

Basic Processes of Turbulent Plasmas

Summer School

23-28 September 2003, Chalkidiki, Greece

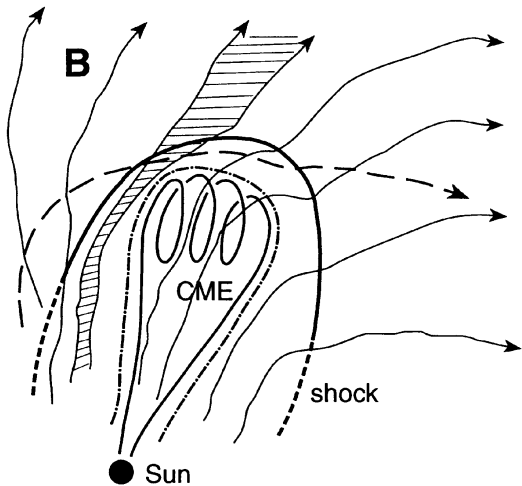
**Plasma-turbulence and wave-particle interaction
downstream of the main interplanetary shock of the
Bastille Day coronal mass ejections**

Reinald Kallenbach

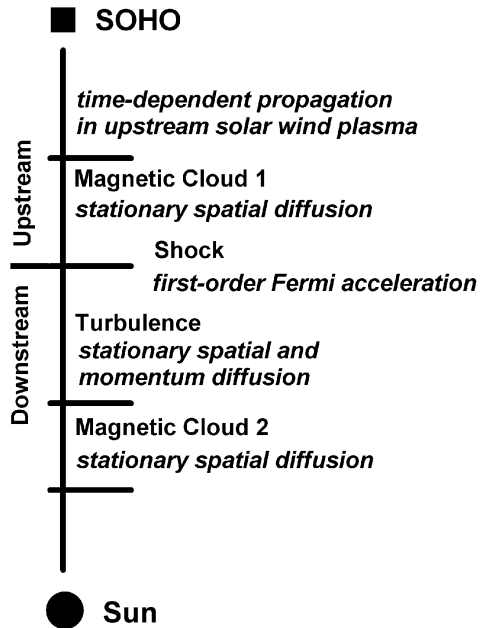
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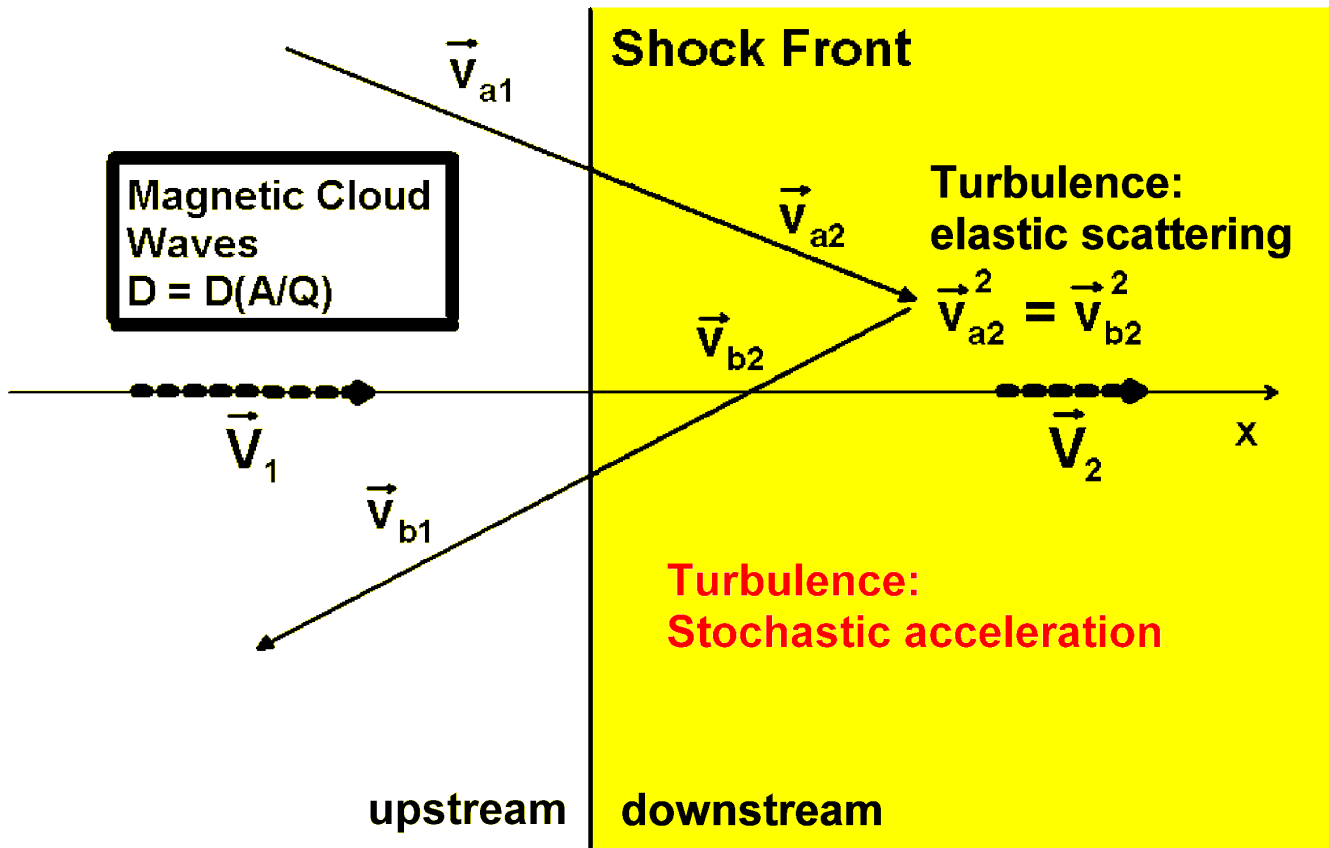
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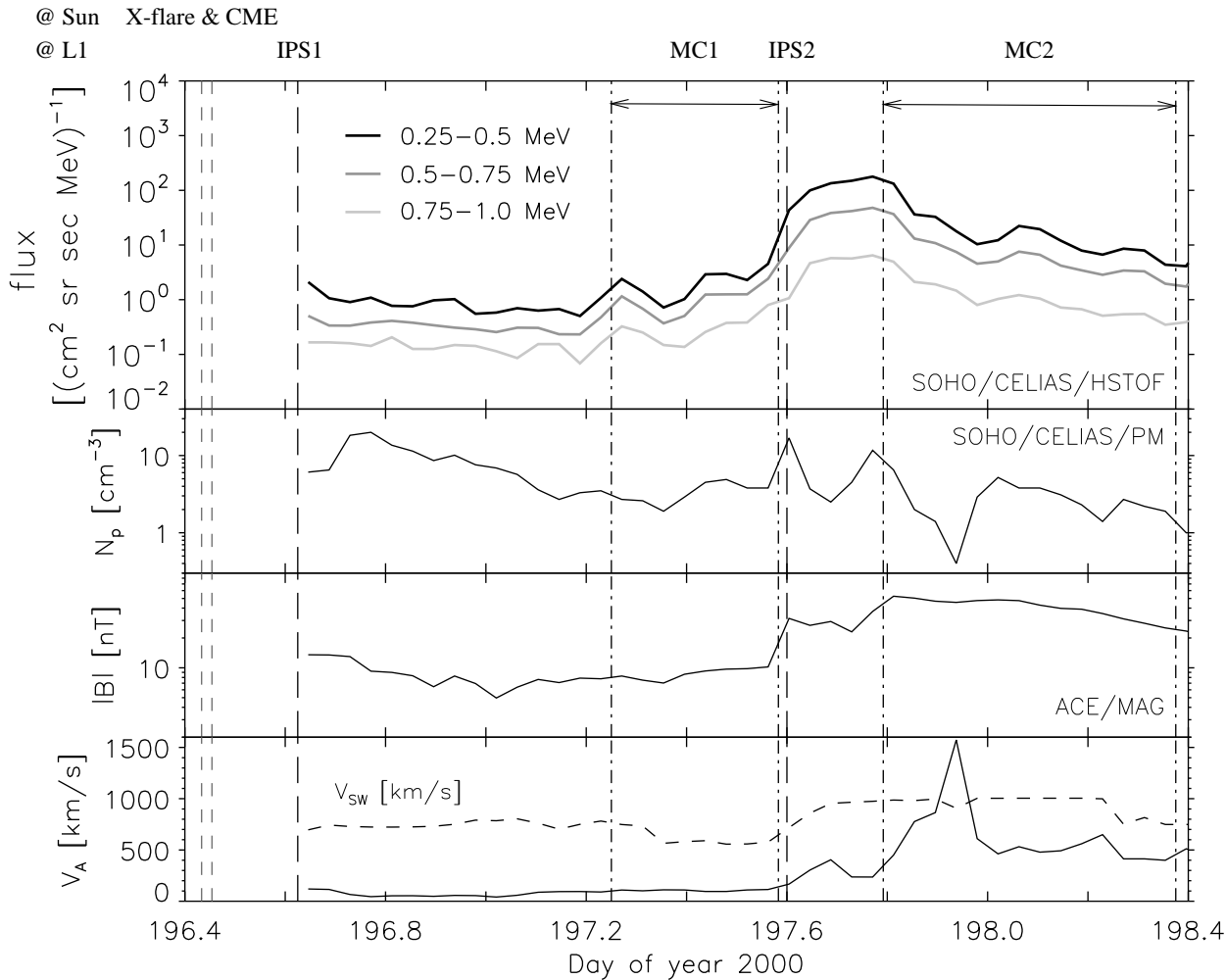


