



CME DYNAMICS AND PHYSICAL CONNECTION BETWEEN CMEs AND FLARES

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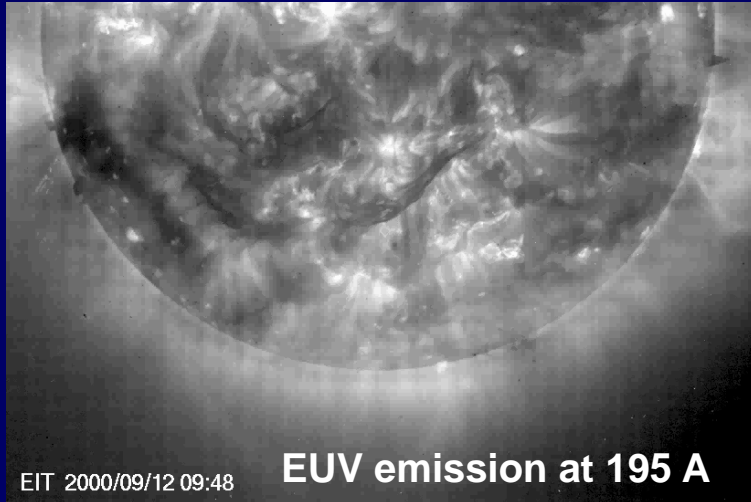
George Mason University

Modern Challenges of Nonlinear Plasma Physics

15–19 June 2009 Sani Resort, Halkidiki, Greece

SOLAR ERUPTIONS

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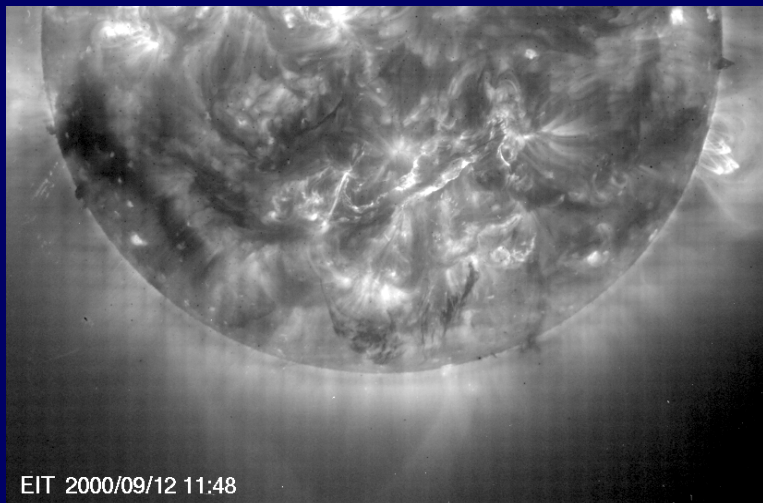


- **Solar eruptions:** Coronal mass ejections (CMEs), flares, prominence eruptions
- Canonical parameters of solar eruptions:
 - KE, photons, particles $\sim 10^{32-33}$ erg
 - Mass $\sim 10^{14-16}$ g
 - Speed $\sim 100 - 2000$ km/s
- *Space Weather.* CMEs are the solar drivers of large *geomagnetic storms*



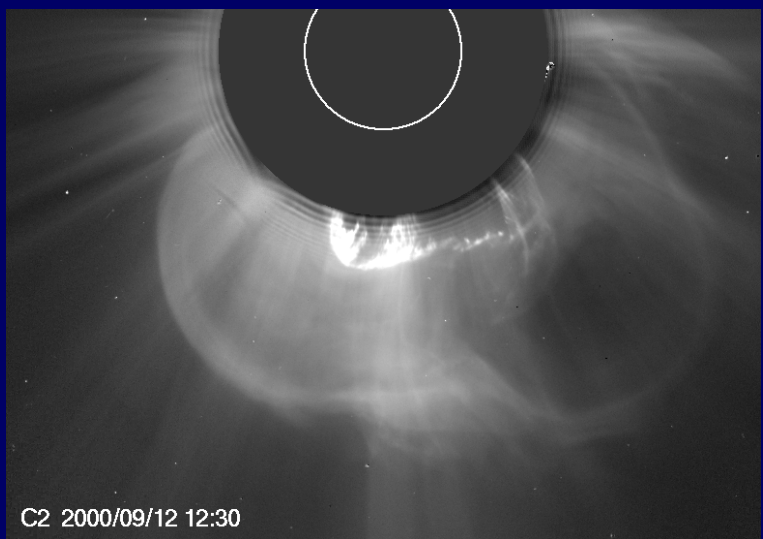
SCIENTIFIC CHALLENGES

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Observational challenges:

- All remote sensing
- Different techniques observe different aspects/parts of an erupting structure
- 3-D geometry not directly observed



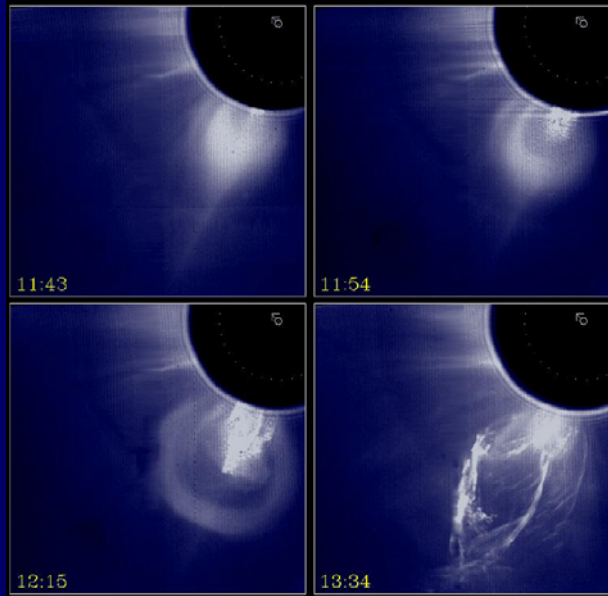
Theoretical challenges:

- An important unsolved question of theoretical physics
- Energy source
- Driving force (“magnetic forces”)
 - Underlying magnetic structure
- Physical relationship between CMEs, flares, and eruptive prominences (EPs)

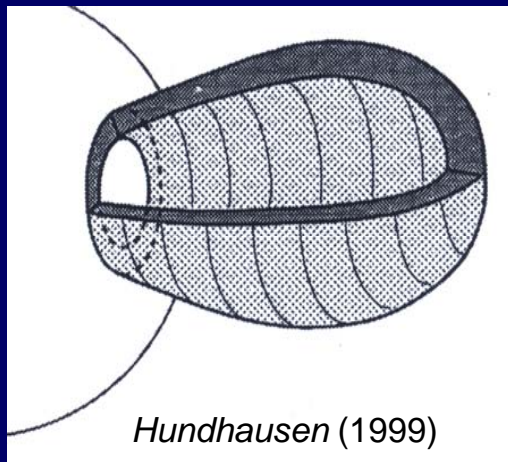
MAGNETIC GEOMETRY UNDERLYING CMES

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Pre-SOHO

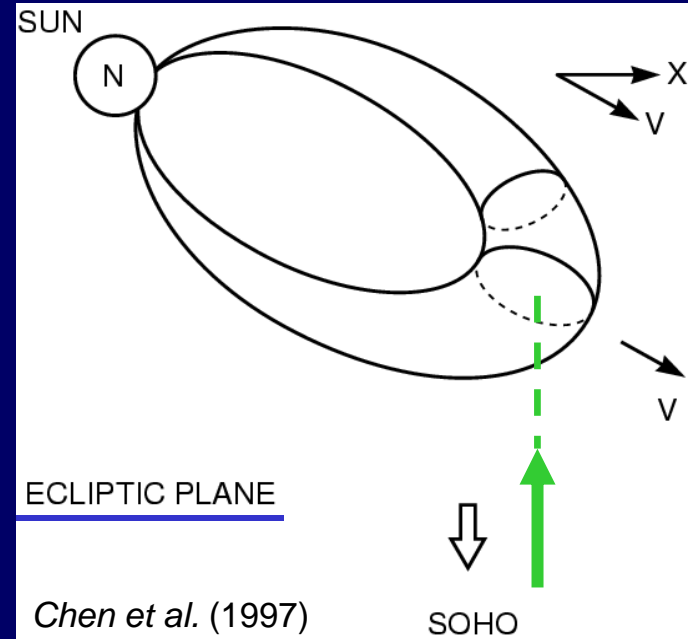


*Illing and
Hundhausen
(1986)*



Hundhausen (1999)

