

Simplified simulations of MHD turbulence in a coronal loop

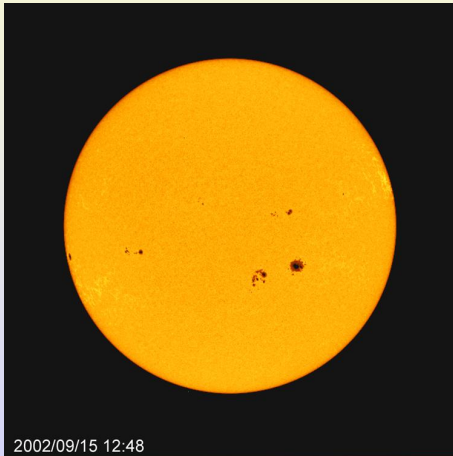
Éric Buchlin




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2 Feb 2005



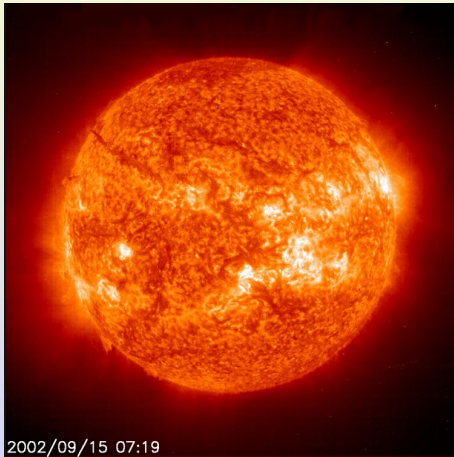
The Solar atmosphere



- ▶  Visible surface (photosphere), 6000 K
- ▶  High chromosphere / transition region, 50 000 K, altitude: 2000 km
- ▶  Corona, $> 10^6$ K, altitude: > 5000 km.



The Solar atmosphere



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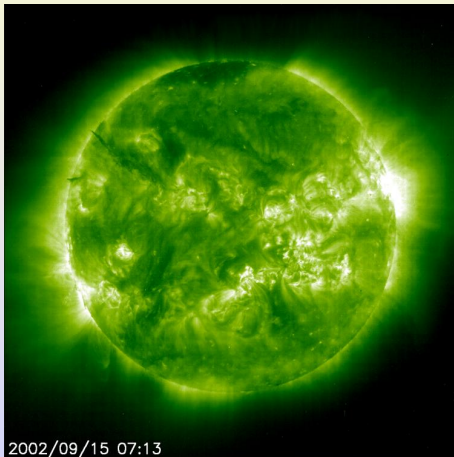
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




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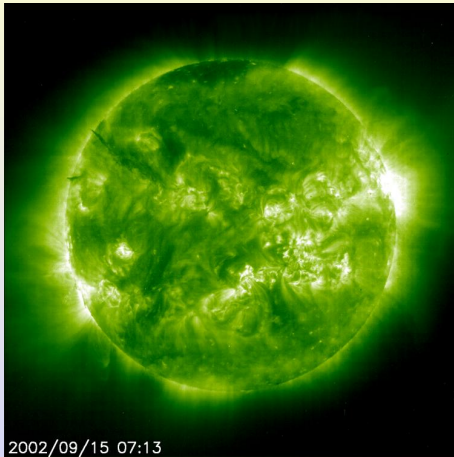
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Coronal structures



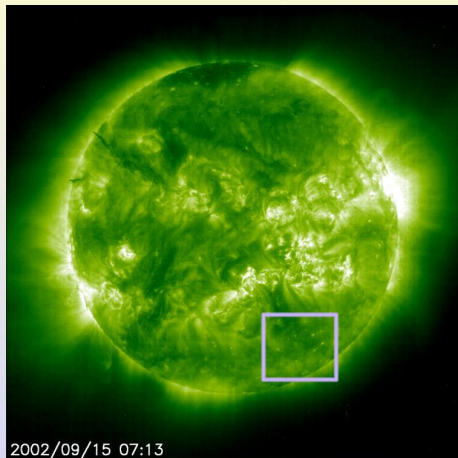
EIT 195, 2002

High energy dissipation \longrightarrow bright structure in UV

- ▶ Active regions
- ▶ Quiet Sun, bright points



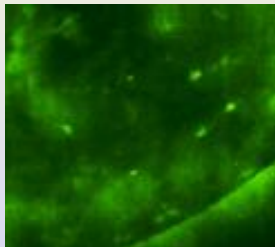
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EIT 195, 2002