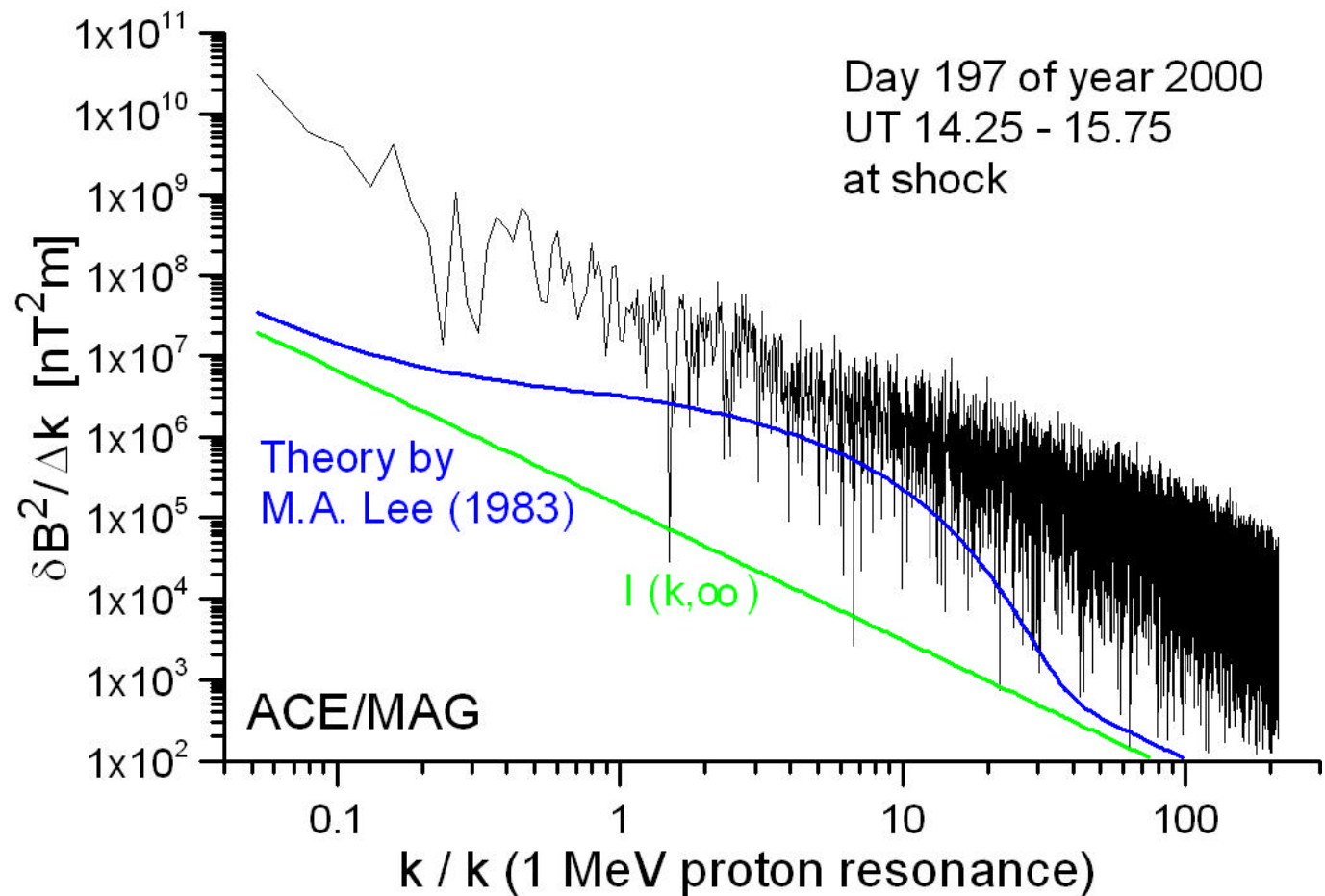
Configuration of a CME
(M.A. Lee, 1997)