Post-SOHO

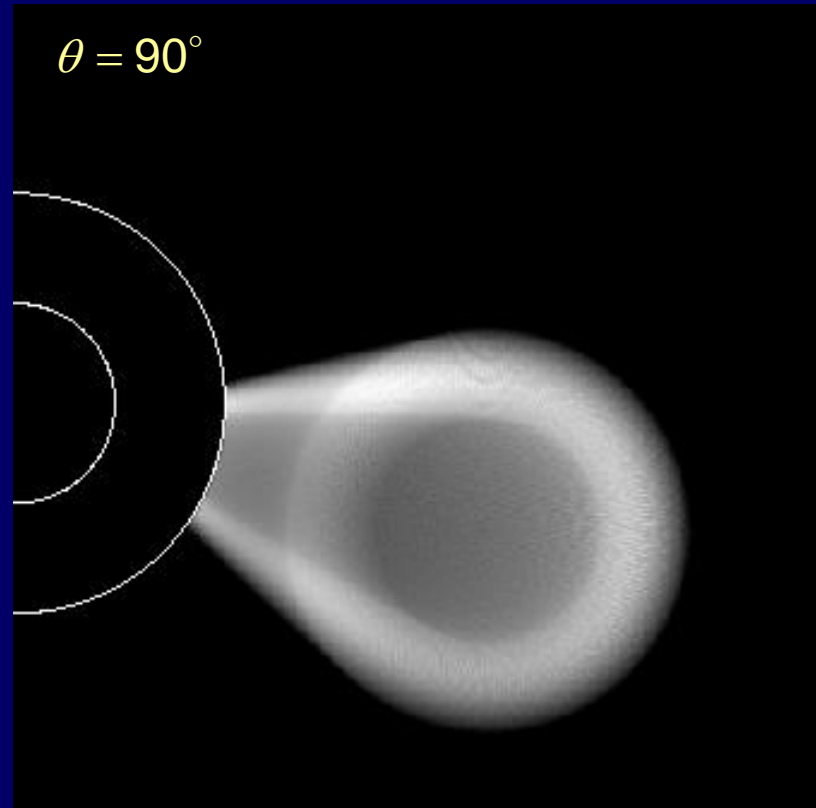
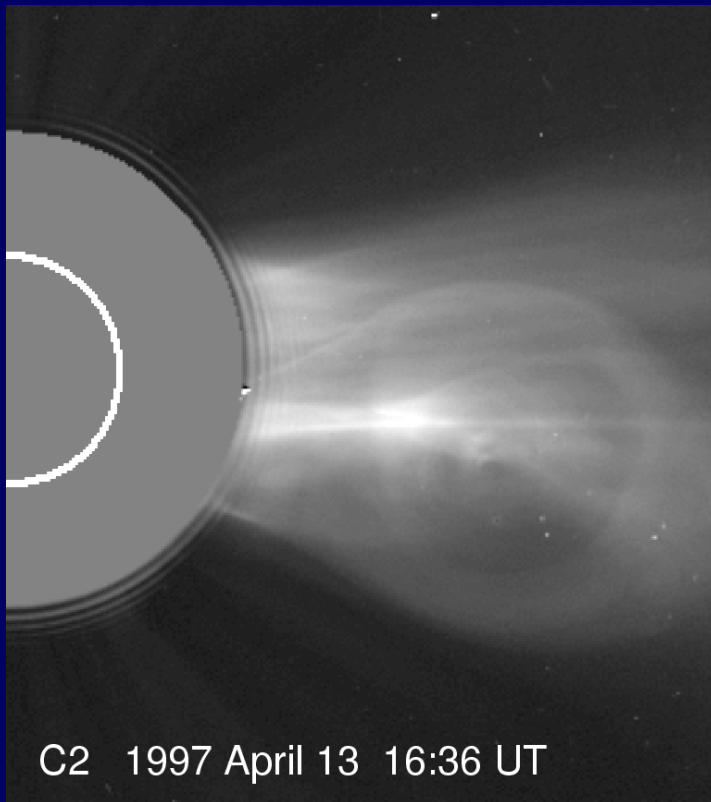


Chen et al. (1997)

SOHO

OBSERVATIONAL EVIDENCE

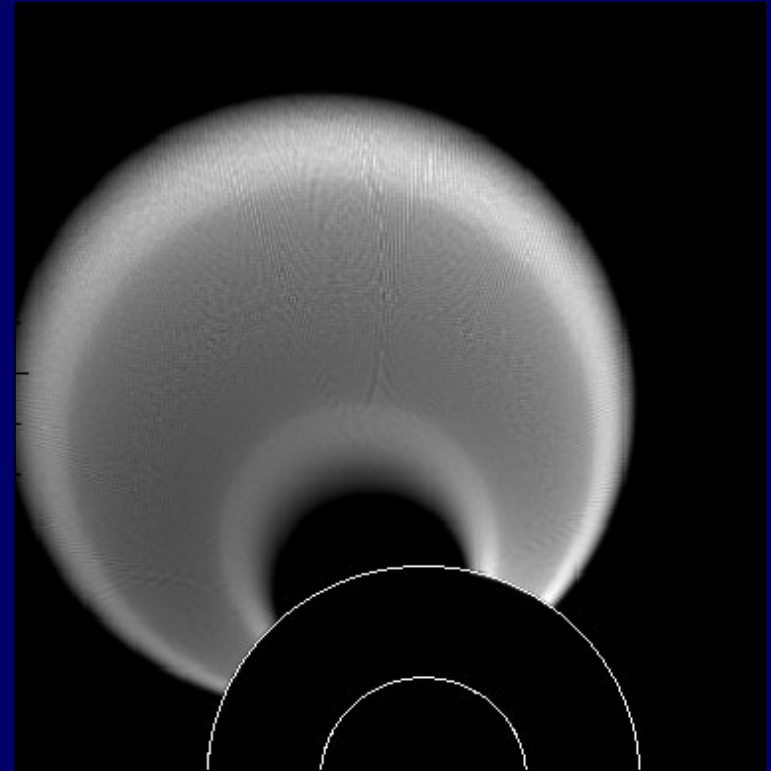
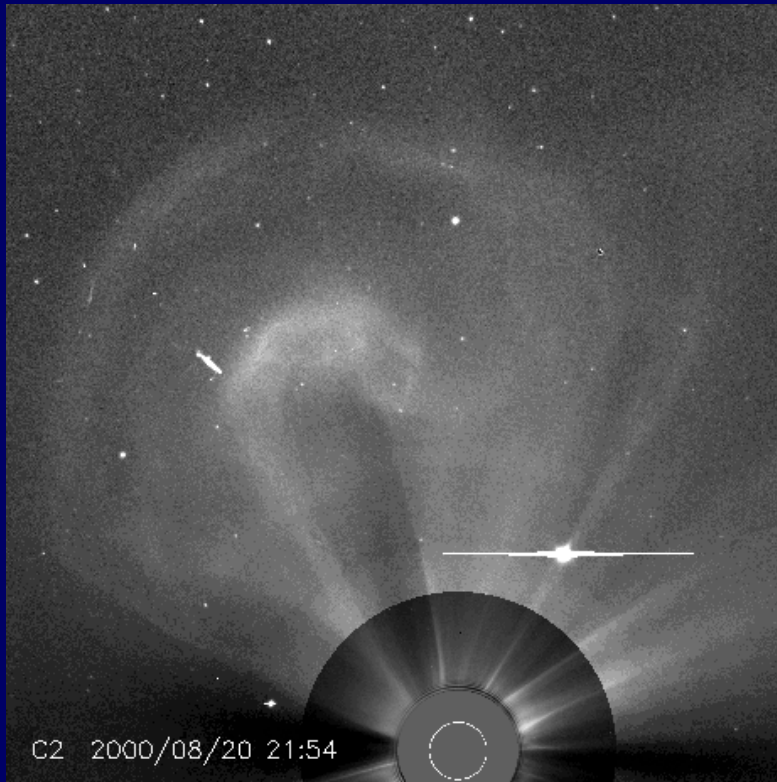
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- Good **quantitative** agreement with a flux rope viewed end-on (*Chen et al. 1997*)
 - No evidence of structural changes attributable to disconnection
- Other examples of flux-rope CMEs (*Wood et al. 1999; Dere et al., 1999; Wu et al. 1999; Plunkett et al. 2000; Yurchyshyn 2000; Chen et al. 2000; Krall et al. 2001; Thernisien et al. 2006*)

OBSERVATIONAL EVIDENCE (cont'd)

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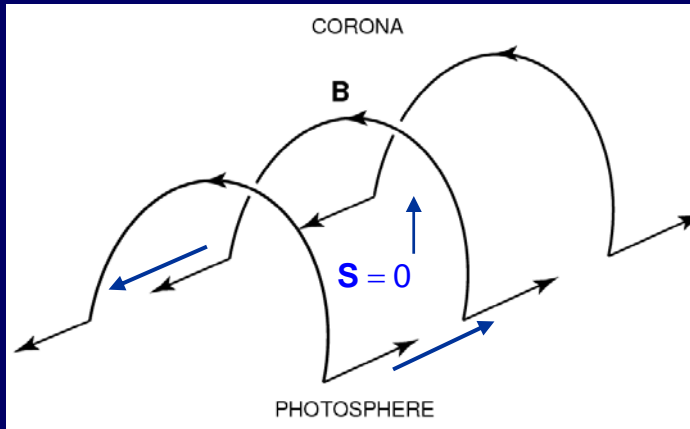


- A flux-rope viewed from the side
- Halo CMEs are flux ropes viewed head on [*Krall et al. 2005*]