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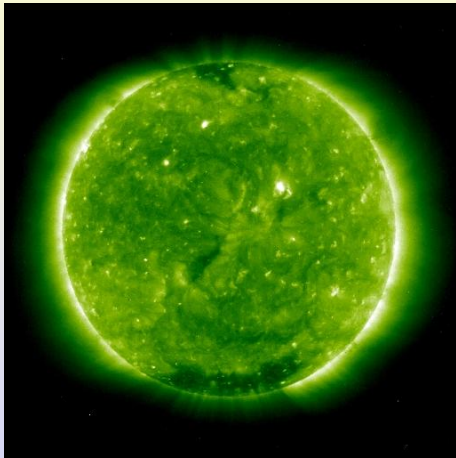
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Structures even smaller may exist



Coronal structures



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EIT 195, 1997

High energy dissipation \longrightarrow bright structure in UV



The coronal heating problem

- ▶ The corona is very hot (known since 1943)
- ▶ The Sun produces enough energy (in its core) to heat it

But:

- ▶ This energy needs to be transported to the corona, in a non-thermal way. Solutions:
 - Sound waves and slow magnetosonic waves excluded (do not reach the corona)
 - Fast magnetosonic and Alfvén waves
- ▶ It needs then to be dissipated, but:
 - The observed events of energy dissipation are not enough
 - The physical mechanisms of dissipation:
 - wave-particle interactions
 - reconnexion, resistivity (Joule)are too slow (not enough efficient)



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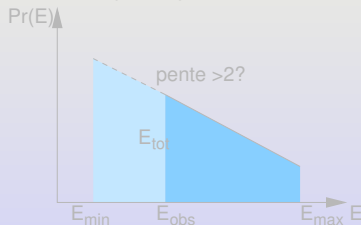
How to (finally) solve the problem?

A possible way: small scales (spatial and temporal)

- ▶ Dissipation mechanisms are more efficient at **small** scales (nanoflares, Parker 1988)
- ▶ The **smallest** events (non-observables) contribute perhaps the most to the heating of the corona

- ▶ Diffusion time: L^2/η , with $\eta \approx 1 \text{ m}^2/\text{s}$

- ▶ Hudson (1991):



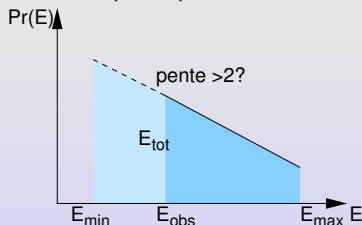
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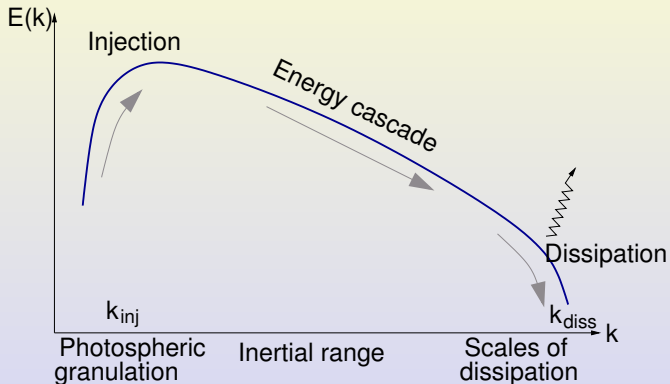
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The origin of small scales

May be created by **turbulence**:



In the corona: $R_e = UL/\eta \approx 10^{13} \gg 1$
 (for $U = 1 \text{ Mm/s}$, $L = 10 \text{ Mm}$, $\eta = 1 \text{ m}^2/\text{s}$)

Smallest (dissipative) scales: 10 m!



