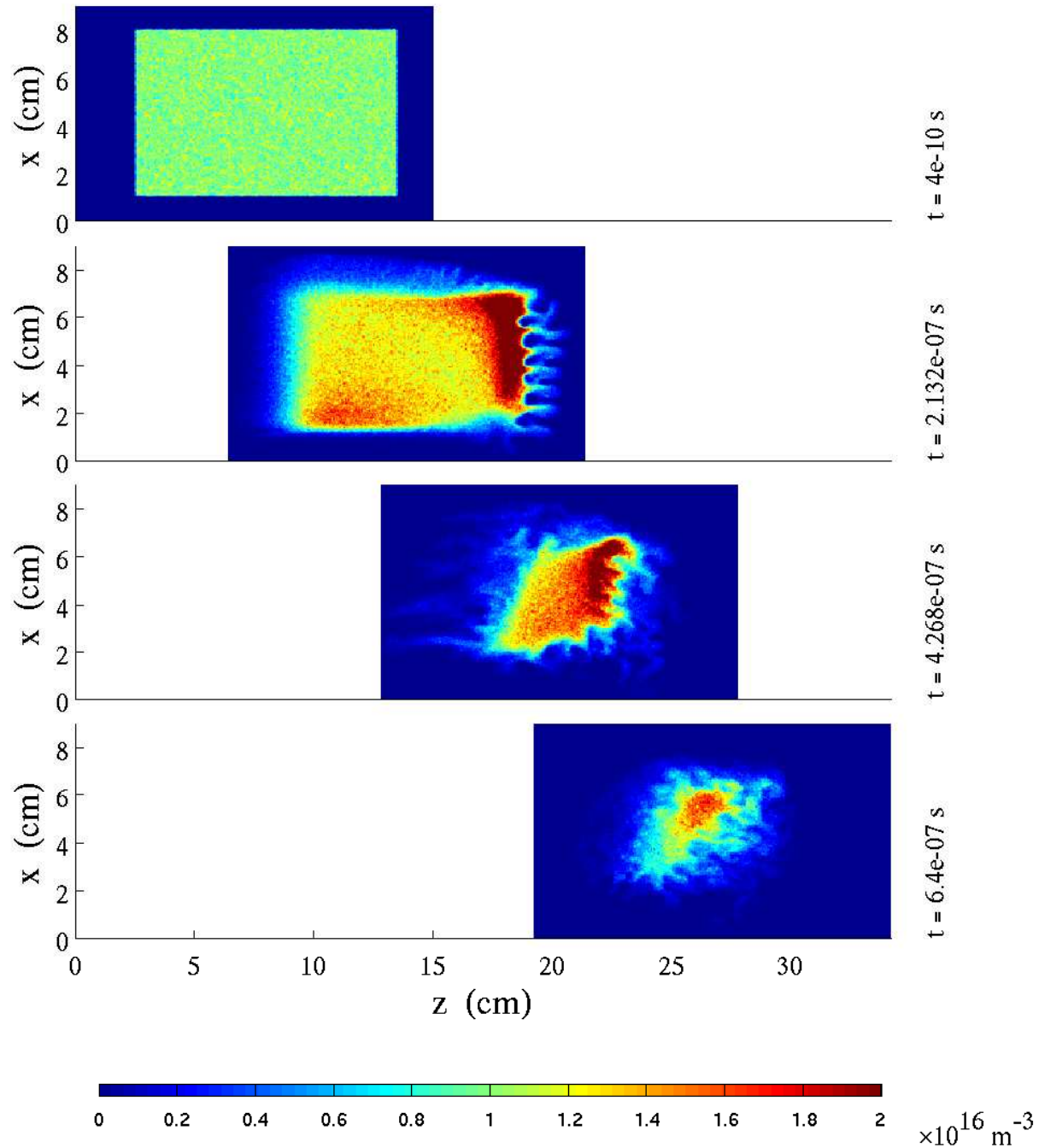
$$ewv_0B_{\perp} \leq \frac{1}{2}m_iv_0^2 \quad \Rightarrow \quad w \leq \frac{1m_iv_0}{2eB_{\perp}} = \frac{1}{2}r_{gi}$$