THEORETICAL CONCEPTS: TWO MODEL GEOMETRIES

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Magnetic Arcades



Magnetic arcade-to-flux rope

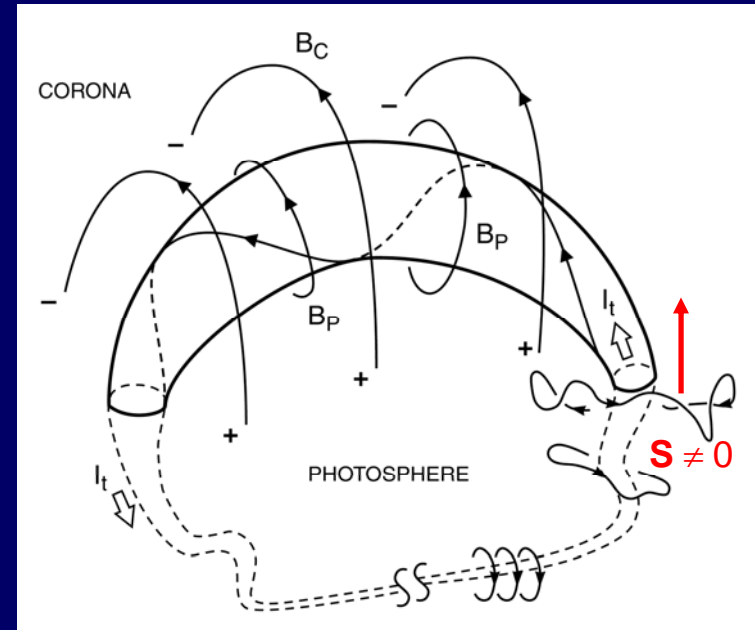
- Energy release and formation of flux rope during eruption

(e.g., *Antiochos et al. 1999; Chen and Shibata 2000; Linker et al. 2001; Lynch et al. 2004, 2009*)

Poynting flux $S = 0$ through the surface

Not yet quantitative agreement with CMEs

Magnetic Flux Ropes



Pre-eruption structure: flux rope with fixed footpoints (S_f) (*Chen 1989; Wu et al. 1997; Gibson and Low 1998; Roussev et al. 2003*)

$S \neq 0$ through the surface (*Chen 1989*)

PHYSICS OF CMEs: Forces

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- “Toroidal” magnetic flux rope *with fixed footpoints* separated by S_f
- Major Radial Forces: integrate $\mathbf{f} = \rho d\mathbf{v} / dt = c^{-1} \mathbf{J} \times \mathbf{B} - \nabla p + \rho \nabla \phi_g$

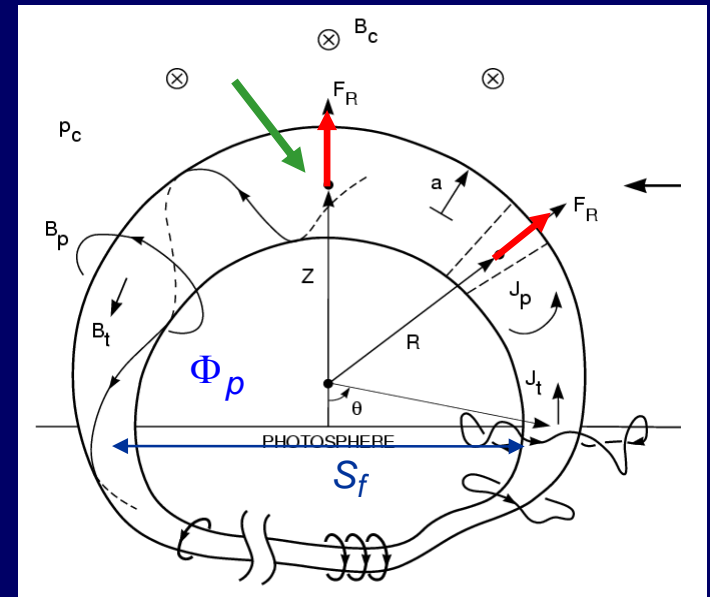
$$\rightarrow M \frac{d^2 Z}{dt^2} = \frac{\Phi_p^2}{c^4 L^2 R} \left[\overset{J_t \times B_p}{\ln\left(\frac{8R}{a}\right)} + \overset{\nabla p}{\frac{1}{2} \beta_p} - \overset{J_p \times B_t}{\frac{1}{2} \frac{B_t^2}{B_p^2}} + \overset{J \times B_c}{2 \left(\frac{R}{a}\right) \frac{B_c}{B_p}} - 1 + \frac{\xi_i}{2} \right] + F_g + F_d$$

[Shafranov 1966; Chen 1989; Garren and Chen 1994]

$$\Phi_p = c L I_t, \quad L = 4\pi \Theta R \left[\ln\left(\frac{8R}{a_f}\right) - 2 \right]$$

- Initiation of eruption:

$$\frac{d\Phi_p(t)}{dt} = \text{poloidal flux "injection"}$$



MINOR RADIAL DYNAMICS

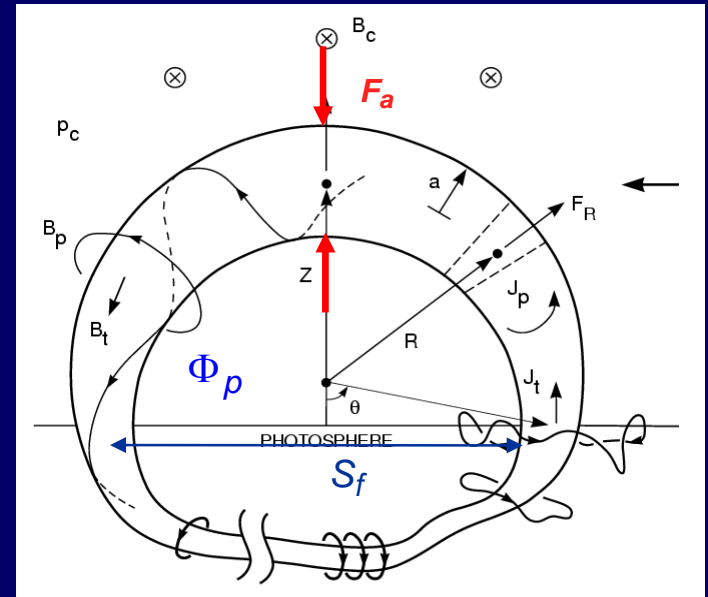
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- Minor Radial Forces: (integrated over a)

$$\rightarrow M \frac{d^2 a}{dt^2} = \frac{a}{4} \left(\overset{J_p \times B_t}{B_t^2} - \underset{J_t \times B_p}{B_p^2} + \overset{\nabla p}{\beta_p B_p^2} \right)$$

- $d^2 a / dt^2 \approx 0$ is a good approximation [Chen 1989]
- Key property of flux-rope Geometry:
Constant S_f is an essential scale length

$$R = \frac{Z^2 + S_f^2 / 4}{2Z}$$



STRUCTURE OF EQUATIONS OF MOTION

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- Shafranov's original work:
 - Forces in major and minor radial directions
 - Axisymmetric toroidal equilibrium
- CMEs: An Extension and New Application
 - Local curvature approximation [*Chen* 1989; *Garren and Chen* 1994]
 - Stationary photospheric footpoints: nonaxisymmetric \longrightarrow additional lengths scales S_f, a_f
 - Dynamical expansion \longrightarrow time scales
 - Momentum coupling to the ambient to the ambient plasma
 - S_f and a_f are directly manifested in *observed* CME acceleration data [*Chen et al.* 2006]
- Comparison with other recent models invoking Shafranov
 - *Wu et al.* [1997] – 2D axisymmetric MHD simulation with stationary footpoints
 - *Lin et al.* [1998], *Titov and Demoulin* [1999], *Kliem and Torok* [2005] – axisymmetric with no footpoints, no minor radial force equation, no coupling to the ambient plasma
 - *Isenberg and Forbes* [2007] – major radial force only, no dynamics
 - *Roussev et al.* [2003], *Torok and Kliem* [2008] -- MHD simulations with fixed footpoints (invoking *Titov and Demoulin* and *Kliem and Torok* but scales are different)

