



Academy of Athens
Office for Space Research



Imperial College
London



Particle pressure radial profile in the
dayside magnetosphere of Saturn
during near-radial parts of Cassini's trajectory.

N. Sergis⁽¹⁾, C.S. Arridge^(2,3), S.M. Krimigis^(1,4),
D.G. Mitchell⁽⁴⁾, D.C. Hamilton⁽⁵⁾, N. Krupp⁽⁶⁾,
E.C. Roelof⁽⁴⁾, M.K. Dougherty⁽⁷⁾ and A.J. Coates^(2,3).

- (1) Office for Space Research and Technology, Academy of Athens, Athens, GR.
- (2) Mullard Space Science Laboratory, University College London, UK.
- (3) Centre for Planetary Sciences, University College London, UK.
- (4) Applied Physics Laboratory, Johns Hopkins University, Laurel, Maryland, USA.
- (5) Department of Physics, University of Maryland, College Park, Maryland, USA.
- (6) Max Planck Institute for Solar System Research, Lindau, Germany.
- (7) Space and Atmospheric Physics Group, Imperial College, London, UK.

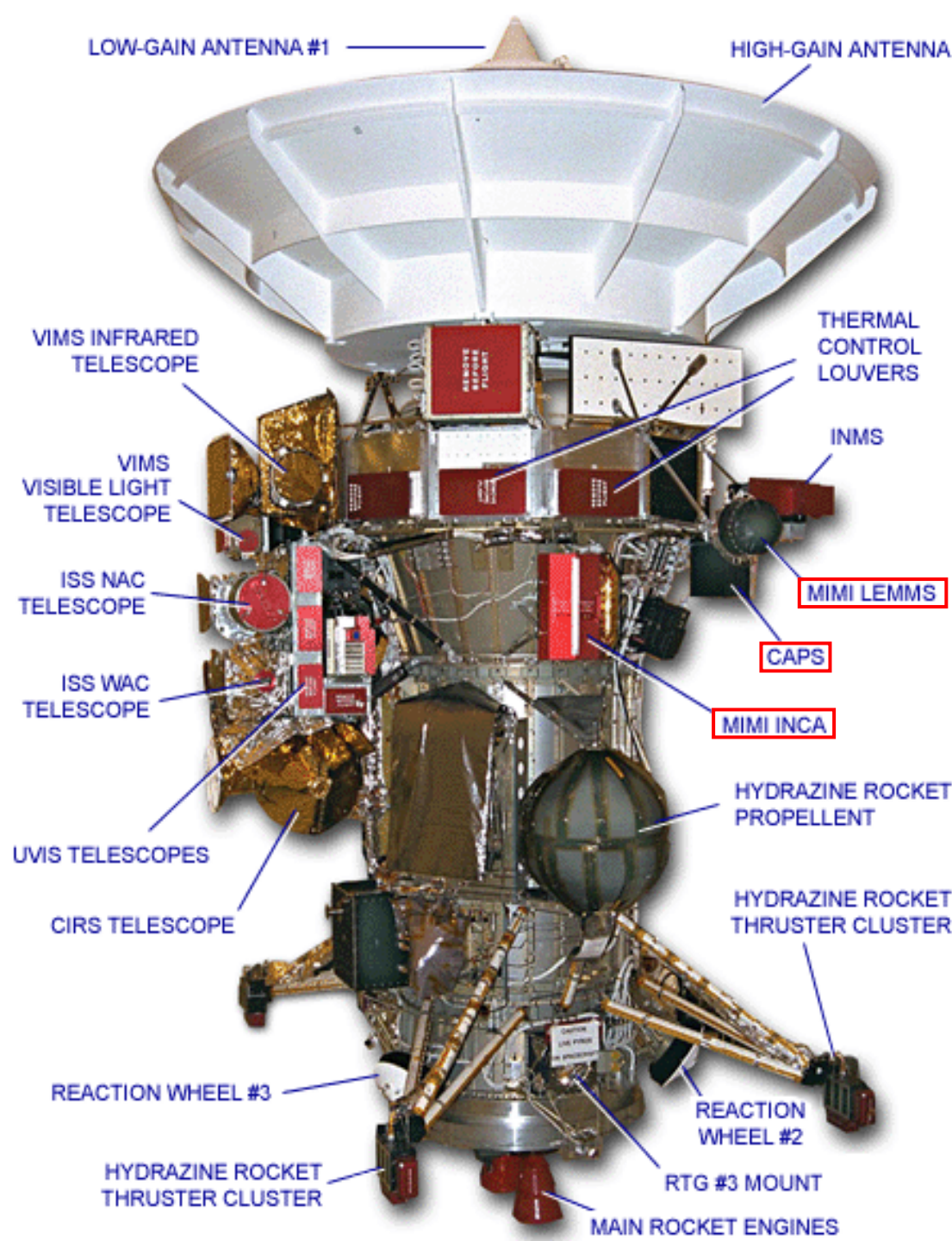
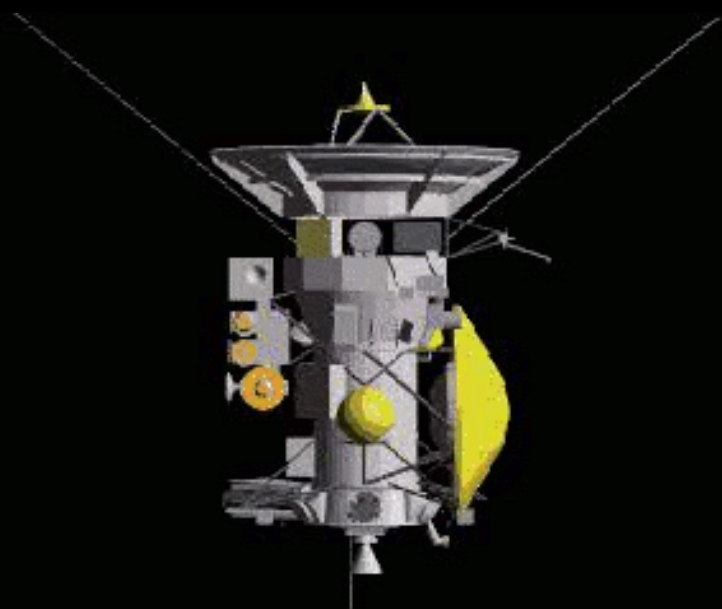
The Cassini spacecraft

Characteristics:

6.7m x 4.0m

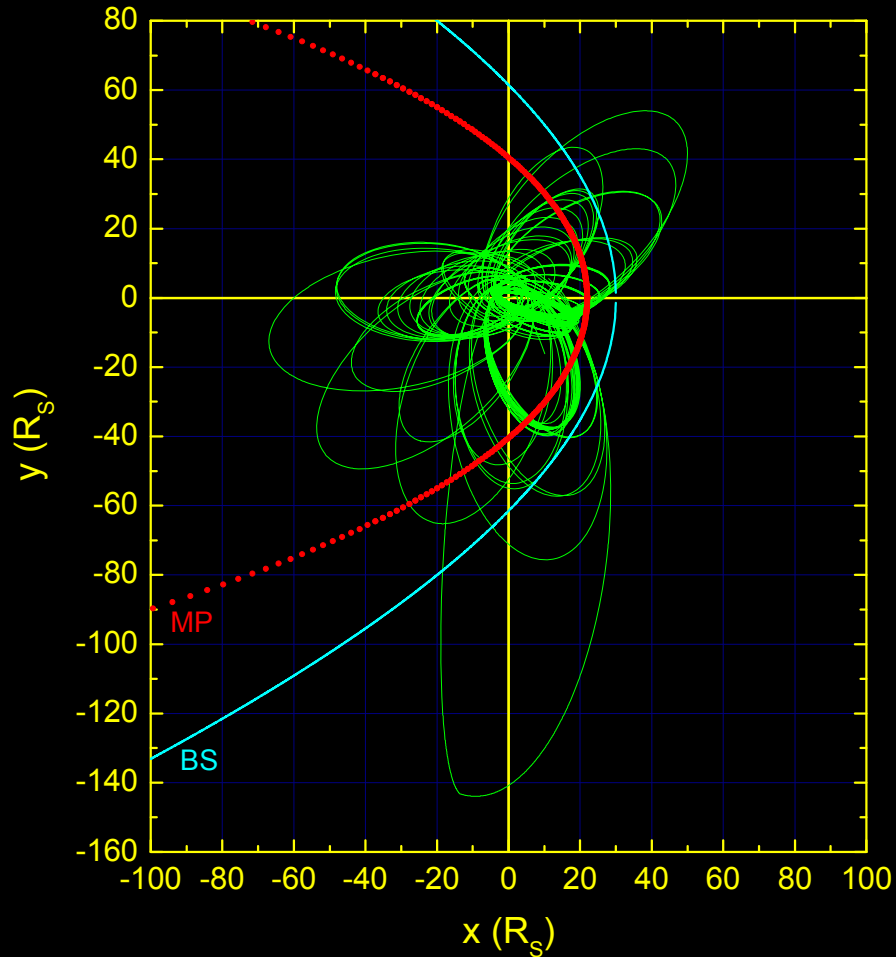
5.7 tons

12 instruments onboard

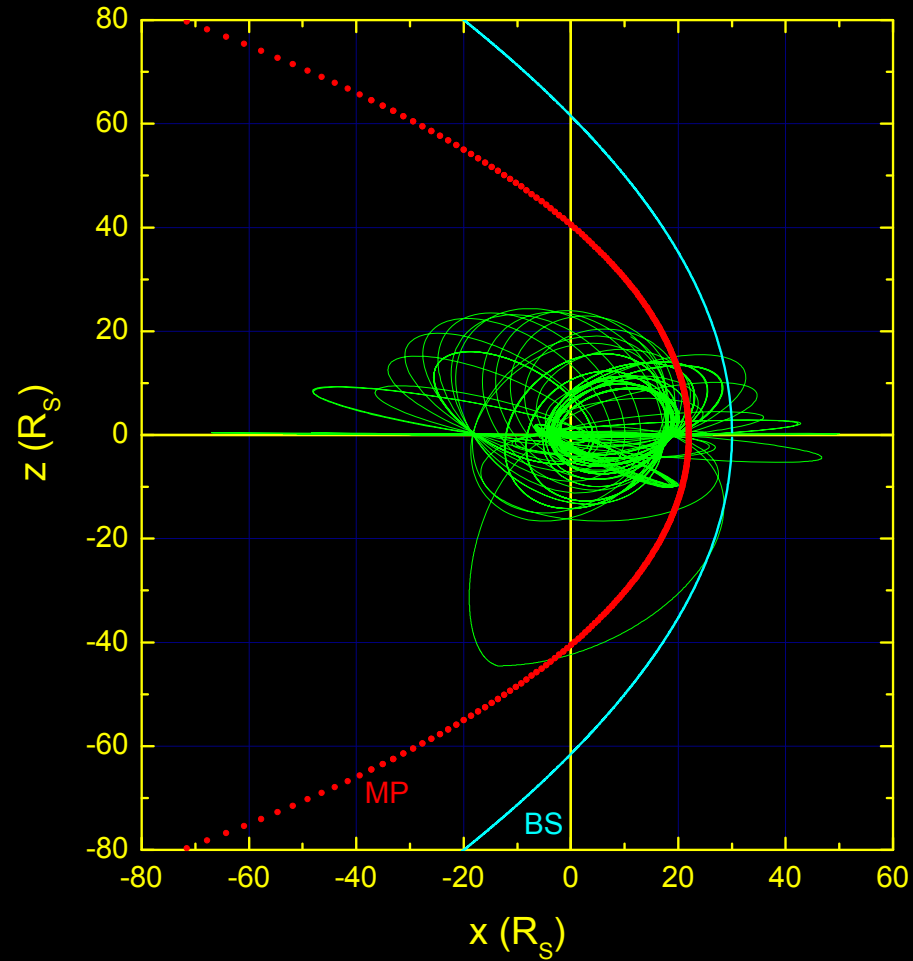


Cassini trajectory during the First 5 years of orbiting Saturn (July 2004-May 2009)

top view



side view



Questions to be addressed

1. Can we obtain a representative radial (total) pressure profile for the dayside magnetosphere of Saturn?
2. How are plasma, suprathermal (keV) and magnetic pressure compared in the Saturnian ring current region?
3. Is the Saturnian ring current inertial or pressure-gradient driven?
4. How well can the ring current density be reproduced by models?

Theory

Assuming that all ion components have the same bulk velocity in the steady-state, the force balance equation can be written as:

$$\rho \mathbf{V} \nabla \mathbf{V} + \nabla \cdot \mathbf{P} - \mathbf{J} \times \mathbf{B} = 0$$

ρ : plasma mass density
 \mathbf{V} : plasma bulk velocity
 \mathbf{P} : total plasma pressure
 \mathbf{J} : current density

Assuming that the pressure is isotropic and that $\mathbf{V} = \Omega \times \mathbf{r}$ (strict corotation), the radial component of the equation will be:

$$-\rho \Omega^2 r + \frac{\partial P}{\partial r} - (J_\theta B_\phi - J_\phi B_\theta) = 0$$

Ω : Saturn's rotational angular velocity
 J_ϕ : azimuthal current density

and as long as $B_\theta \gg B_\phi$ (equatorial plane orbits)

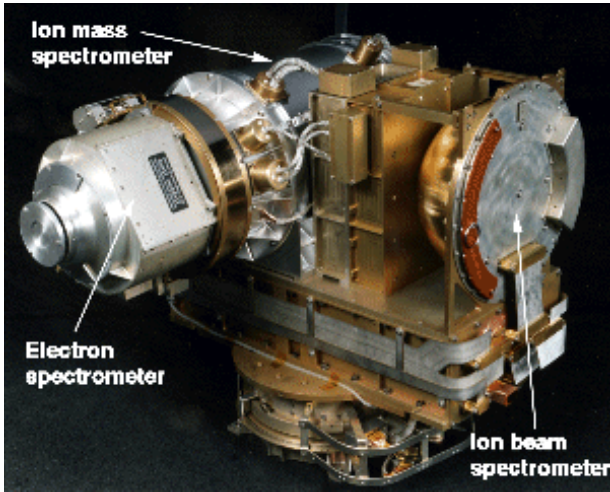
$$J_\phi \approx \frac{1}{B_\theta} \left(\rho \Omega^2 r - \frac{\partial P}{\partial r} \right)$$

Inertial
contribution

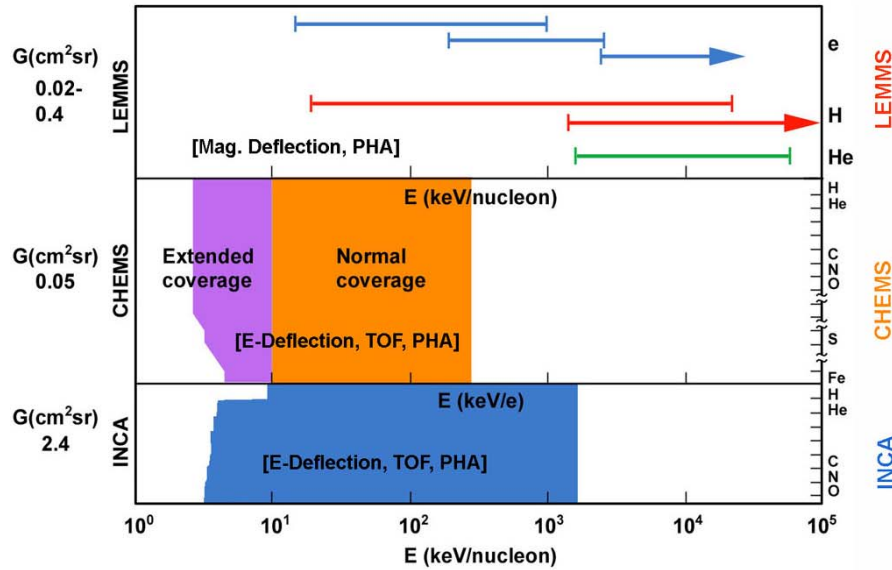
Pressure gradient
contribution

Instrumentation

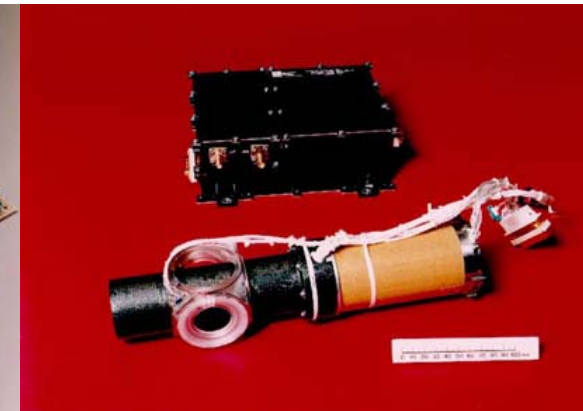
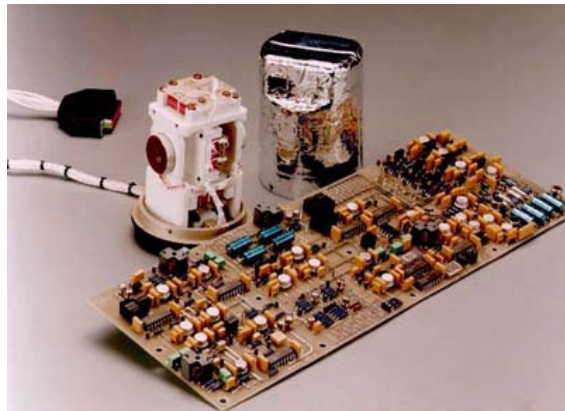
Plasma: CAPS instrument (IMS, ELS) energy range 1 eV to few keV.



Energetic particles: MIMI instrument (CHEMS, LEMMS and INCA). $E > 3$ keV coverage (ions), $E > 20$ keV (e^-), composition, directional intensities (pitch angle measurements).