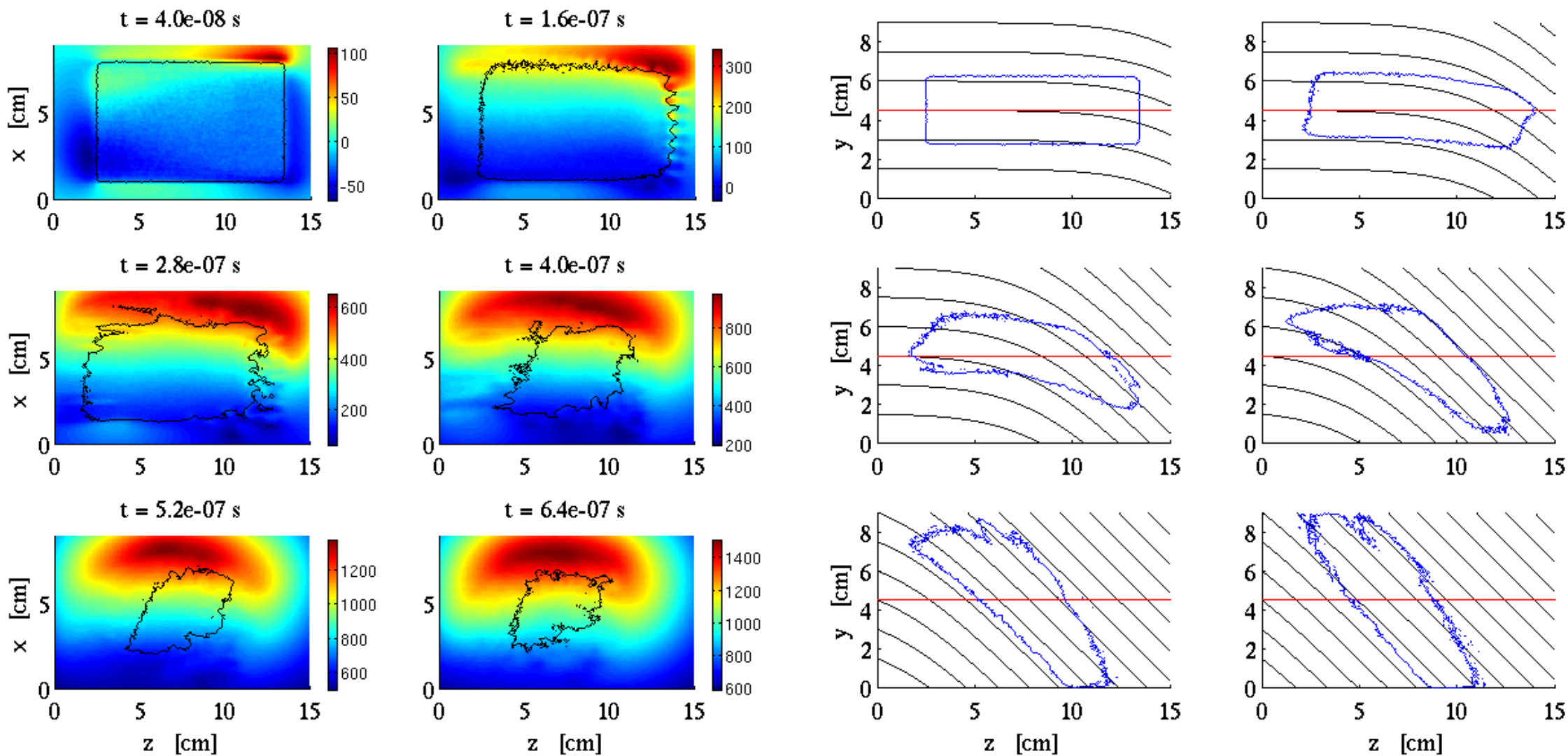
# The wide plasmoid – density (vertical slice)



# The wide plasmoid – density (horizontal slice)



# The wide plasmoid – penetration



# Conclusions

- Lower hybrid frequency waves are seen at the magnetopause, in laboratory experiments and in simulations.
- For  $w < r_{gi}/2$  all cross sectional shapes are compressed into vertically aligned structures.
- Plasmoids with  $w \approx r_{gi}$  can penetrate with the aid of a backward propagating potential and through compression to smaller widths.
- The finger-like structures that develop at an early stage for  $w \approx r_{gi}$  plasmoids are a result of the same instability that gives rise to the lower hybrid frequency waves along the flanks.
- Next step:  $w \gg r_{gi}$