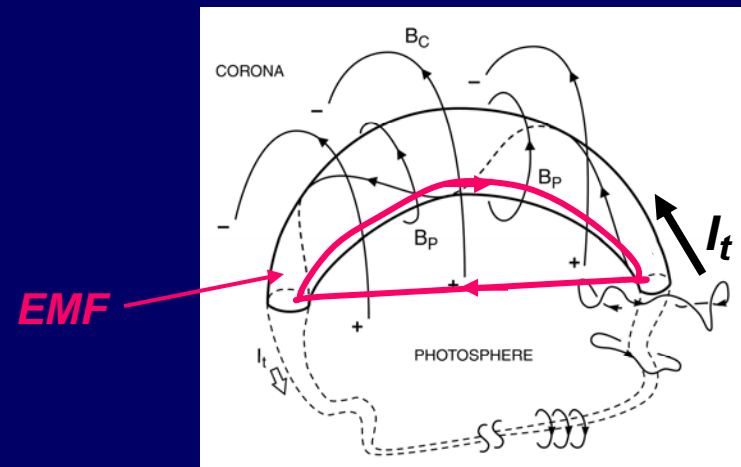
DIRECT COMPARISON OF THEORY AND DATA

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- Previous comparison of theoretical predictions and directly observable quantities
 - Good agreement with observed height and acceleration data
 - Agreement of predicted Sf -scaling law and observed CME acceleration profiles (17 events)
- A new theoretical prediction: the temporal form of $d\Phi_p(t)/dt$ for a CME should be correlated with that of the X-ray emission profile of the associated flare
 - Physics: $-(1/c)d\Phi_p(t)/dt = \text{electromotive force (EMF)} \propto \text{electric field}$

$$M \frac{d^2 Z}{dt^2} = \frac{\Phi_p^2(t)}{c^4 L^2 R} \left[\ln\left(\frac{8R}{a}\right) + \frac{1}{2} \beta_p - \frac{1}{2} \frac{B_t^2}{B_p^2} + 2 \left(\frac{R}{a}\right) \frac{B_c}{B_p} - 1 + \frac{\xi_i}{2} \right] + F_g + F_d$$

$$\boxed{EMF(t) = -\frac{1}{c} \frac{d\Phi_p}{dt}}$$



BEST-FIT SOLUTION

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- Define goodness of fit with respect to height-time data: G

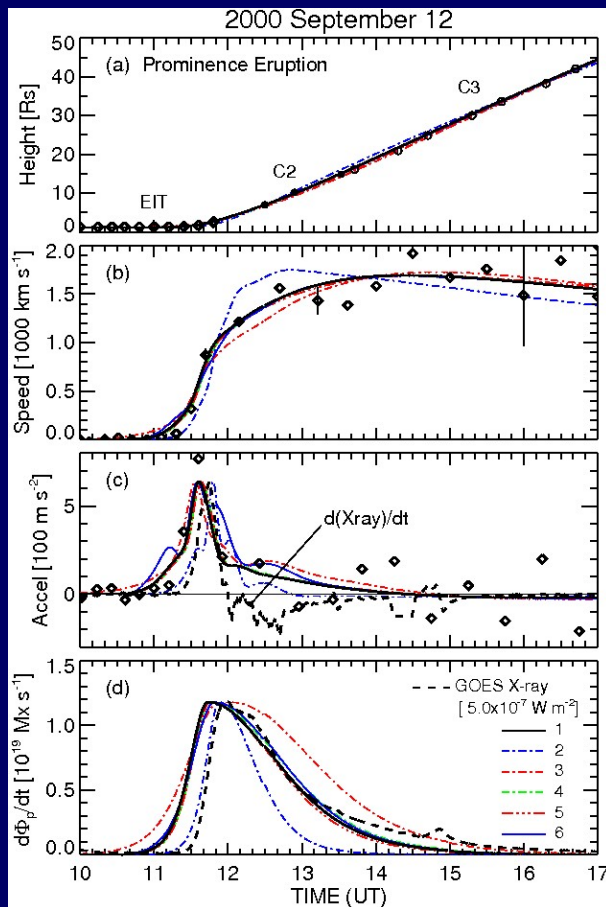
$$G \equiv \frac{1}{T} \sum_{i=1}^N \frac{|Z_{data}(t_i) - Z_{th}(t_i)|}{\Delta Z(t_i)} \delta t_i$$

- Adjust $d\Phi_{\rho}(t)/dt$ to find theoretical solutions that best fit the observed CME height-time data and compare the *calculated* $d\Phi_{\rho}(t)/dt$ with *observed* GOES X-ray data
- Results:
 - The form of $d\Phi_{\rho}(t)/dt$ is strongly constrained by the height data with little freedom
 - Agreement is good for both short- and long-duration flare events
- *Goodness of fit is determined with no regard to speed, acceleration, and X-ray emissions.*

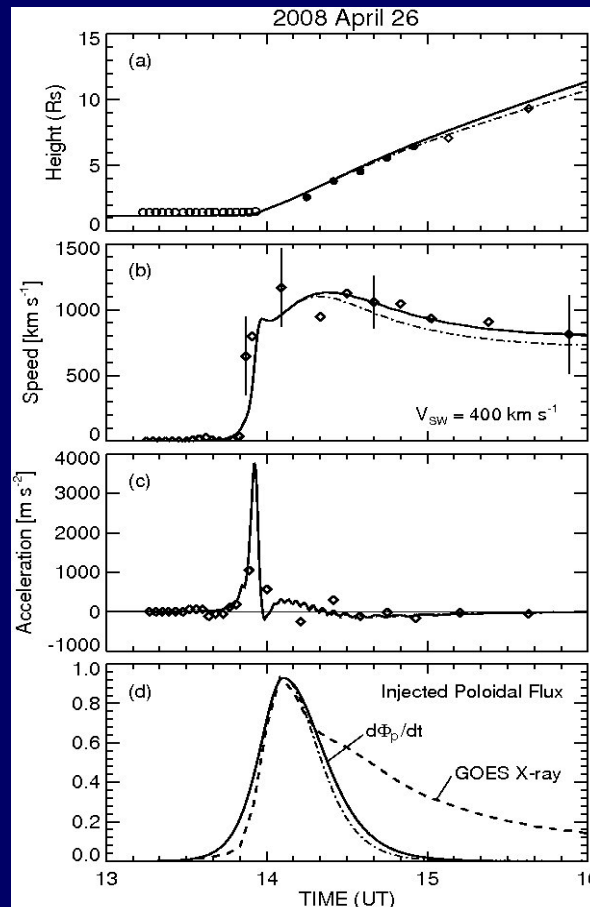
THEORY-DATA COMPARISON

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- Set up initial equilibrium flux rope according to available observational proxies: e.g., S_f , footpoint separation distance, $B_c(Z_0)$. Adjust $d\Phi_p(t)/dt$

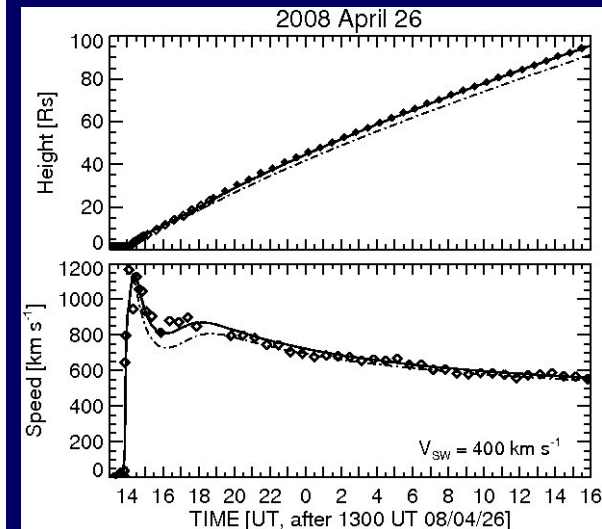


$G = 0.85$ $S_f = 4.5 \times 10^5$ km $E \sim 1$ V/cm



$G = 0.42$ $S_f = 2.0 \times 10^5$ km $E \sim 15$ V/cm

Chen and Kunkel (2009)



Error: 2% of height

Goodness of fit: $G \sim 0.5 - 1.0$

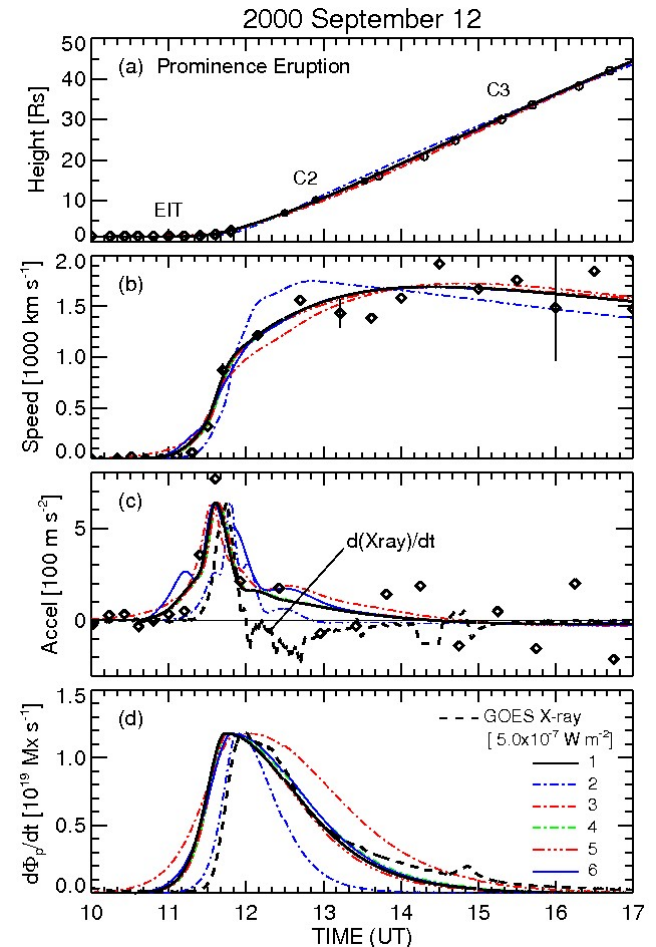
PARAMETER STUDY

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- For each set of parameters, adjust $d\Phi_p(t)/dt$ to obtain the best-fit solution
 - All “best-fit” solutions ($G \sim 0.85-1.2$ for this case) have similar FWHM durations
 - For LASCO heights, the fit is sensitive to the duration but not to V_{SW}