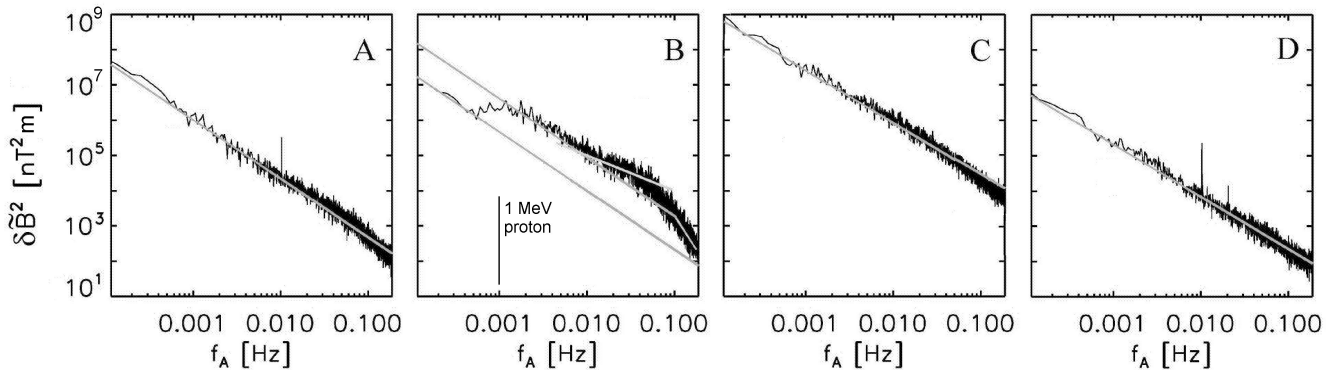
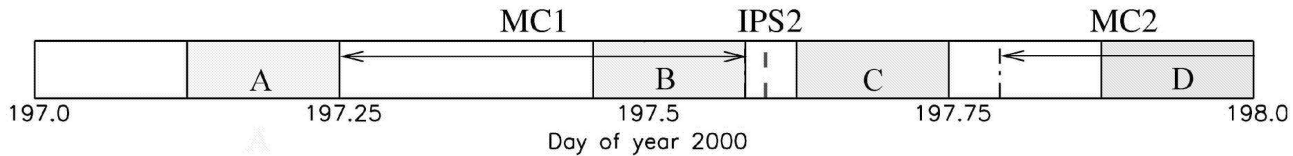