Need for statistics of coronal heating

- ▶ Small flares, small scales:
 - cannot be observed directly
 - Hudson... → small-scale statistics extrapolated from observations
- ▶ Theoretical description of turbulence: statistics are a privileged means of tackling its **complexity**

When used at the same time to analyze **observations** and **simulations**:
a **comparison** is possible



Outline

- 1 Introduction
- 2 Numerical models of a coronal loop
- 3 Coupled shell-models
- 4 Conclusions



Numerical models of a coronal loop

- 1 Introduction
- 2 Numerical models of a coronal loop
 - Magnetohydrodynamics
 - Simplifications of MHD
 - Models of coronal loops
 - First model: based on cellular automata
- 3 Coupled shell-models
- 4 Conclusions



MagnetoHydroDynamics (MHD)

Starting-point: incompressible MHD

Equations, with $\rho_0 = 1$, $\nabla \cdot \mathbf{v} = 0$ and $\nabla \cdot \mathbf{B} = 0$:

$$\partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} + \nu \nabla^2 \mathbf{v} \quad (1)$$

$$\partial_t \mathbf{B} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B} \quad (2)$$

Includes velocity \mathbf{v} , magnetic field \mathbf{B} , and:

- ▶ **Non-linear** terms, which allow the creation of small scales by turbulence
- ▶ **Diffusion** terms, which allow energy dissipation (mainly at small scales)
- ▶ Alfvén waves (as fluctuations)



Need for simplifications

Direct numerical simulations (DNS) of MHD:

- ▶ Small grid sizes ($\approx 1024^3$ max)
- ▶ Reynolds numbers too small (1000)
- ▶ Too slow (long computations for just a few events)

⇒ Need for simplified simulations of MHD:

- ▶ Reduce number of spatial dimensions
→ Einaudi, Velli, Georgoulis; Galtier...
- ▶ Reduce the number of active modes in turbulence, and the complexity of their interactions
→ shell-models: Carbone, Giuliani...
→ cellular automata: Lu & Hamilton; Isliker & Vlahos, Krasnoselskikh...



Coronal loops

Our models represent a **coronal loop**: loop of magnetic field, containing plasma.

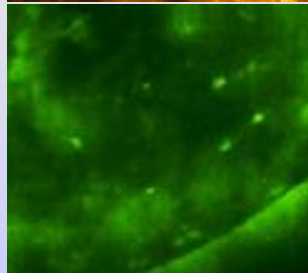
Large loops in active regions:

↑
150 000 km

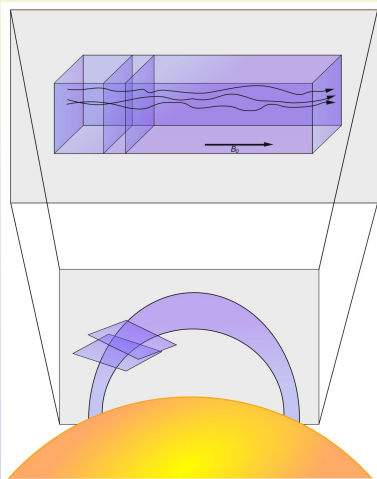


We are mainly interested in **small loops** of the quiet Sun (**bright points** and smaller):

↑
120 000 km



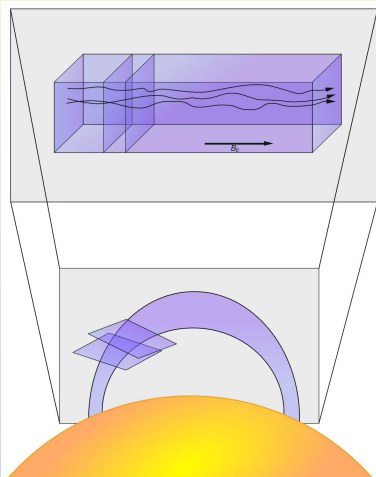
Common framework to our loop models



- ▶ The box represents a loop, MHD
- ▶ Weak forcing at the “photosphere”,
 $\Pi = -(\mathbf{v}_{\perp, \text{ph}} \cdot \mathbf{B}_{\perp}) \mathbf{B}_0 / \mu_0$
- ▶ Propagation of Alfvén waves along \mathbf{B}_0
 (along the loop)
- ▶ Non-linear interactions between these waves
- ▶ Energy dissipation



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Simplification: simplification of the non-linear interactions in each cross-section of the loop:

- 1 Cellular automata
- 2 Shell-models

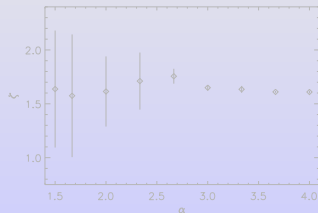
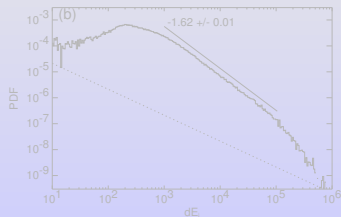
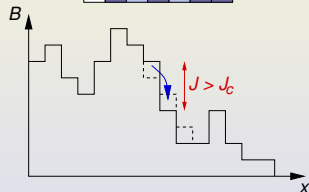
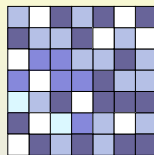


Coupled cellular automata

Each cross-section is a
cellular automaton.