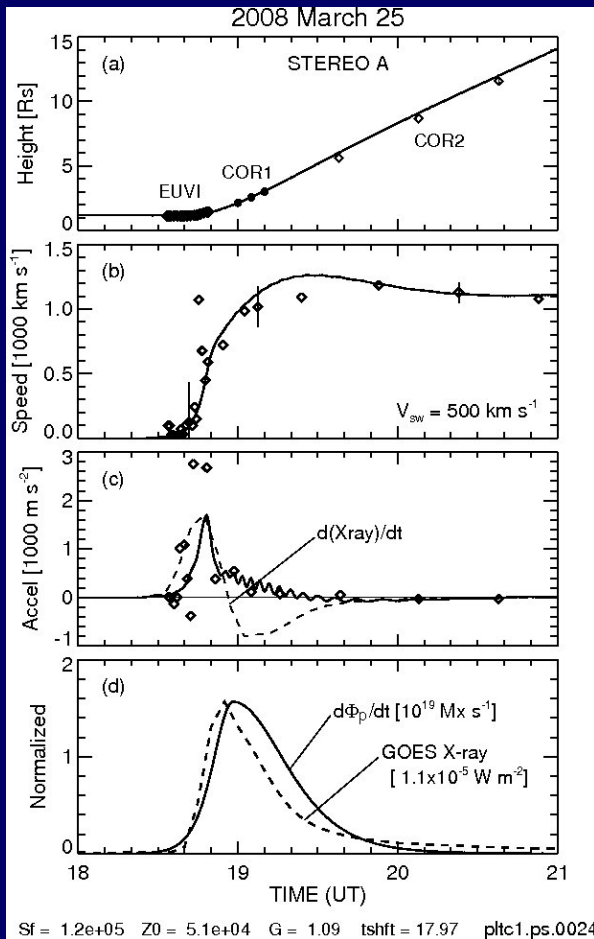
Form of $d\Phi_p(t)/dt$

	File ID	GFIT	Z0	Sf	ASPCT	XPP	NP	BS0	PHIA	DTC	Dtpeal	TSC1	TSC2	Vsw
	240+0	1.03	2.0	4.5	2.5	1.0	1.0	-0.5	9.400	76.0	1.0	25.0	73.0	400
	240+1	0.94	2.0	4.5	2.5	1.0	1.0	-1.5	5.300	76.0	1.0	25.0	75.0	400
	240+2	1.13	2.0	4.5	2.5	1.0	1.0	-1.5	4.500	76.0	2.0	20.0	72.0	400
	250	0.96	1.8	4.5	2.5	1.0	1.0	-1.0	5.850	76.8	0.0	20.0	74.6	400
1	250-1	0.85	2.0	4.5	2.5	1.0	1.0	-1.0	5.600	71.0	0.0	20.0	75.0	400
	250-2	0.97	2.0	4.5	2.0	1.0	1.0	-1.0	6.400	70.0	0.0	20.0	72.5	400
	250-3	0.99	1.8	4.5	2.5	1.0	1.0	-1.0	6.430	70.0	0.0	22.0	73.3	400
	250-4	0.87	2.0	4.5	2.5	1.0	1.0	-1.0	6.150	70.0	0.0	22.0	73.3	400
	252+5	1.18	2.0	4.5	2.0	0.5	2.0	-1.0	6.460	72.0	1.0	20.0	67.9	400
	260+0	0.87	2.0	4.5	2.5	1.0	1.0	-1.0	6.700	71.0	0.0	25.0	76.7	400
	260+1	1.09	2.0	4.5	2.5	1.0	1.0	-1.0	8.000	71.0	0.0	30.0	76.7	400
6	616+2	1.29	2.0	5.5	2.5	1.0	1.0	-1.0	5.600	73.0	0.0	25.0	74.5	400
	0530+0	1.09	2.0	5.0	2.5	1.0	2.0	-1.5	4.500	90.0	0.0	25.0	83.0	400
	0530+1	1.18	2.0	5.0	2.5	1.0	2.0	-1.0	5.530	90.0	0.0	25.0	82.0	400
3	0540+1	1.45	2.0	5.0	2.5	1.0	1.0	-1.0	5.400	130.0	0.0	35.0	89.5	400
	0540+2	2.89	2.0	5.0	2.5	1.0	1.0	-1.5	5.400	110.0	0.0	15.0	30.0	400
2	0540+3	2.19	2.0	5.0	2.5	1.0	1.0	-1.0	5.900	110.0	0.0	15.0	38.0	400
4	0550+1	0.87	2.0	4.5	2.5	1.0	1.0	-1.0	5.590	70.7	0.0	20.0	78.2	300
	0550+2	0.88	2.0	4.5	2.5	1.0	1.0	-1.0	5.590	70.7	0.0	20.0	79.7	250
5	0550+3	0.84	2.0	4.5	2.5	1.0	1.0	-1.0	5.580	71.5	0.0	20.0	70.1	600

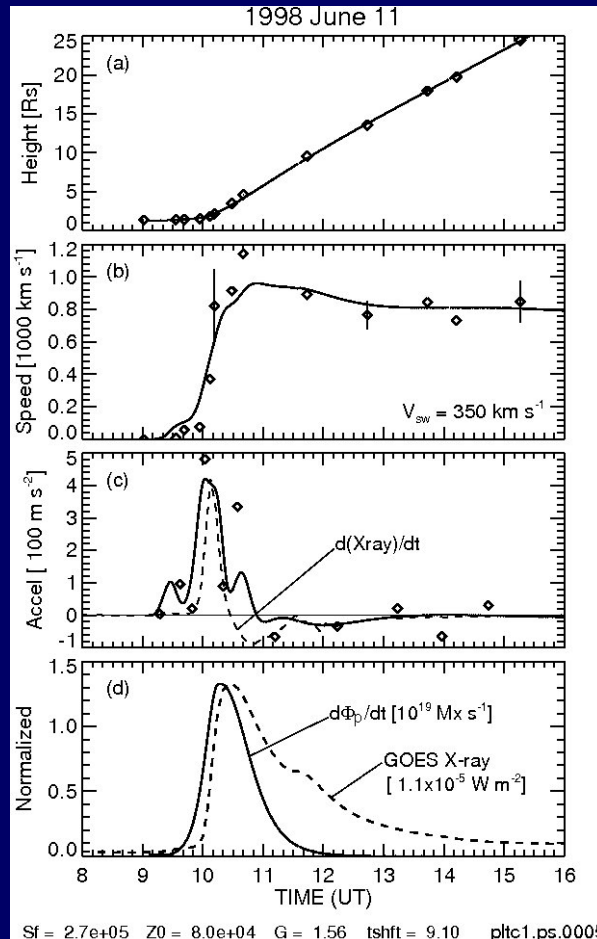


THEORY-DATA COMPARISON

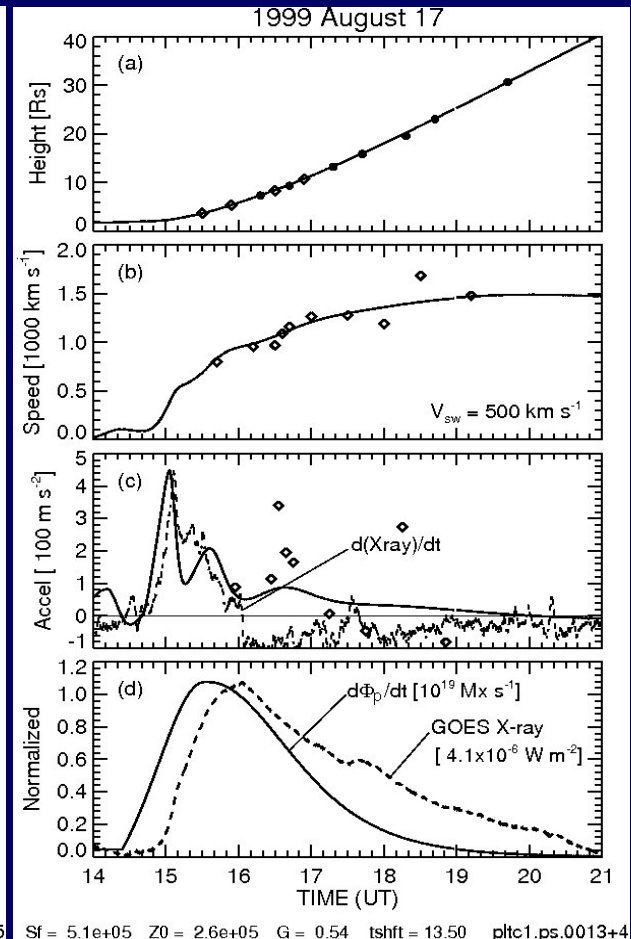
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$E_{max} \sim 12 \text{ V/cm}$