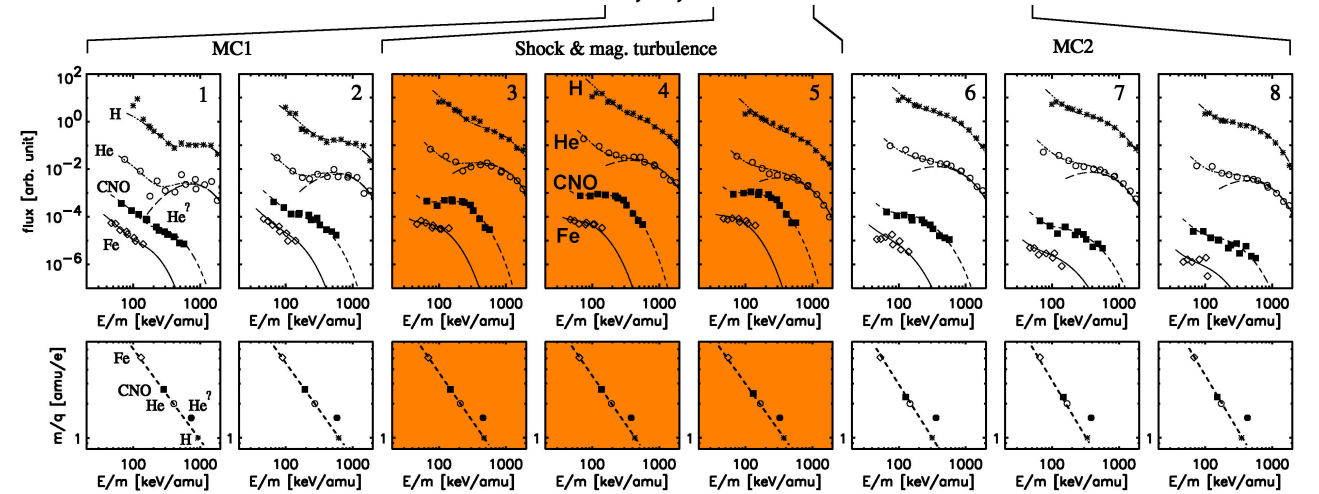
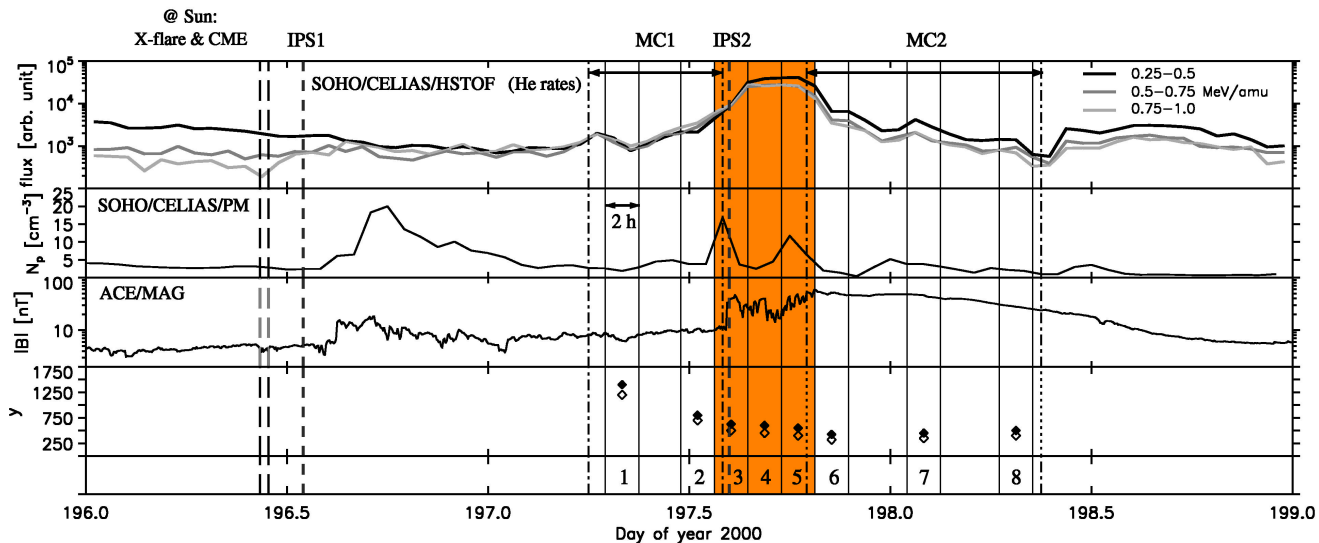
The Bastille Day Event (July 14-16, 2000)







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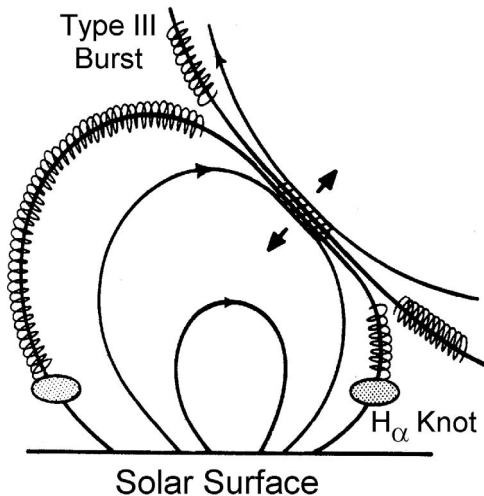


Figure 136. Erzeugung von Elektronenstrahlen oder Asymmetrien in der Elektronentemperatur ($T_{e\parallel} \gg T_{e\perp}$) durch die Rekonnexion in impulsiven Flares auf der Sonnenoberfläche.