Magnetic field:
Cassini magnetometer (MAG)
High resolution magnetic field
Vector measurements.
(4 sec sampling, pT level)



Data selection

All dayside Cassini passes were examined

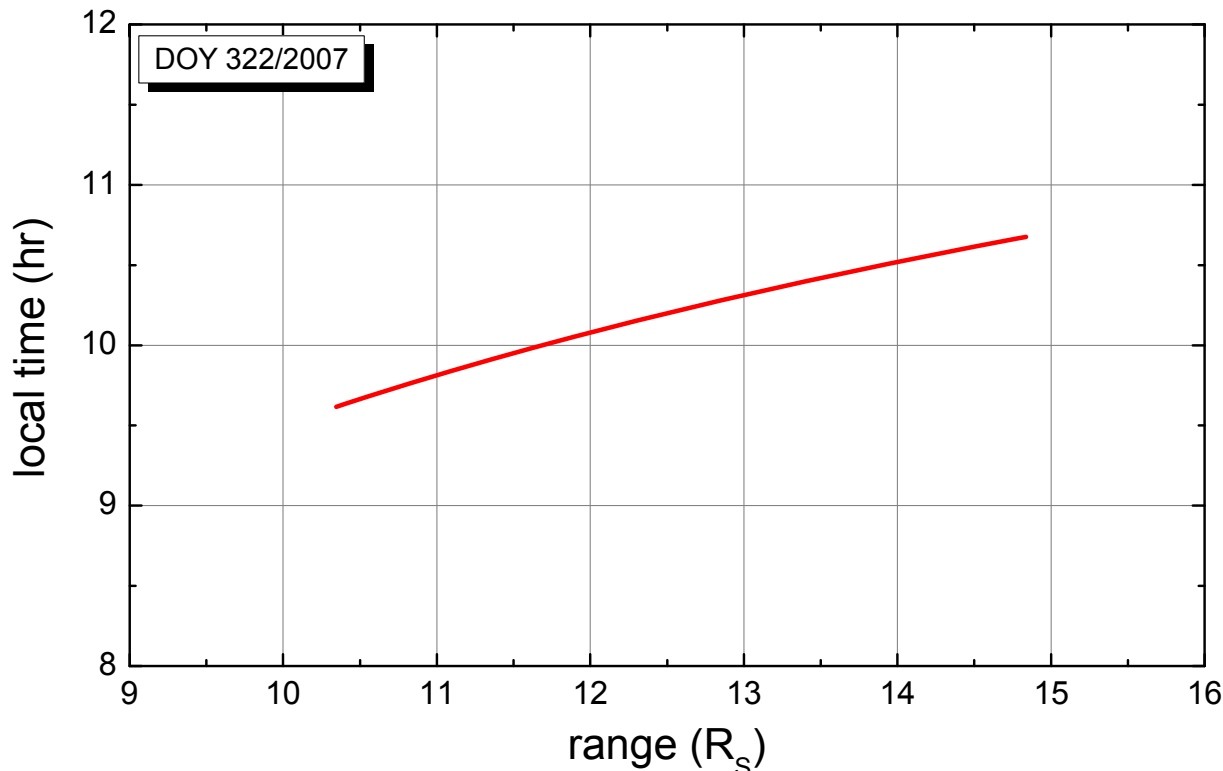
→ Equatorial plane orbits ($|z| < 1 R_S$)

→ Minimum local time change $\frac{d(LT)}{dr}$

→ Available CAPS electron moments

Less than 5 cases

DOY 322/2007



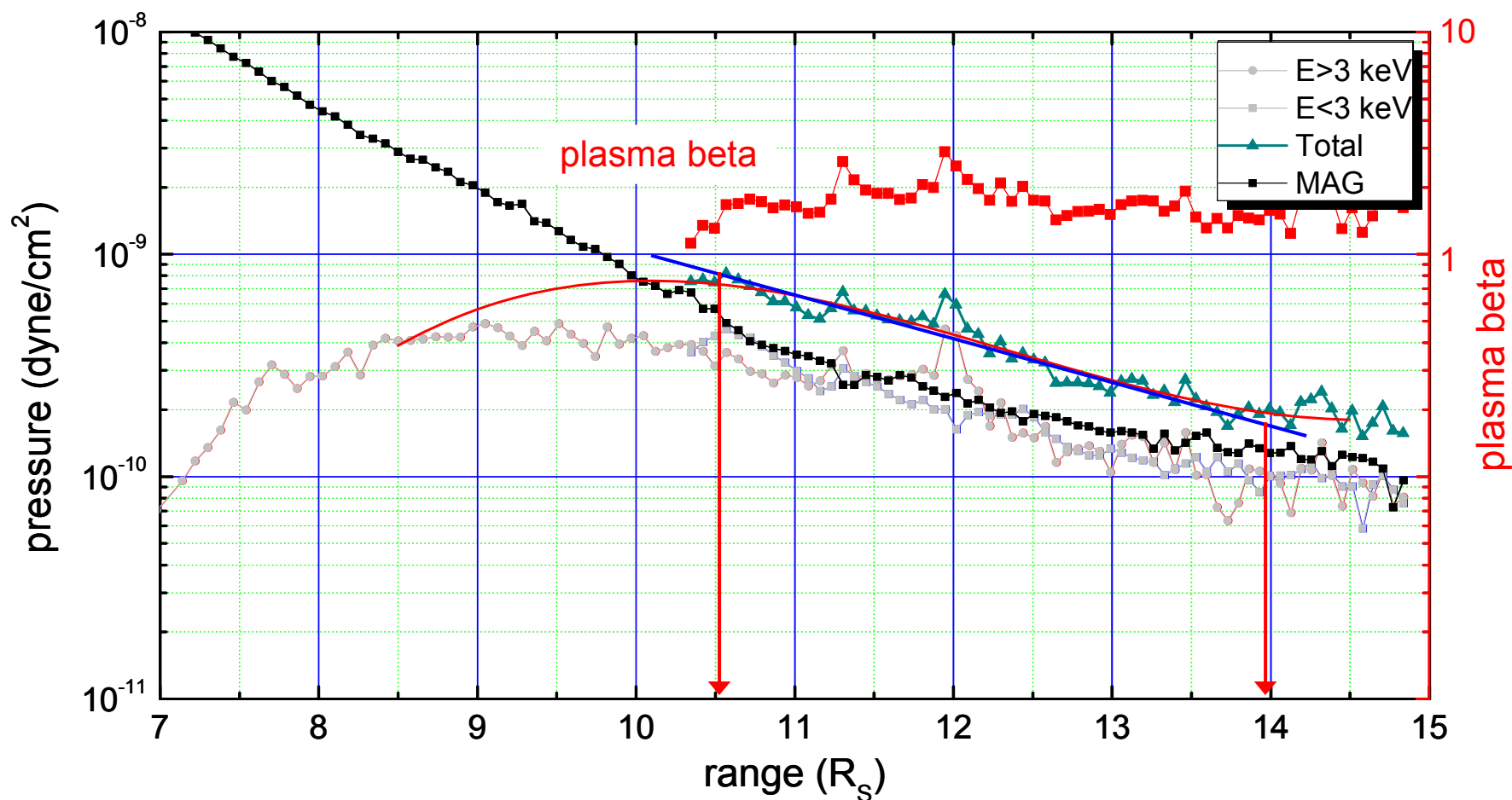
$$\frac{d(LT)}{dr} \approx 0.20 - 0.25 \text{ hr}/R_S$$

Magnetosheath or near-magnetosheath data have been excluded.

Results

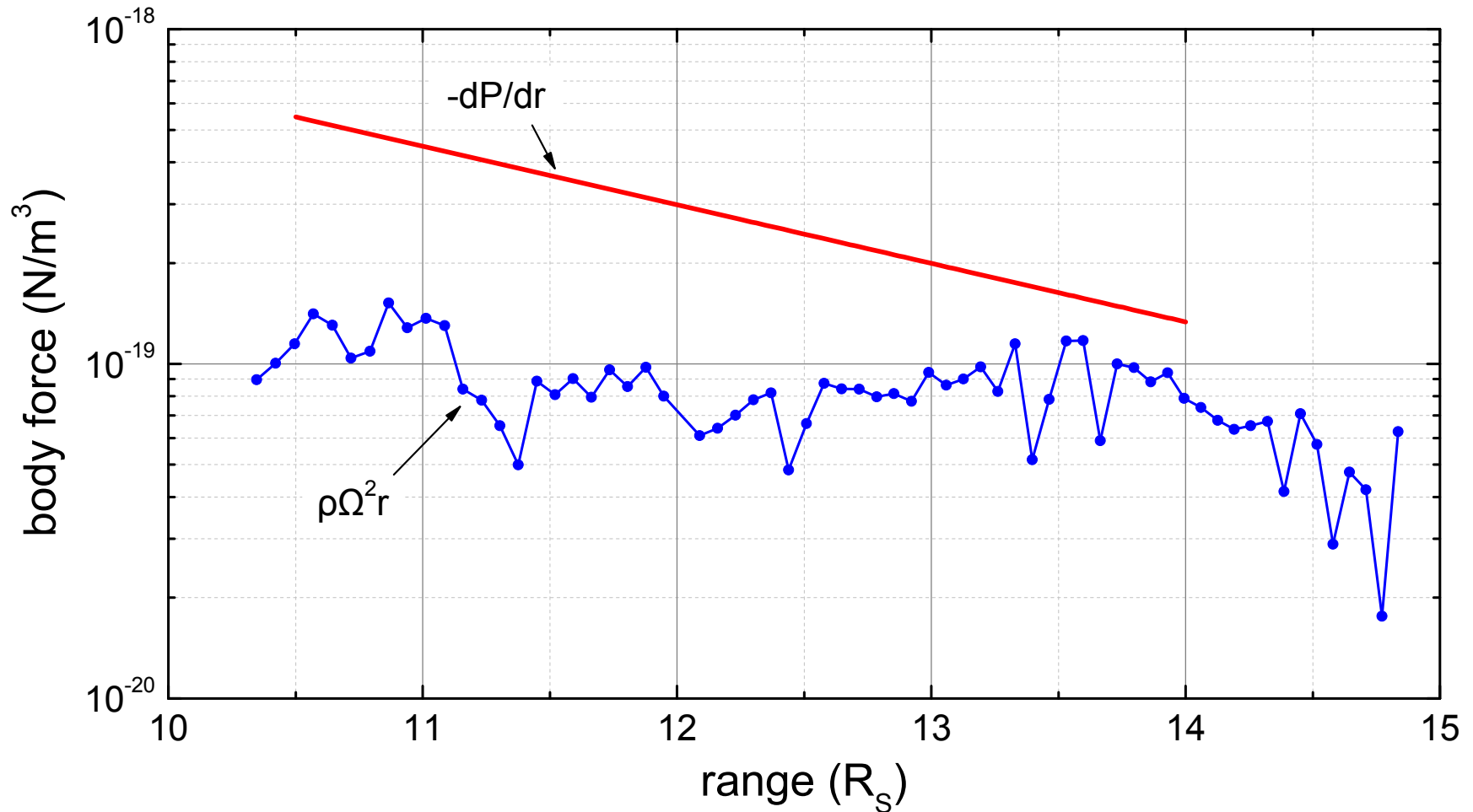
Particle and magnetic radial pressure profiles DOY 322/2007

Plasma: $P=3nkT$ (lower limit), $\langle m \rangle = 12m_p$ (25% protons - 75% O+).