$E \sim 5 \text{ V/cm}$



$E \sim 2 \text{ V/cm}$

- Consistent with observational studies of temporal relationship between acceleration and derivative of soft X-ray: Zhang et al. (2001), Maricic et al. (2007), Temmer et al. (2008)

PHYSICAL INTERPRETATION OF $d\Phi_p(t)/dt$

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- In the toroidal flux rope model, $d\Phi_p(t)/dt$ is a prescribed mathematical function
 - A direct proxy for electric field (super Dreicer) for DC acceleration: $E \sim 0.4\text{--}15 \text{ V cm}^{-1}$
 - Agreement with form of observed X-ray emission profiles is evidence of physical connection linking $d\Phi_p(t)/dt$, CME acceleration, and flare soft X-rays
- Physical interpretation of $d\Phi_p(t)/dt$:
 - (1) Subphotospheric origin via poloidal flux transport from deep source
 - (2) Coronal origin via macroscopic reconnection [*Antiochos et al. 1999; Amari et al. 2000*]
 - Neither has been theoretically or observationally verified
- Comparison with arcade-based coronal storage scenario:
 - 2-D MHD simulation with J -dependent resistivity [*Cheng et al. 2003*]: temporal relation between flux-rope acceleration and inferred energy release with $E \sim 10 \text{ V cm}^{-1}$
 - Estimates of reconnected flux based on photospheric magnetograms:
 - 90 V cm^{-1} [*Qiu et al., 2002*]
 - $0.2\text{--}5 \text{ V cm}^{-1}$ with reconnected flux of $\sim 0.5\text{--}10 \times 10^{18} \text{ Mx s}^{-1}$ [*Jing et al. 2005; Qiu et al. 2007*]

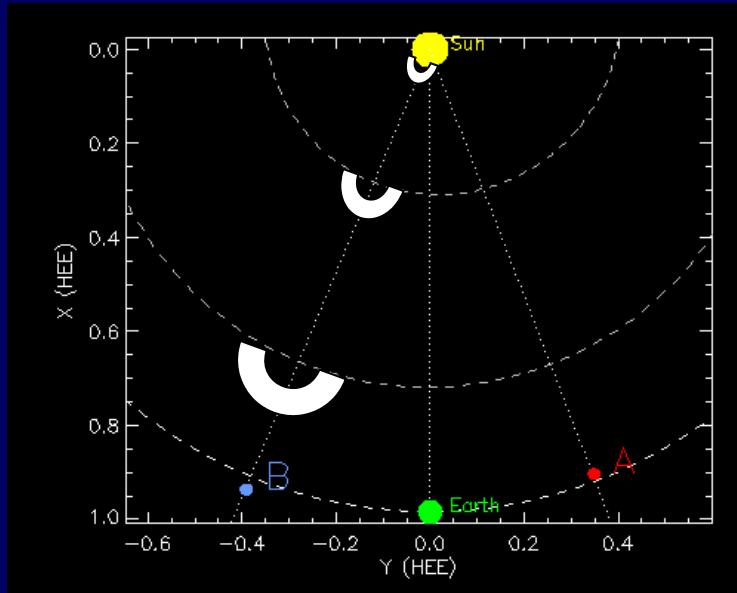
PROPAGATION OF CMEs

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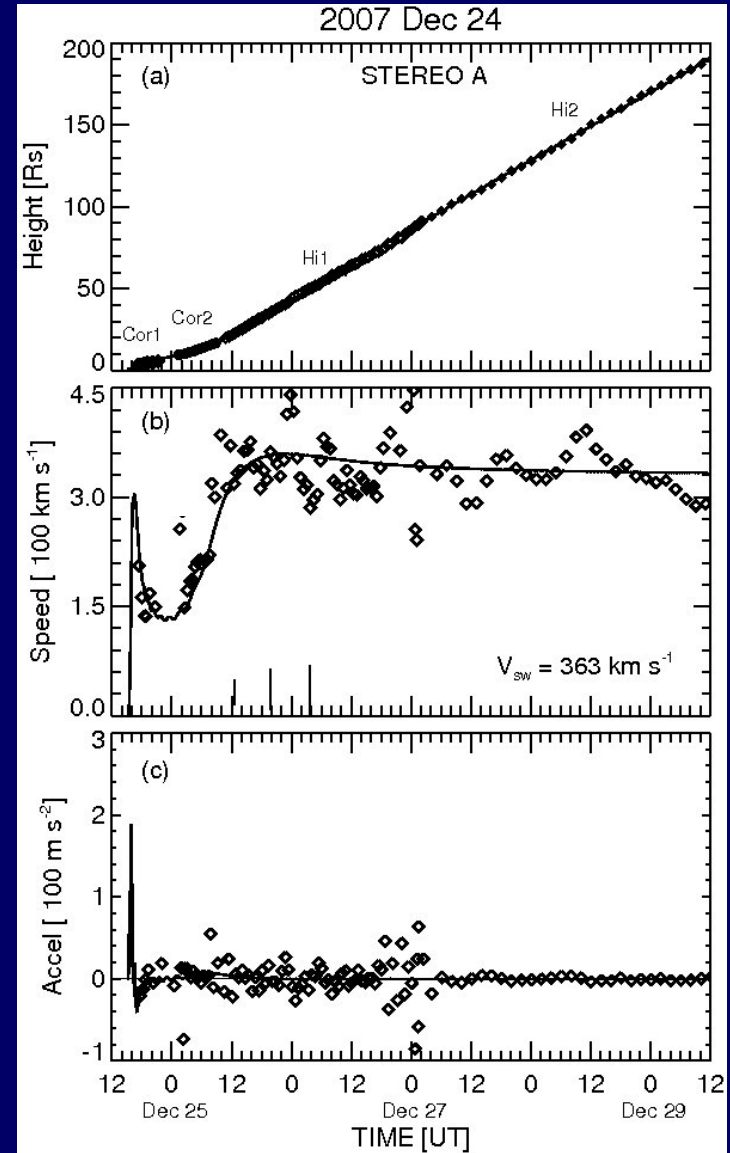
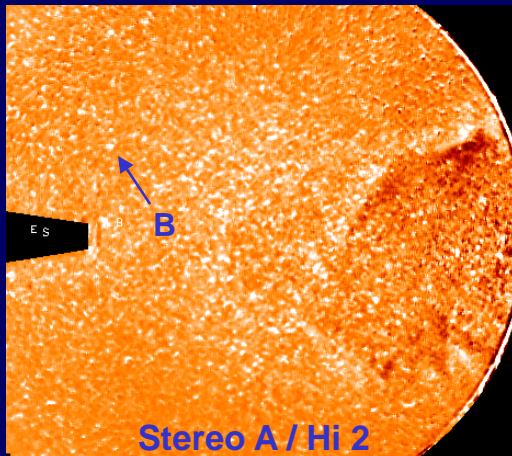
- Model the dynamics of a CME (2007 Dec 24) from initiation to 1 AU (STEREO A data). Predict magnetic field at 1 AU and compare with data (STEREO B data). [*Kunkel and Chen*, in preparation, 2009]
- Two situations:
 - Source region can be observed—obtain proxies for S_f , Z_0 , etc.
 - Source region not observed—adjust S_f , Z_0 , and fit model solutions to HI1/HI2 height-time data \longrightarrow predict \mathbf{B} field at 1 AU
 - For both situations, $d\Phi_\rho(t) / dt$ is an adjustable parameter that can be validated using GOES X-ray data