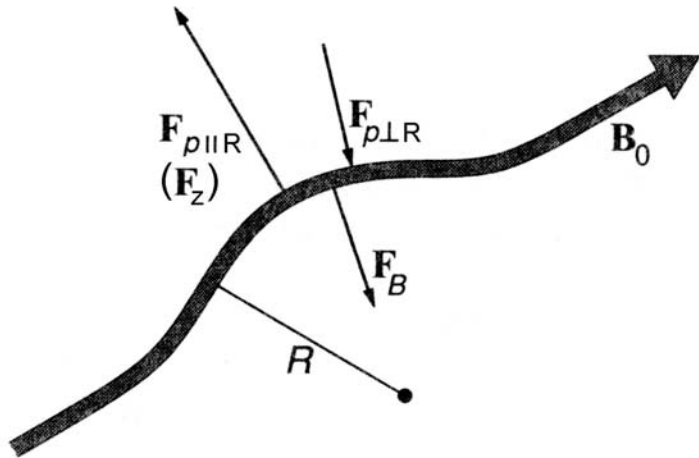
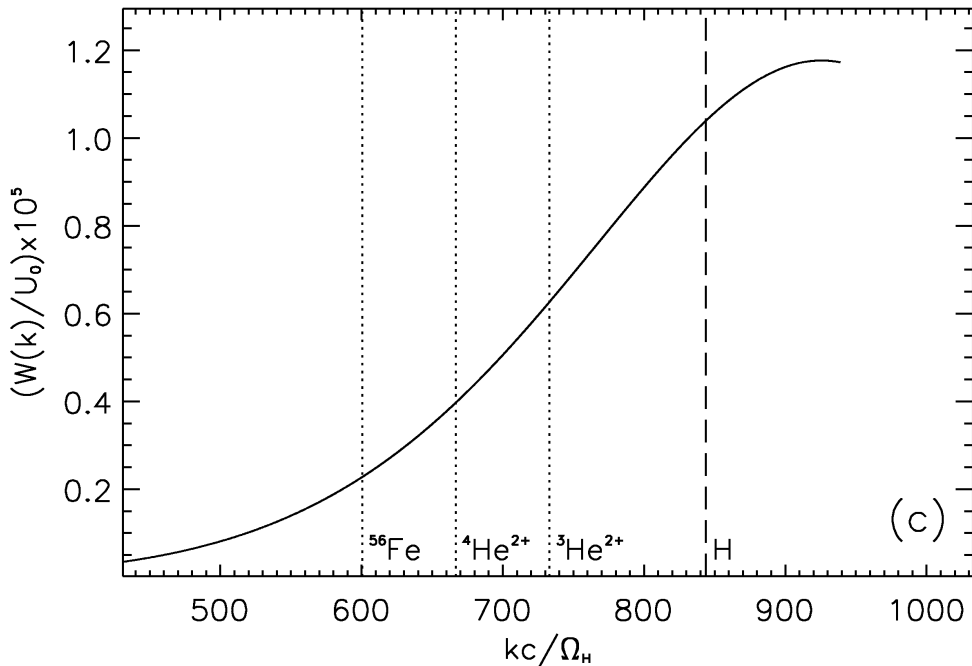


Figure 135. Schematische Darstellung der Feuerwehrschlauchinstabilität (Treumann and Baumjohann, 1997; Seite 51).

Electron Firehose Instability in impulsive solar flares

$$T_{e,\parallel} / T_{e,\perp} = 10, \quad E_{\text{wave}} = 0.01 B_0^2 / (2\mu_0)$$

Paesold et al., 2003



cyclotron resonance

$$\xi = \Omega_s + k_{\parallel} v_{\parallel} - \omega$$

$$v_{\perp} \dot{v}_{\perp} \approx \frac{2\pi q^2 E_{\text{rms}}^2}{m^2 \delta\omega}$$

$$\frac{\dot{W}_{\perp}}{m} \approx 6.1 \left[\frac{\text{Mev}}{\text{amu s}} \right] \left(\frac{E_{\text{rms}}}{100 \text{ V m}^{-1}} \right)^2 \left(\frac{10^6 \text{ rad s}^{-1}}{\delta\omega} \right) \left(\frac{Q}{A} \right)^2$$