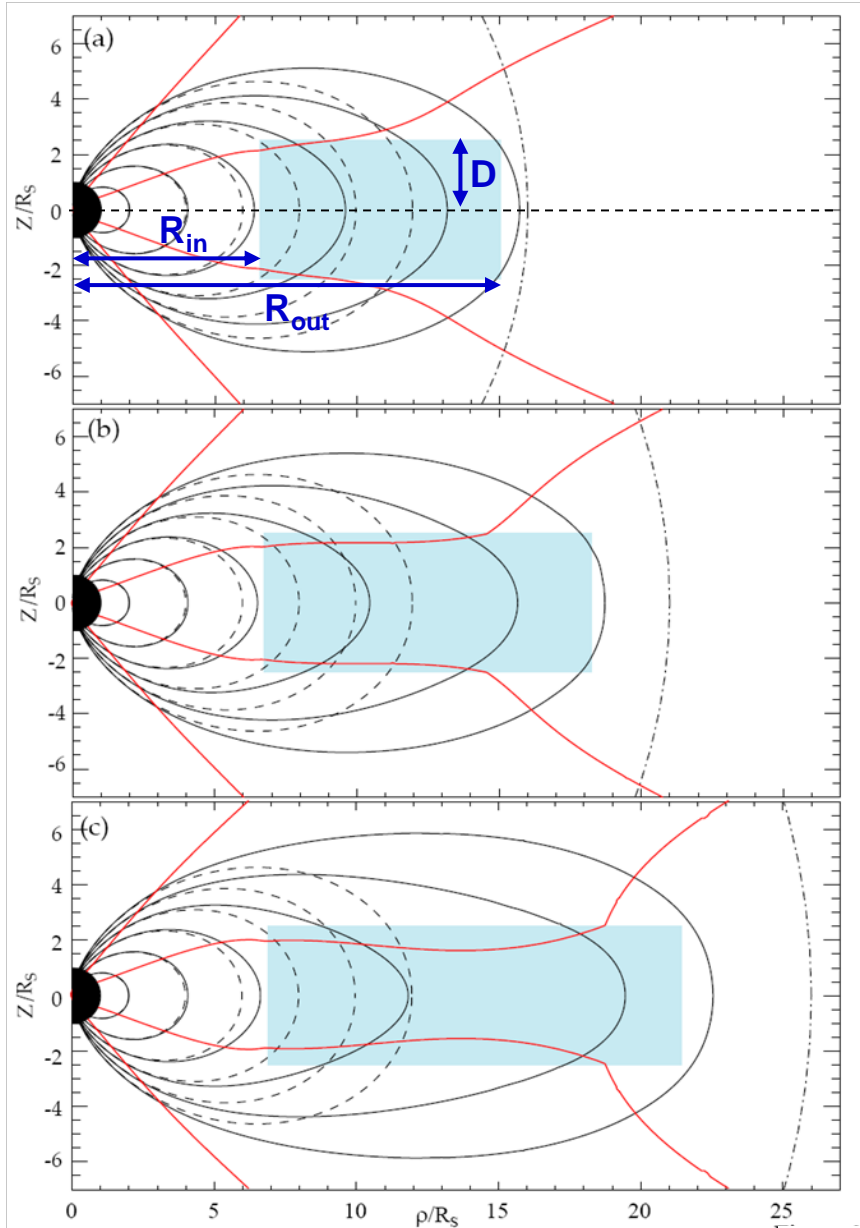
This is a typical example !

$\frac{dP}{dr}$ turned out to be $\times 2$ to $\times 5$ compared to $\rho\Omega^2 r$



Modeling

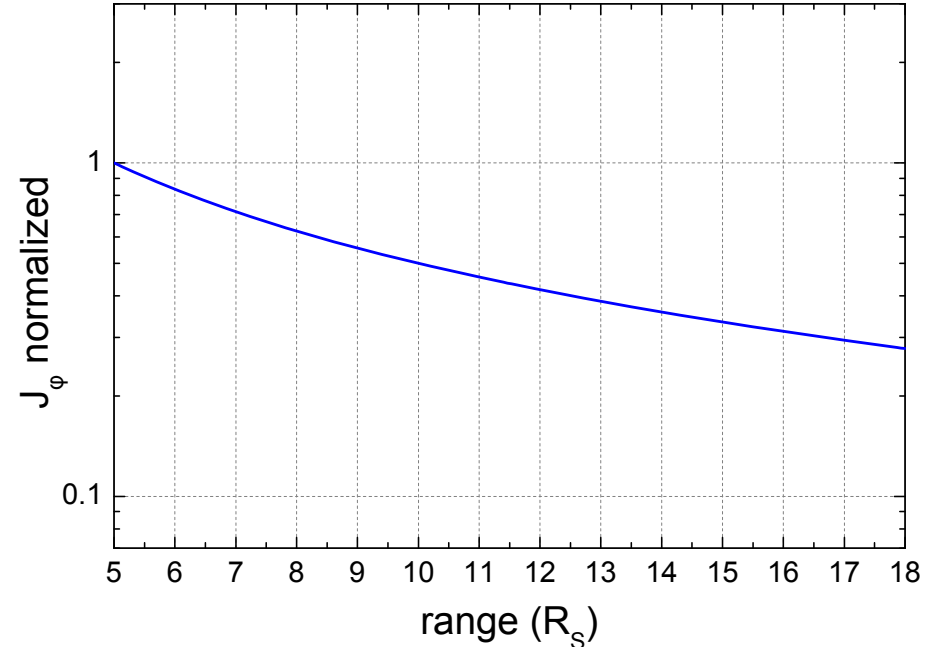
Disk current model, Bunce et al., 2007

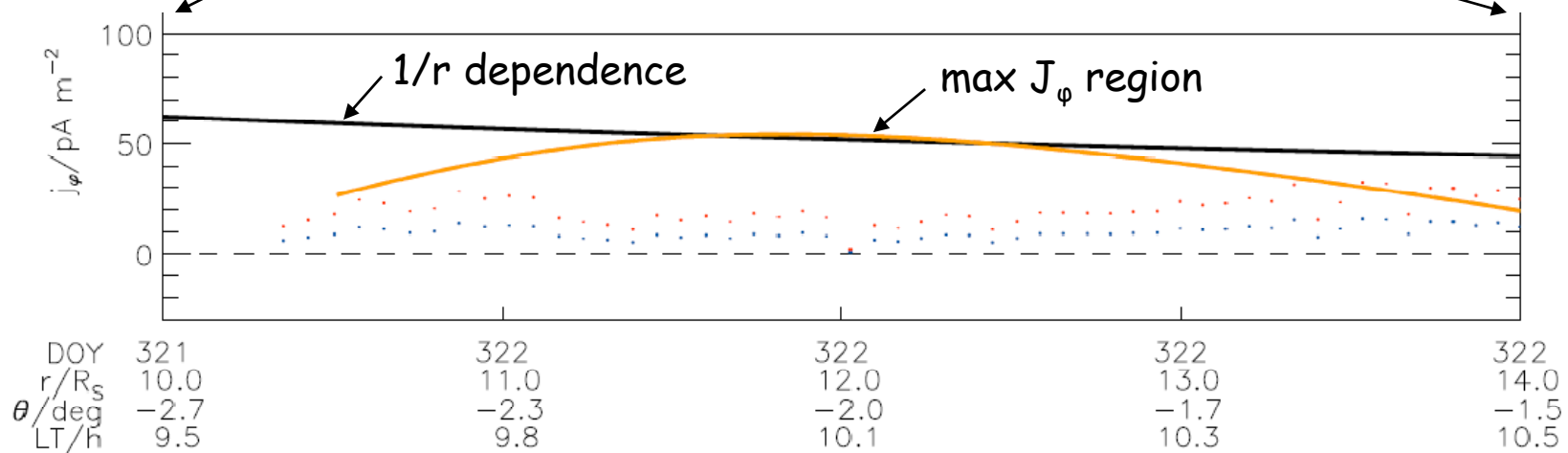
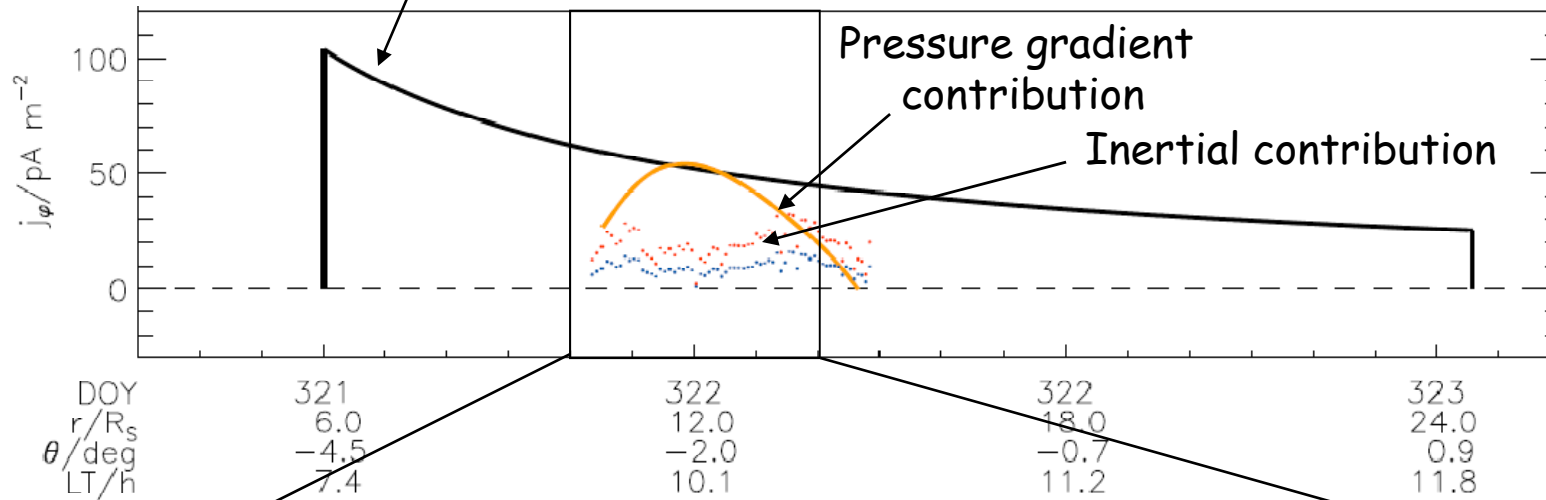


Existing models fit the magnetic field data using 4 free parameters:

- Inner radius R_{in}
- Outer radius R_{out}
- Half thickness D
- Current strength $\mu_0 I_0$

Assuming square cross-section and a $J_\phi \sim 1/r$ dependence.