24 DEC 2007 CME

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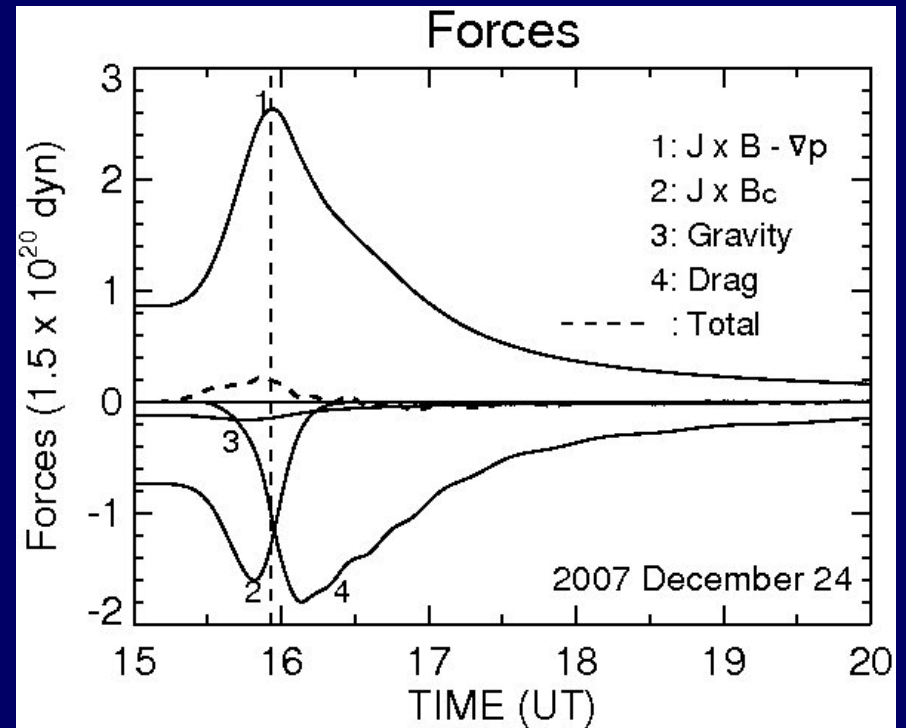
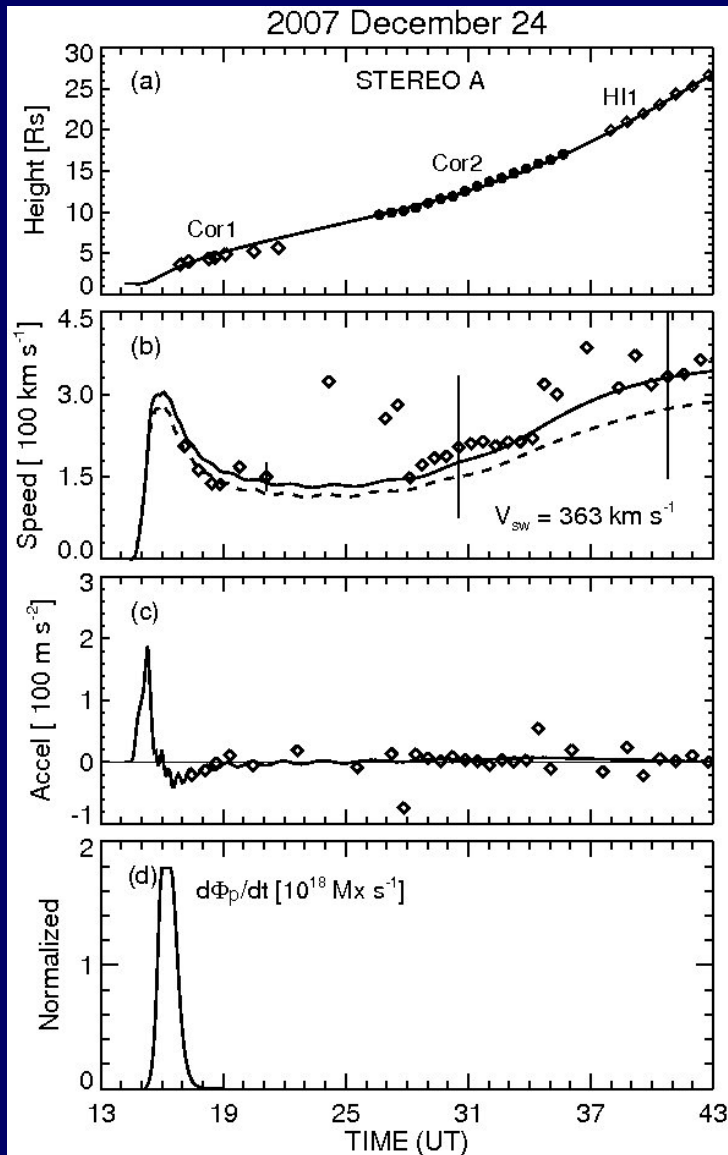


Separation angle Stereo A and B is 44 degree



CME TRAJECTORY: NEAR SUN

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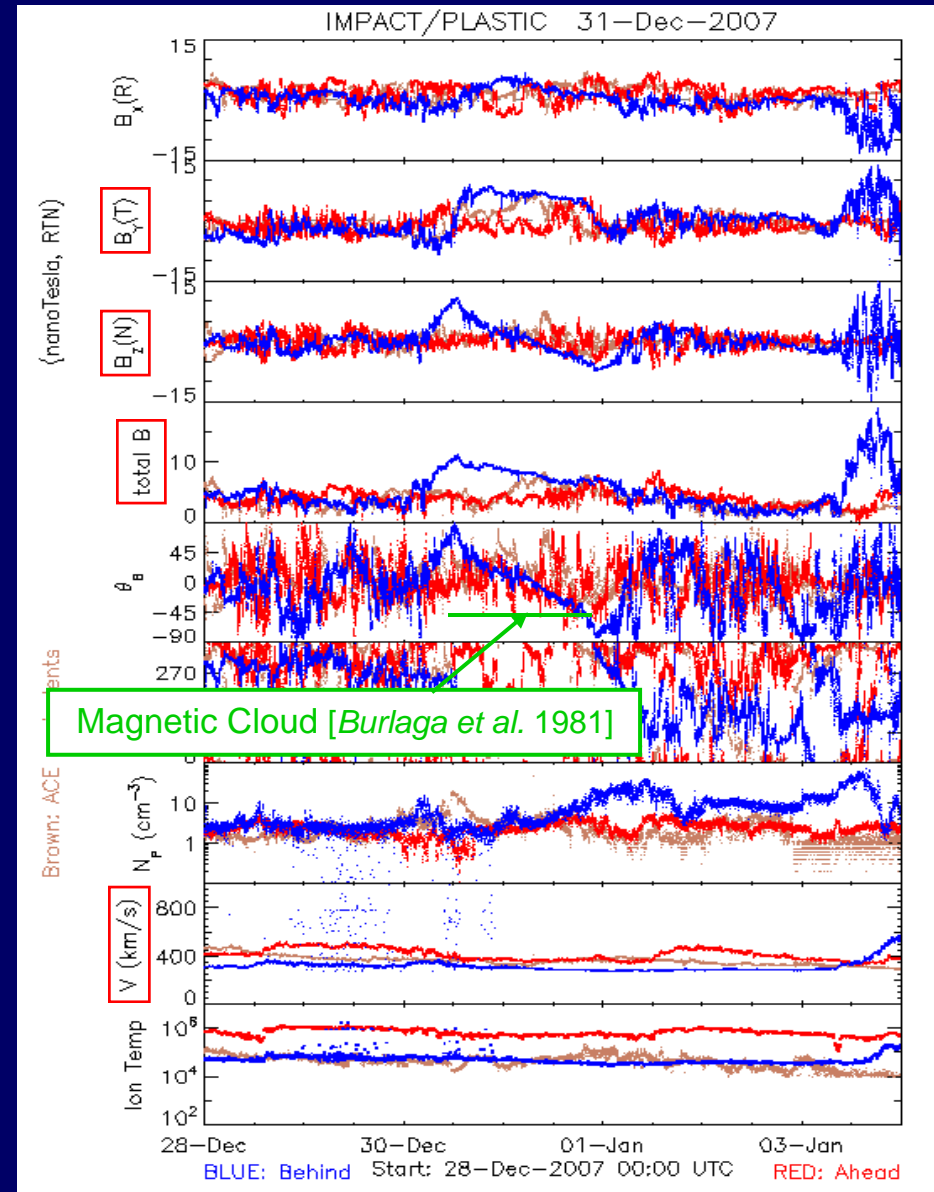
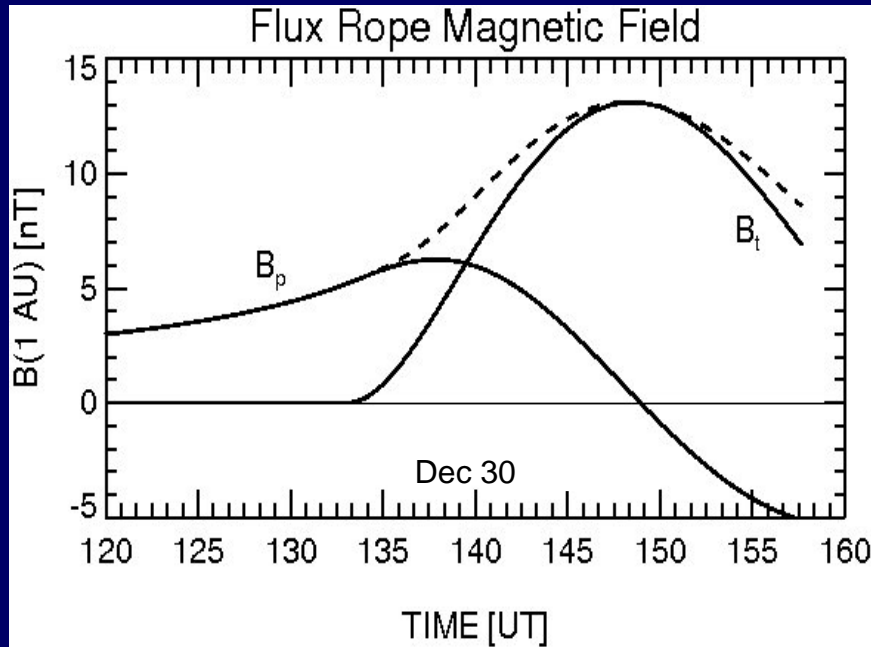
Inclusion of drag in the force equation is essential for the long-time propagation

PREDICTED MAGNETIC FIELD AT 1 AU

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Calculated magnetic field at 1 AU