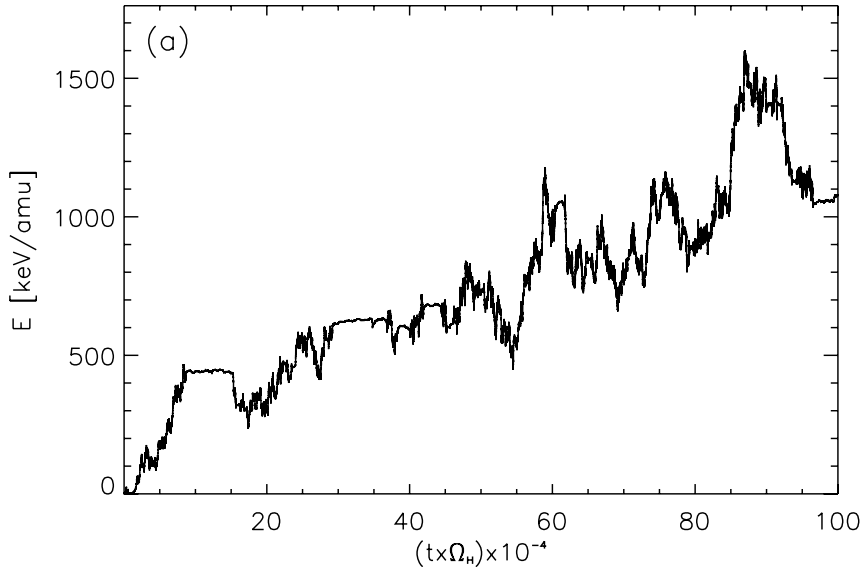
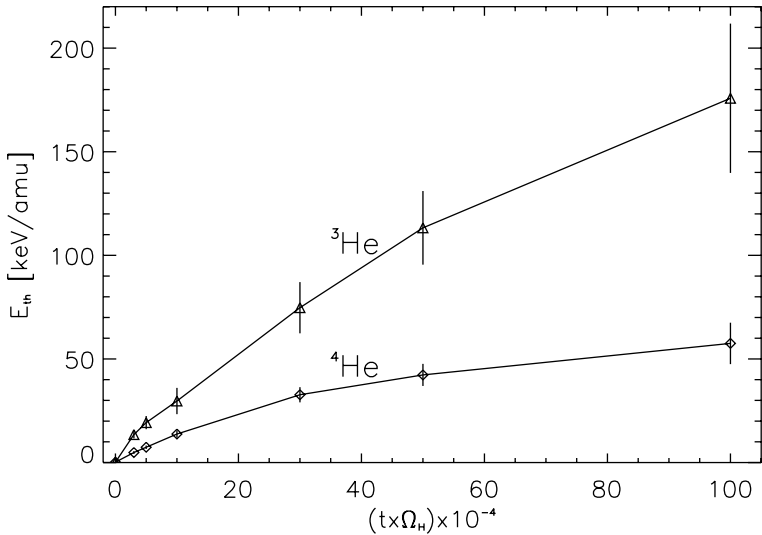
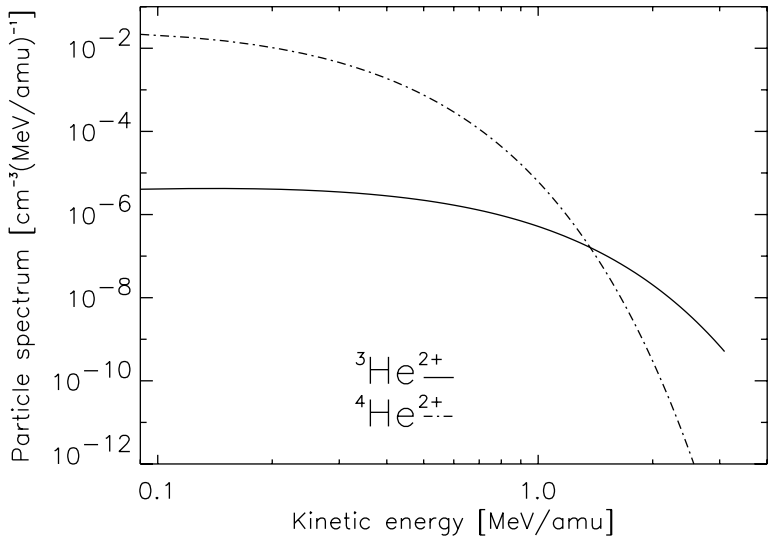
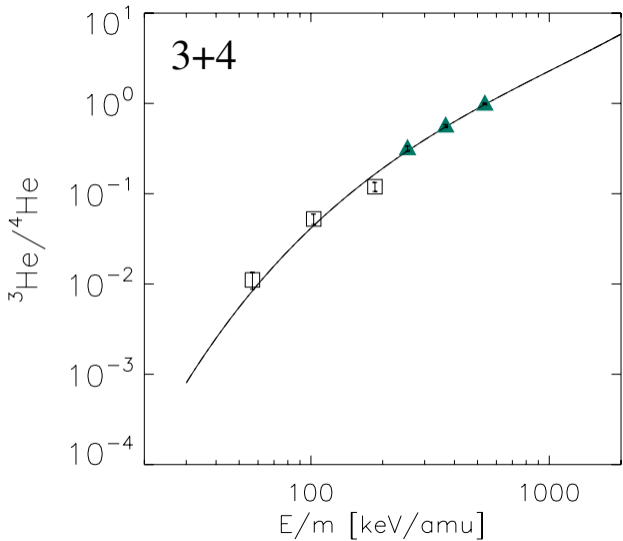