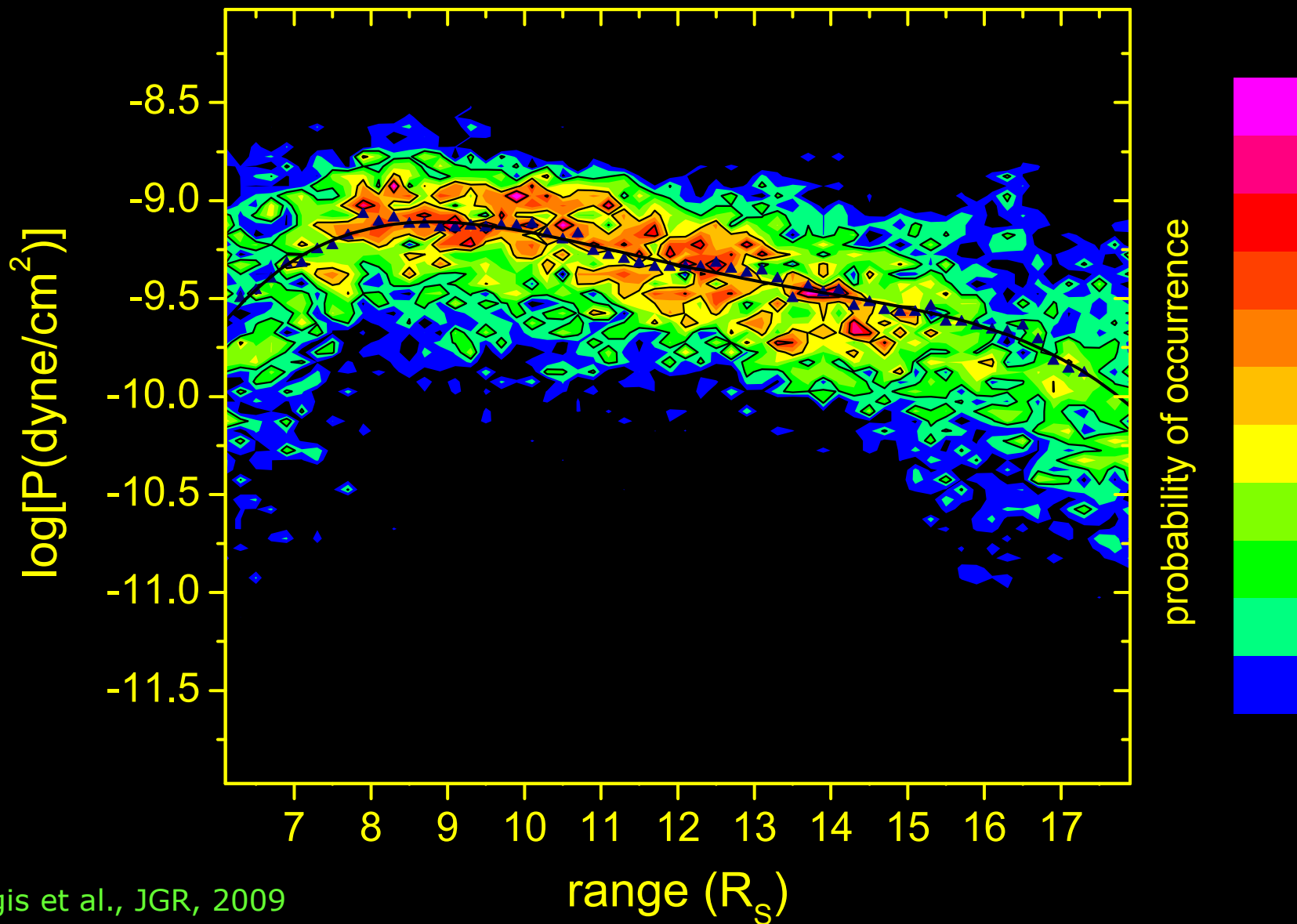
Model J_ϕ 

Kellet et al., 2008

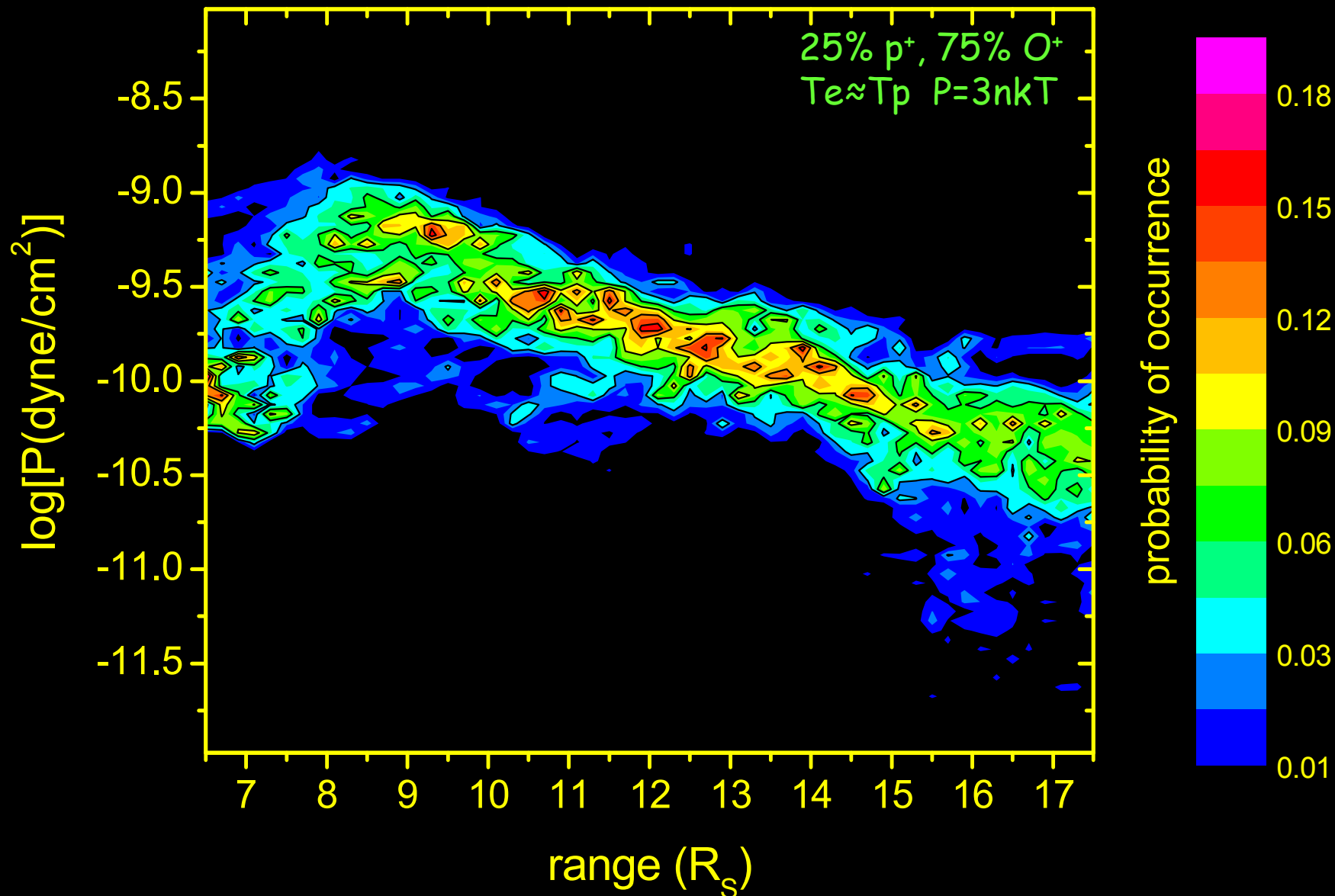
When the pressure gradient dominates over the inertial term, J_ϕ develops a clear maximum and cannot be reproduced by a disk current model.

Statistical approach. What happens more often?