Non-linear interaction between Alfvén waves:
avalanches, with threshold J_c on current density
 J .

→ Buchlin et al. A&A 2003

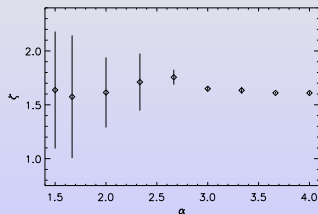
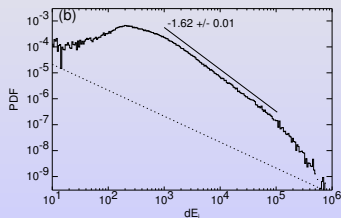
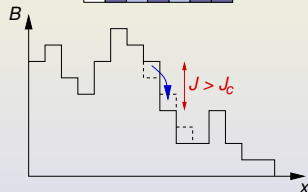
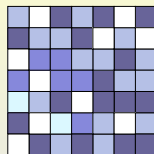


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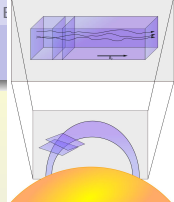


Coupled shell-models

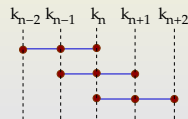
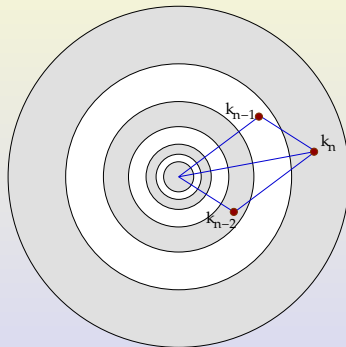
- 1 Introduction
- 2 Numerical models of a coronal loop
- 3 Coupled shell-models**
 - The model
 - Results and statistics
 - Statistics of events
- 4 Conclusions



Coupled shell-models



Each cross-section $\perp \mathbf{B}_0$ is a **shell-model** (Giuliani and Carbone, 1998):



- ▶ Logarithmic spacing of shells (modes) (2D Fourier space)
- ▶ Non-linear interactions between neighboring modes (triads)



Advantages of shell-models

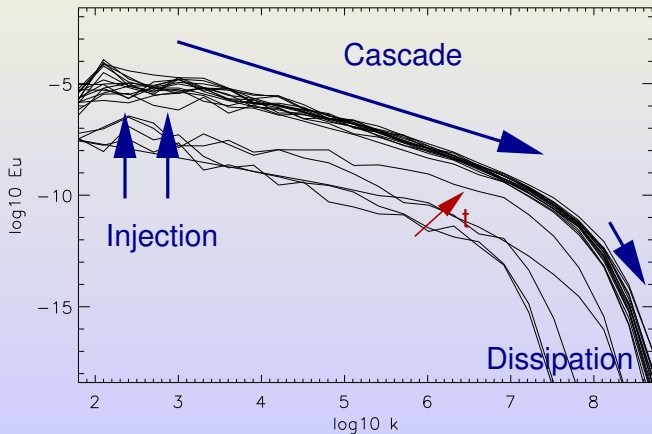
- ▶ Large spread of wavenumbers (of scales) with just a few modes (24)
 - large Reynolds numbers (10^6 instead of ≈ 1000 for DNS)
 - **intermittency** is possible
- ▶ Good model of local **non-linear interactions** between modes of MHD
- ▶ No free parameters (coefficients determined by conservation of **2D MHD invariants**)