Figure 141. Zunahme der Gesamtenergie eines Ions mit der Zeit (von Paesold *et al.*, 2003).







Transport equation:

$$\frac{\partial f}{\partial t} + \mathbf{V} \cdot \nabla f = \nabla \cdot (\kappa \nabla f) + I - S + \frac{p}{3} \frac{\partial f}{\partial p} \nabla \cdot \mathbf{V} + \frac{1}{p^2} \frac{\partial}{\partial p} \left(p^2 D_{pp} \frac{\partial f}{\partial p} \right)$$

Roll-over of velocity distribution:

$$f(v) \propto \exp \left[-v^2/v_{\text{ro}}^2 \right] ,$$

$$v_{\text{ro}}^2 \approx v_{1\text{MeV}}^2 L_{\text{turb}} k_{1\text{MeV}}^2 (2 - \alpha) \left(\frac{Q}{A} \right)^{2-\alpha} \frac{\delta \tilde{B}^2(k_{1\text{MeV}})}{B_0^2} \approx v_{1\text{MeV}}^2 0.01 \left(\frac{Q}{A} \right)^{2-\alpha} (2 - \alpha)$$

$$\text{for } L \approx 1.4 \times 10^{10} \text{ m} ; k_{1\text{MeV}} \approx 1.4 \times 10^{-8} \text{ m}^{-1} ; \frac{\delta \tilde{B}^2(k_{1\text{MeV}})}{B_0^2} \approx 3 \times 10^5 \text{ m}$$

$-\alpha < 2$: spectral index power spectral density $\delta \tilde{B}^2(k)$

Momentum diffusion:

Quasi-linear theory, “1D” ;

2D-MHD (le Roux et al., 2002):

$$D_{vv;1D} = \pi \Omega_s^2 v |\mu| \delta \tilde{B}^2(k_{\text{res}}) / B_0^2 ;$$

$$D_{vv;2D} = \eta 2\pi^3 \Omega_s \mu^2 v^2 V_A^2 \langle \delta \tilde{B}_\perp^2 \rangle^2 / B_0^4 / V^2$$

μ : cosine of pitch-angle ; $\eta = r_A - \frac{1}{4} \sigma_c^2 (1 + r_A)^2$ σ_c : cross-helicity ;

r_A : ratio spectral energy density of velocity fluctuations to magnetic field fluctuations

$$\frac{D_{vv;2D}}{D_{vv;1D}} \approx \eta (L_{\text{turb}} k)^{\alpha-1} = \eta (L_{\text{turb}} k_{1\text{MeV}})^{\alpha-1} \left(\frac{k}{k_{1\text{MeV}}} \right)^{\alpha-1}$$

$$\approx 40\eta \left(\frac{k}{k_{1\text{MeV}}} \right)^{2/3} ; L_{\text{turb}} \approx 1.4 \times 10^{10} \text{ m} ; ; k_{1\text{MeV}} \approx 1.4 \times 10^{-8} \text{ m}^{-1} ; \alpha = 5/3$$

Conclusions

- Second-order Fermi acceleration for super-Alfvénic particles in the heliosphere is important, if data of Bastille Day event are due to stochastic acceleration in downstream 2D-MHD turbulence.
- Energetic (suprathermal) particles are probes for plasma turbulence.
- Need a theory for sub-Alfvénic particles.