- Comparison with IMPACT/PLASTIC data



THEORY SUCCESSES

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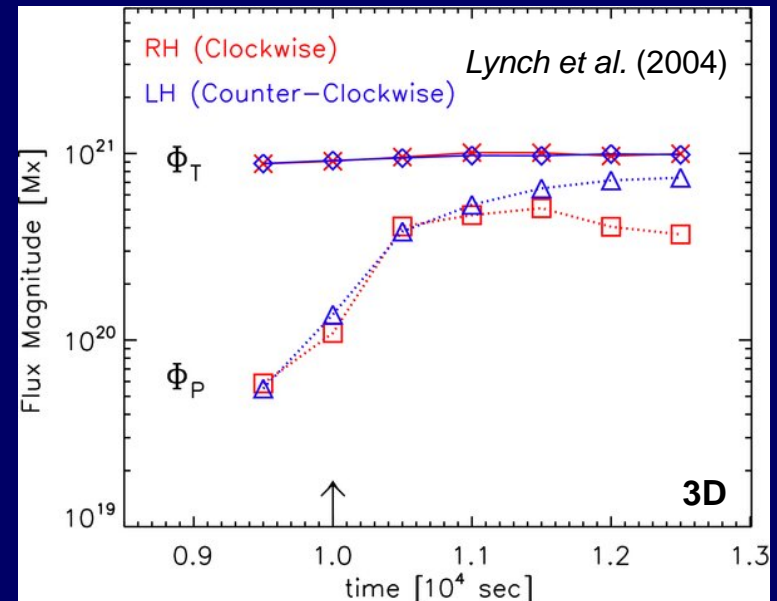
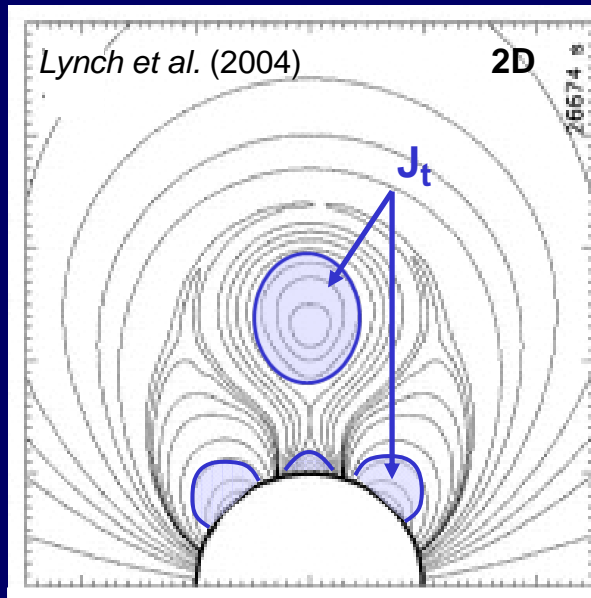
- CME dynamics are described by a set of two ODEs
 - Calculated dynamics have been compared with LASCO and STEREO data
 - Both major radial and minor radial expansion is correctly described by the theory
 - The main acceleration and the subsequent propagation to 1 AU are correctly captured
 - *Drag coupling between CMEs and the ambient SW is essential*
 - The calculated \mathbf{B} field at 1 AU is in agreement with *in situ* measurement at 1 AU (1 event)
 - The best-fit solution yields a temporal profile of $d\Phi_p(t)/dt$ in agreement with the **observed** profile of GOES soft X-ray emissions (five CME-flare events)
- Suggests a new theoretical framework of understanding CME dynamics and flare energy release
 - An initial flux rope is set into motion by injection of **poloidal** flux, which generates an EMF and attendant electric field to accelerate particles to X-ray energies

PHYSICS OF POLOIDAL FLUX INJECTION

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- Currently, $d\Phi_p(t)/dt$ is a specified function of time
- Two physical interpretations are possible:
 - Coronal origin: macroscopic reconnection is required. All models use numerical and/or artificial dissipation. Not yet simulated acceleration in agreement with data.
 - Subphotospheric origin: Not yet observed. Observable photospheric signatures not yet modeled. Favorable if coronal reconnection is not fast enough
- $d\Phi_p(t)/dt$ is a point of overlap between the two basic paradigms (arcade v. flux rope)
[Chen 1996; Chen and Krall 2003]

Breakout
model
simulation



OPEN ISSUES

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- Both arcade models and erupting flux-rope model with poloidal flux injection require further work
- ***Major Physics Issues***
- Arcade models:
 - Physical reconnection on macroscopic scales
 - Demonstration of specific realistic photospheric motion for observed eruptions
 - Calculation of acceleration and speed in agreement with observed CMEs
- Poloidal flux injection model:
 - Demonstration of photospheric signatures in agreement with ***well-resolved*** observation
 - Simulation of subphotospheric plasma dynamics

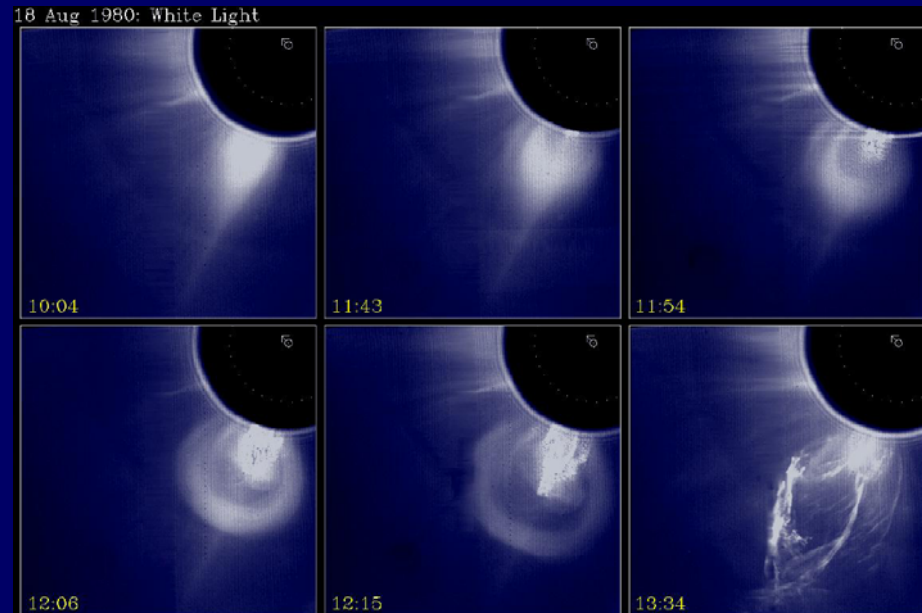
END

3-D GEOMETRY OF CMEs

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- “Coronal transients” (1970’s: OSO-7, Skylab)
- “Thin” flux tubes
(Mouschovias and Poland 1978; Anzer 1978)
- Halo CMEs (*Solwind*) (Howard et al. 1982)
 - Fully 3-D in extent
- CME morphology (*SMM*):
(Illing and Hundhausen 1986)
 - A CME consists of 3-parts: a bright frontal rim, cavity, and a core
 - Conceptual structure: **rotational symmetry** (e.g., ice cream cone, light bulb) (Hundhausen 1999)
- *SOHO* data: 3-D flux ropes (Chen et al. 1997)
 - 3-part morphology is only part of a CME

FOV: $1.7 - 6 R_s$



Source: High Altitude Observatory/Solar Maximum Mission Archives

HAO A-013

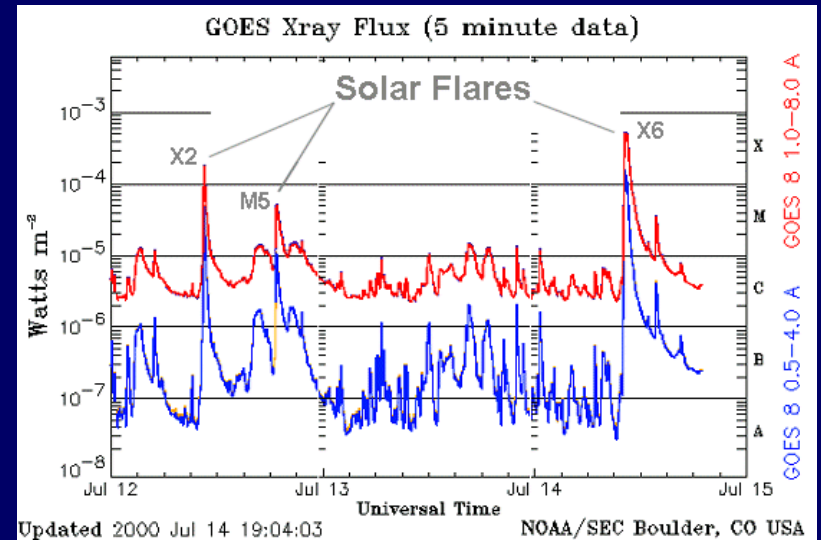
(Illing and Hundhausen (1986))

SMM (1980-1981; 1984-1989)

SOLAR ERUPTIONS: PHENOMENOLOGY

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- Sporadic eruptions
 - Solar flares seen in X-rays, EUV, $H\alpha$, etc.
 - Filament/prominence eruptions seen in $H\alpha$ or white light
 - CMEs in white light
 - All can be accompanied by solar energetic particles (SEPs)



- Solar flares are usually identified by the disk-integrated X-ray emissions detected by GOES satellites
- Stellar flares are recognized by similar X-ray light curves