Energetic ($E > 3$ keV) particle pressure in the equatorial plane
All available Cassini data 2004-2008

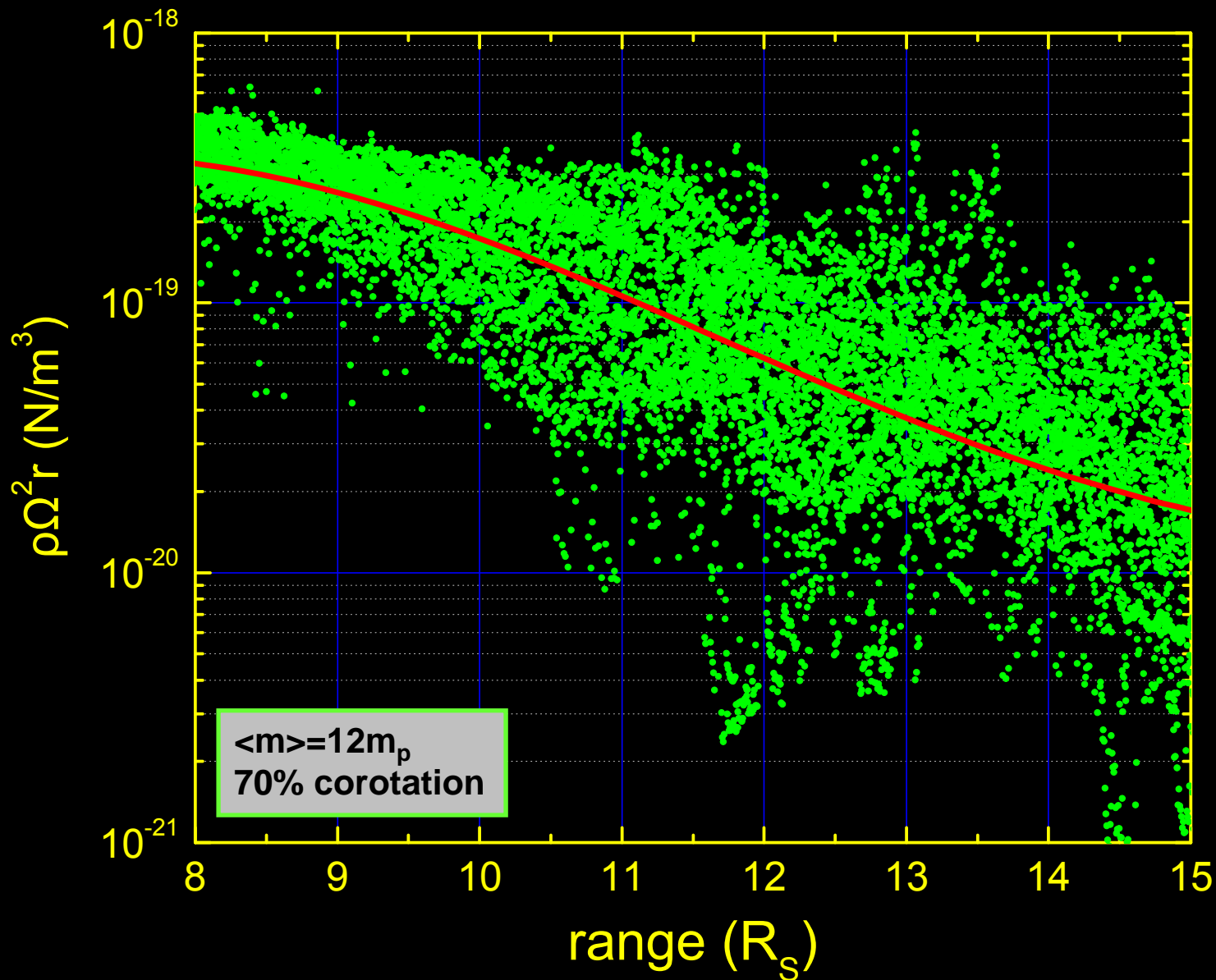


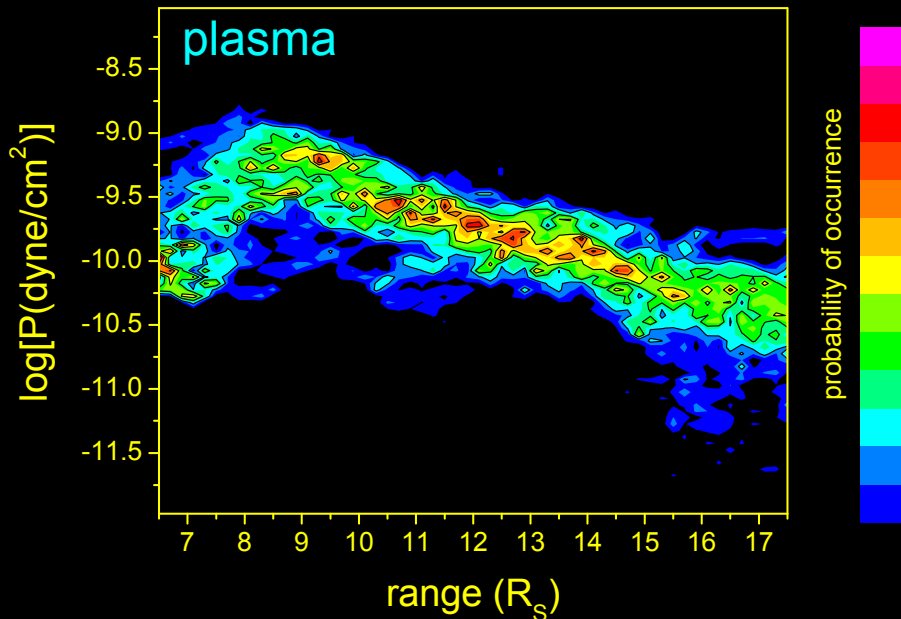
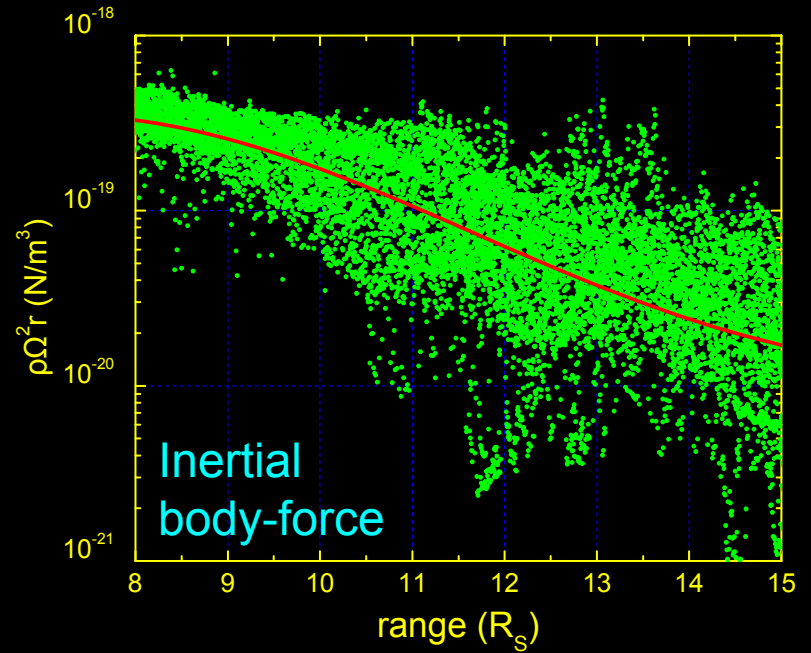
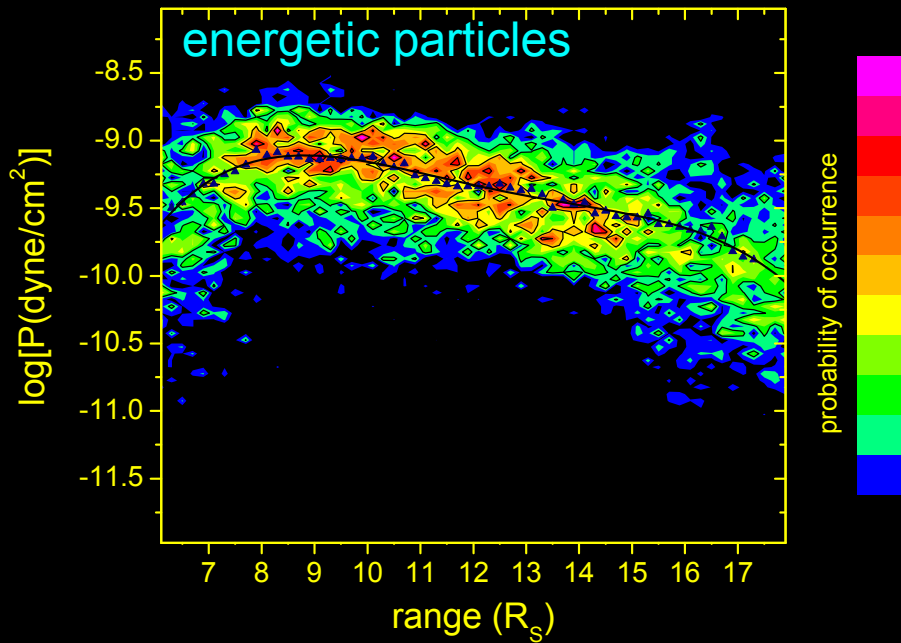
Plasma ($E < 3$ keV) pressure in the equatorial plane
All available Cassini data 2004-2008



Inertial (centrifugal) body force radial profile

All available Cassini data 2004-2008





After 5 years in orbit we have sufficient data to look into the statistical behavior of the system

However, the dynamic nature of the Saturnian magnetosphere appears almost overwhelming!

We can compare the average $-dP/dr$ to the average $\rho\Omega^2 r$ in the context of the radial force balance equation

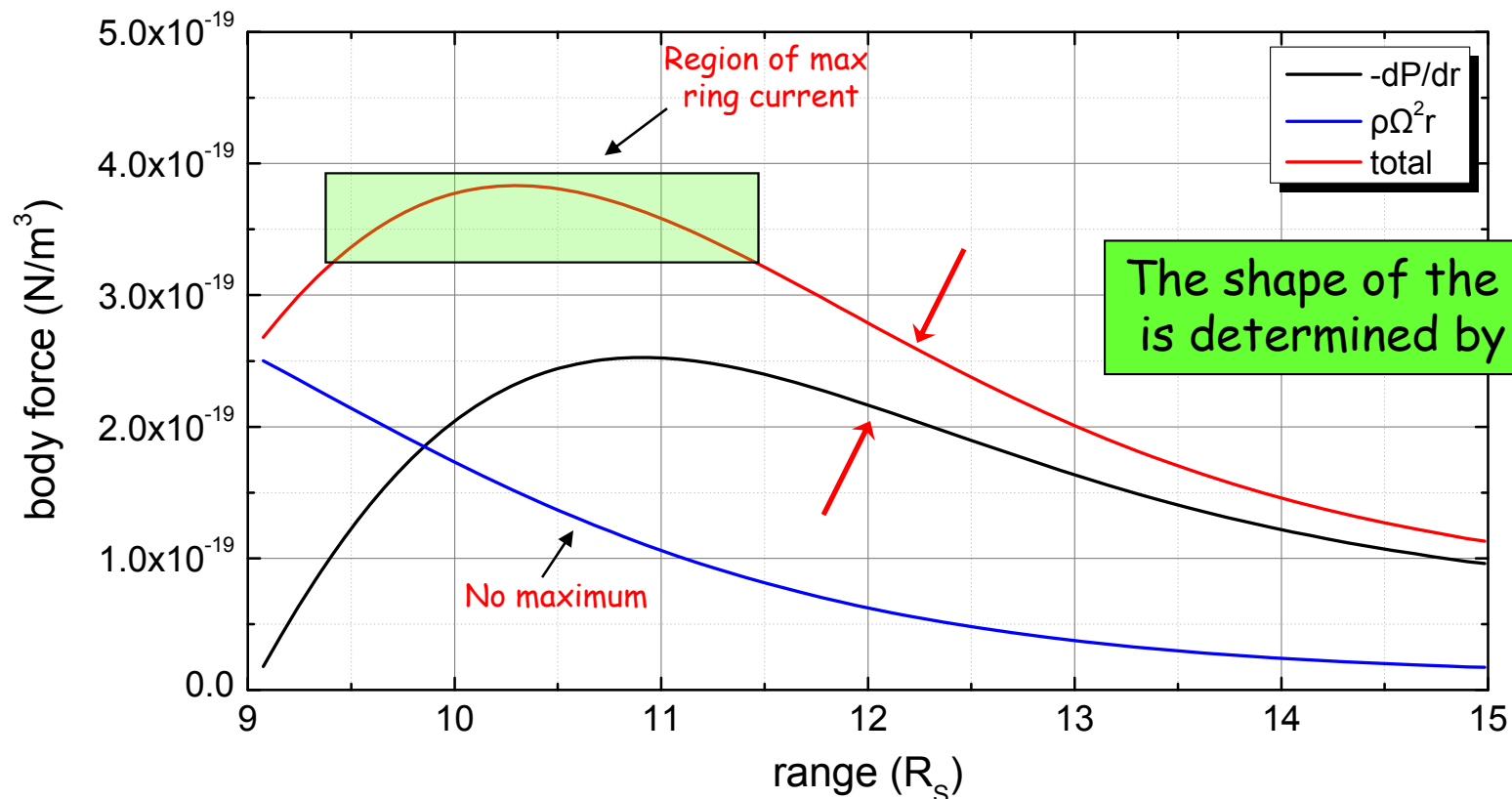
$$\rho\Omega^2 r - \frac{\partial P}{\partial r} \approx J_\phi B_\theta$$

(under the adopted assumptions)

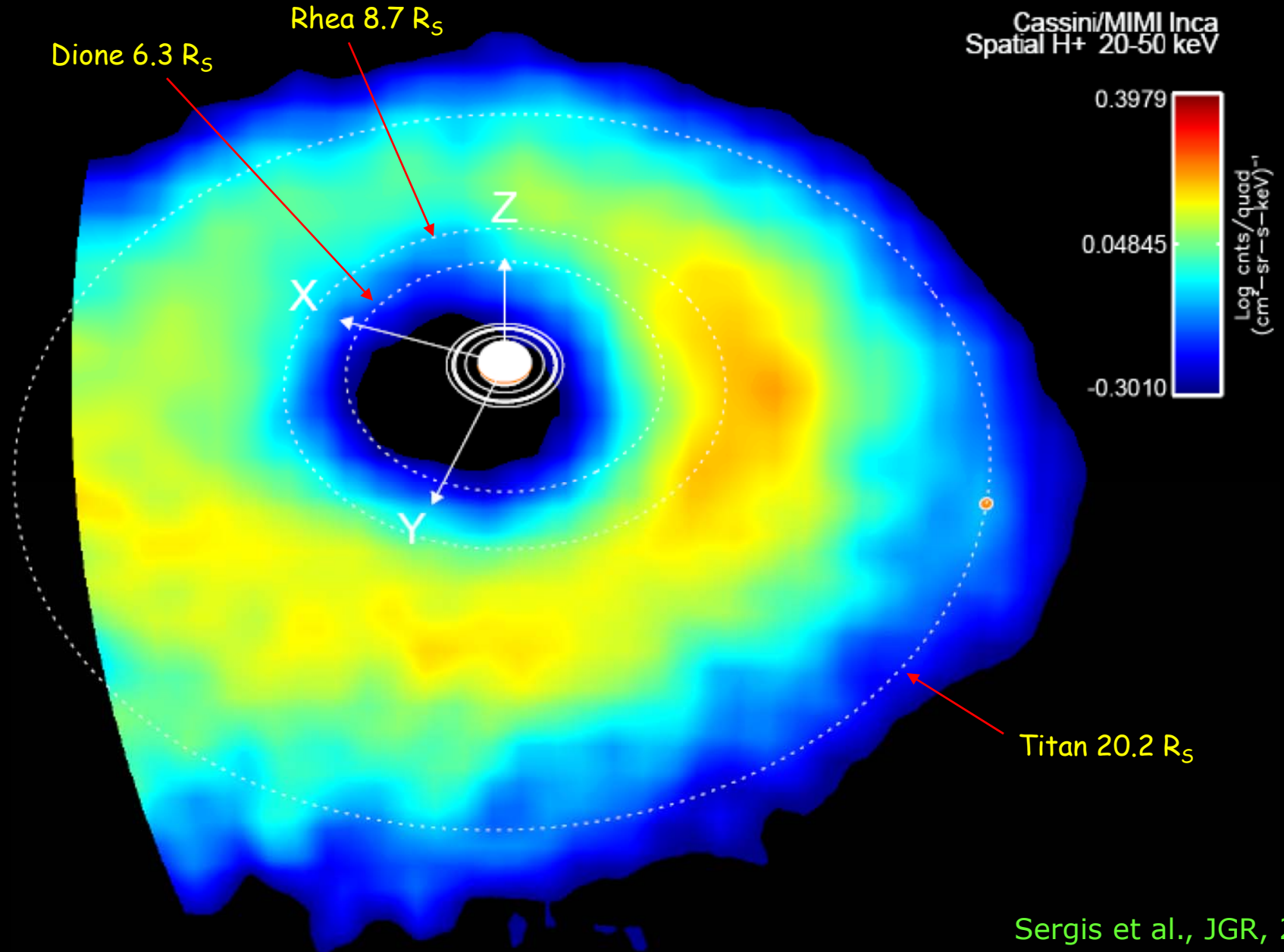
Both terms are of the same order of magnitude and comparable within their uncertainty range.

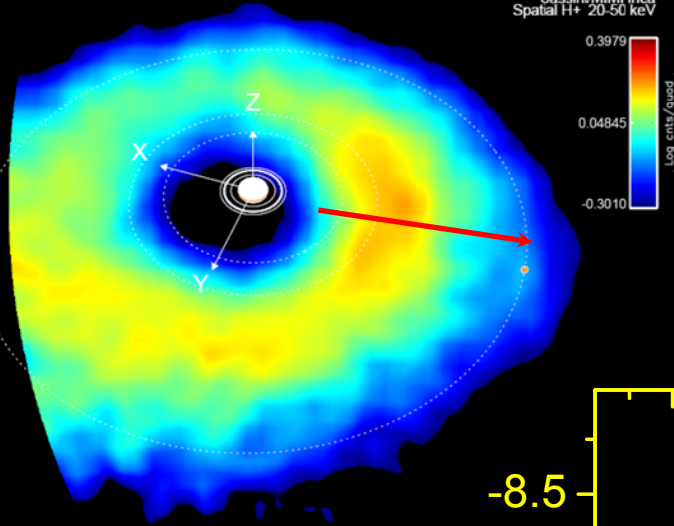
However...

outside 10 R_s , the pressure gradient is almost always higher (x2 or x3) than the inertial term.

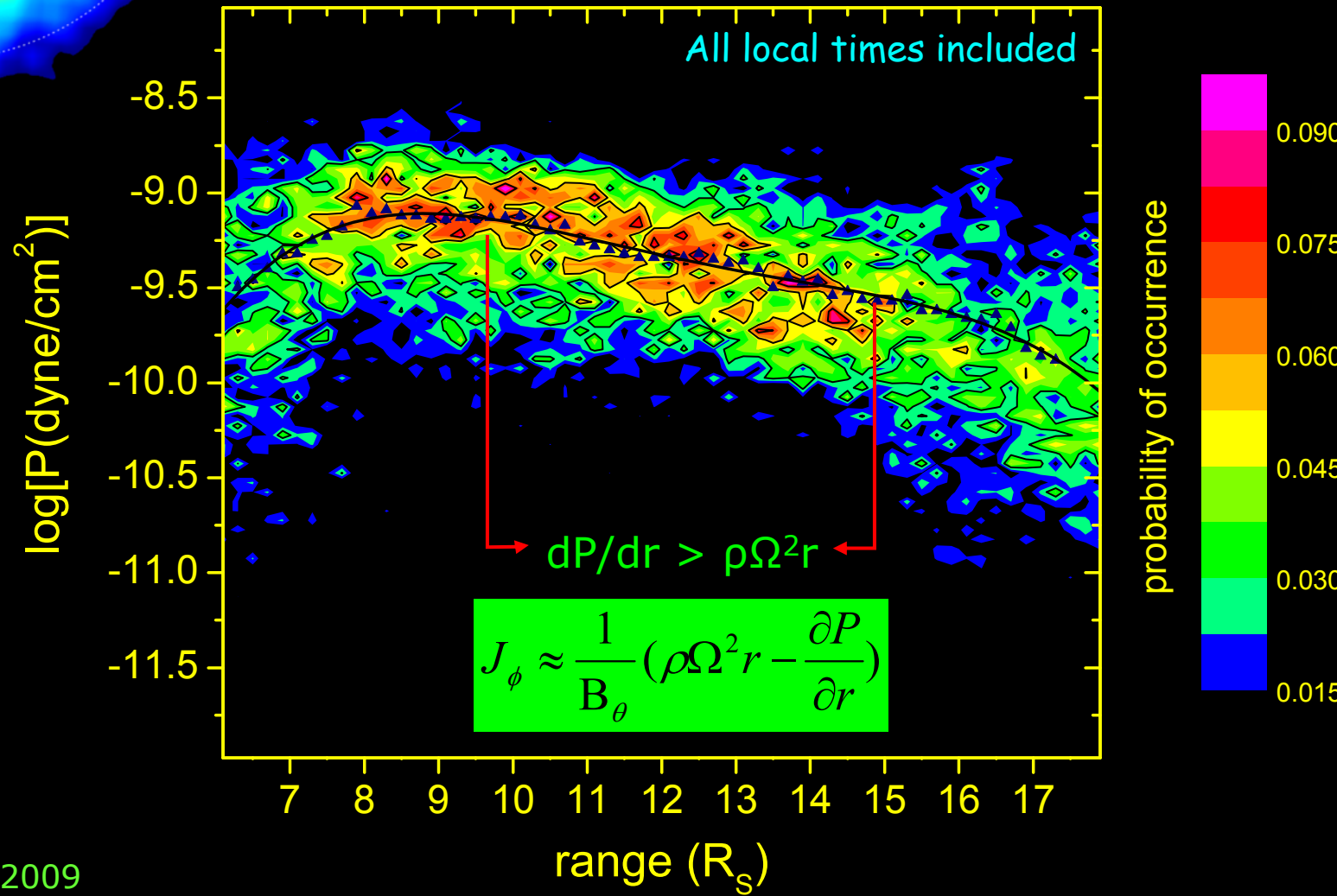


The Saturnian ring current has been captured by MIMI/INCA





To sum up...
A dynamic, O⁺-rich, typically pressure-driven ring current at keV energies in Saturn.



Conclusions

1. Can we obtain a representative radial (total) pressure profile for the dayside magnetosphere of Saturn?

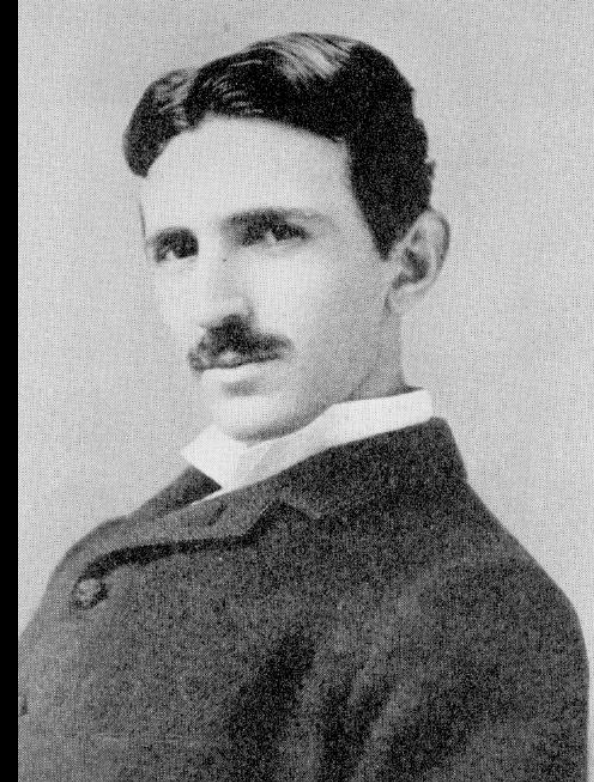
Still **very difficult**. Cassini's orbit, measurement uncertainties and intense dynamics are the basic problems.

2. How are plasma, suprathermal and magnetic pressure compared between 8 and 15 R_S ?
Particle pressure and plasma β are **highly variable** (one order of magnitude)
Half the particle pressure in the keV energy range.
Plasma beta > 1 outside of 9 R_S .

3. How is the radial pressure gradient (dP/dr) compared to the centrifugal body-force ($\rho\Omega^2r$)? Is the Saturnian ring current inertial or pressure-gradient driven?
Outside of 10 R_S , the pressure gradient is almost always higher than the inertial term resulting a **pressure gradient driven ring current**, with a **vital role played by the energetic (keV) particles**.

4. How well can the ring current density be reproduced by existing models?

Disk current models cannot reproduce the ring current when dP/dr is greater than or comparable to $\rho\Omega^2r$, which is usually the case.



“I have been inspired by Tesla. The man thought big, he had revolutionary ideas. He was a risk taker, he had high risk, high payoff ideas. You expect, if you're lucky, to have one percent of these ideas be true, then you've made a tremendous contribution. Tesla had much more than one percent of his ideas being true. **I would be lucky if I had one percent of my ideas being utilized,** even one hundredth of what Tesla has succeeded.”

Dennis Papadopoulos, December 2000.

Thank you

Suprathermal pressure calculation

$$P_i = \frac{8\pi}{3} \int_{E_{i-\min}}^{E_{i-\max}} dE \frac{E_i}{v_i} j_i(E) \Rightarrow P_{part.} = \frac{8\pi}{3} \sum_i \left[\Delta E_i \left(\frac{E_i}{v_i} \right) j_i(E) \right]$$

P_i : pressure supplied by the i -energy channel,

$P_{part.}$: partial particle pressure,

P_{mag} : magnetic pressure,

$E_{i-\min}$, $E_{i-\max}$: lower and upper limits of each i -energy channel of central energy E_i

ΔE_i : channel energy width,

j_i : differential intensity,

v_i : velocity of either of the ion species (H^+ or O^+ in our case) in a particular i -channel.

Plasma pressure calculation

$$\left. \begin{array}{l} T_p \approx T_e \quad T_w \approx \alpha T_e \quad (3 < \alpha < 5) \\ n_p \approx n_e \quad n_w \approx 3n_e \end{array} \right\} \Rightarrow P \approx 4n_e kT_e$$

Assuming ion components of the same bulk velocity (\mathbf{V}) in the steady state, the force (momentum) balance relation can follow from the first (velocity \mathbf{v}) moment of the collisionless (source-free) Vlasov equation:

$$\rho \mathbf{V} \nabla \mathbf{V} + \nabla \cdot \mathbf{P} - \mathbf{J} \times \mathbf{B} = 0$$

For non-relativistic particles ($v \ll c$ for ions and electrons) we can use velocity (rather than momentum) space distribution functions for each component.

$$\mathbf{P} = \int d^3\mathbf{v} (\mathbf{v} - \mathbf{V})(\mathbf{v} - \mathbf{V}) \sum m_i f_i(\mathbf{r}, \mathbf{p}, t) = \int d^3\mathbf{v} (\mathbf{v}\mathbf{v}) \sum m_i f_i(\mathbf{r}, \mathbf{p}, t) - \rho \mathbf{V}\mathbf{V}$$

$$\rho \mathbf{V} = \int d^3\mathbf{v} (\mathbf{v}) \sum m_i f_i(\mathbf{r}, \mathbf{p}, t) \quad \rho = \int d^3\mathbf{v} \sum m_i f_i(\mathbf{r}, \mathbf{p}, t)$$

The radial component of this equation in the equatorial plane expressed in spherical polar coordinates (r, θ, φ) can be written as:

$$-\rho \Omega^2 r + \frac{\partial P}{\partial r} - (J_\theta B_\varphi - J_\varphi B_\theta) = 0$$

We assume that the pressure is isotropic (P) and that the plasma is corotating with a constant angular velocity, so that $\mathbf{V} = \boldsymbol{\Omega} \times \mathbf{r}$. If the plasma bulk velocity does not obey strict corotation, other terms will appear in the radial component of $\mathbf{V} \cdot \nabla \mathbf{V}$.

In Saturn's equatorial plane, J_θ is field-aligned and small, $|\mathbf{B}_\varphi / \mathbf{B}_\theta| \ll 1$ and $\mathbf{B}_\theta \approx \mathbf{B}$, therefore:

$$-\rho\Omega^2 r + \frac{\partial P}{\partial r} + J_\phi B \approx 0$$

At the radius where the pressure is near its maximum ($\partial P/\partial r \cong 0$), the centrifugal body force must become comparable to the radial component of the $\mathbf{J} \times \mathbf{B}$ force. Elsewhere:

$$J_\phi \approx \frac{1}{B} \left(\rho\Omega^2 r - \frac{\partial P}{\partial r} \right)$$

Both $\rho\Omega^2 r$ and $\partial P/\partial r$ must be taken into account. Outside $r \approx 9 R_s$, $\partial P/\partial r < 0$, and both terms add together to contribute to the J_ϕ .