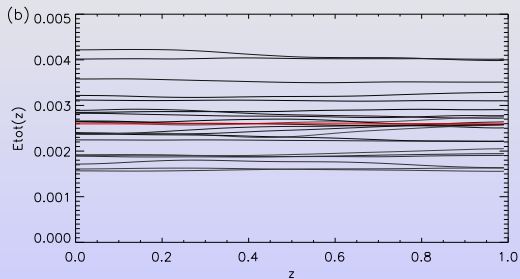
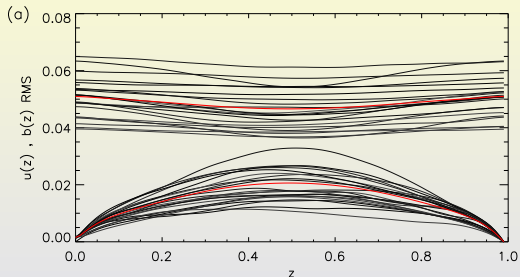
Spectrum development

Small scales, created by turbulent cascade
(due to non-linear terms of MHD)

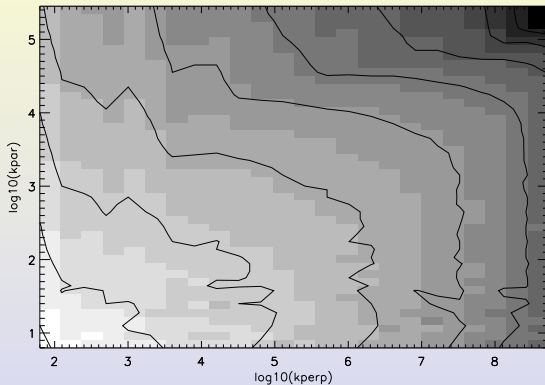
Spectra of \perp kinetic energy in a loop cross-section as a function of time:



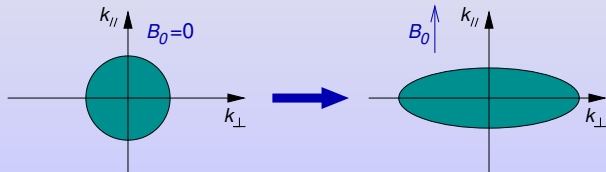
Profile of kinetic and magnetic energy along the loop



2D spectra



Obtained by Fourier transform along the loop.

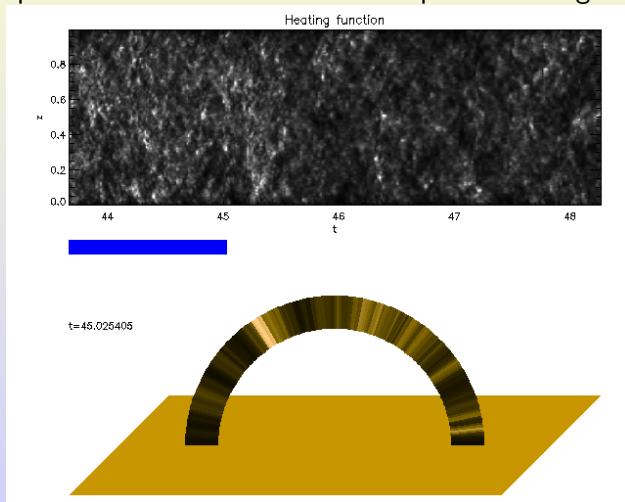


Anisotropic?

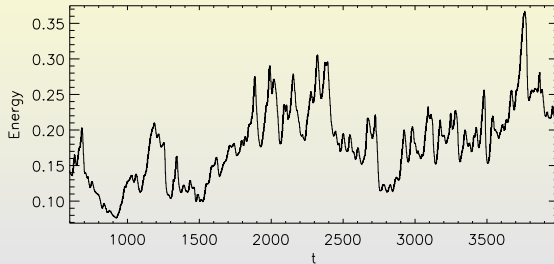


Heating function

Energy dissipation as a function of time and position along the loop:

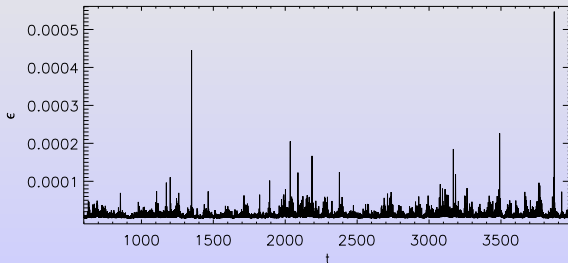


Time series: energy and dissipation



Physical values of model units:

- ▶ Time: 10 s
- ▶ Energy: 10^{17} J
- ▶ Power: 10^{16} W

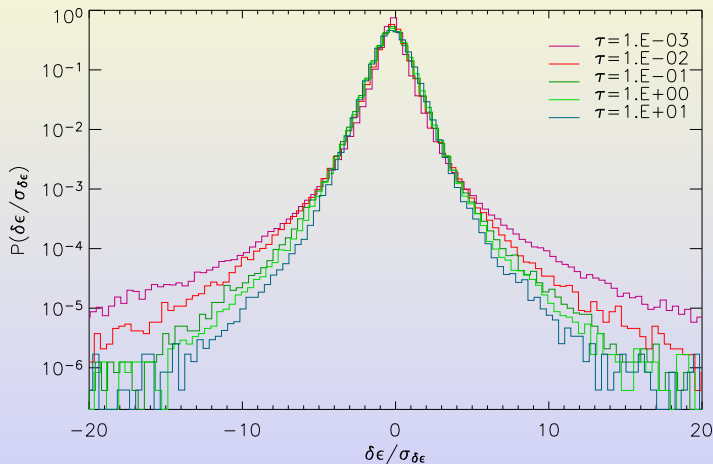


4 days of computation
 $5 \cdot 10^7$ time steps
 100 planes
 (cross-sections)



Intermittency of the dissipation power time series

Distributions of $\delta_\tau \epsilon \equiv \epsilon(t + \tau) - \epsilon(t)$ for different τ 's:



(normalized by their standard deviations)



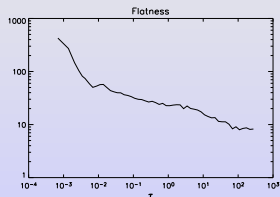
Intermittency of the dissipation power time series (2)

$$\delta_\tau \epsilon \equiv \epsilon(t + \tau) - \epsilon(t)$$

Intermittency:

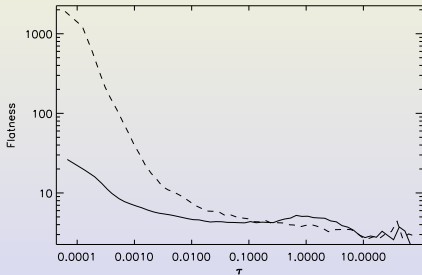
- ▶ Distributions of $\delta_\tau \epsilon$ have a shape which depends on the scale τ ,
Deviation from the Kolmogorov 1941 turbulence theory:
Turbulence is not self-similar (fractal), but **multi-fractal**
- ▶ Exponents ζ_q of the **structure functions** $S^q(\tau) \equiv \langle |\delta_\tau \epsilon|^q \rangle \propto \tau^{\zeta_q}$ are a non-linear function of q .

- ▶ In particular, the **flatness**
 $F(\tau) \equiv S^4(\tau)/(S^2(\tau))^2$
grows at small scales:

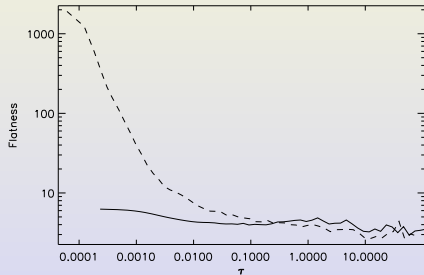


Intermittency as a function of the parameters

Reference run (---): $\nu = 10^{-13}$, $a = L/\ell = 10$



$\nu = 10^{-11}$

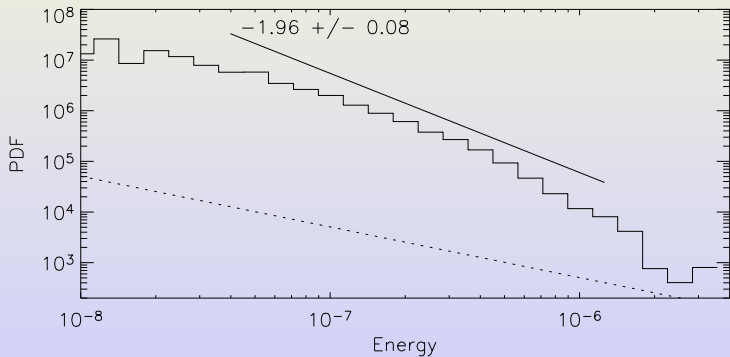
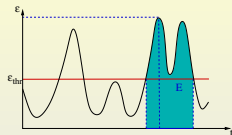


$a = 80$



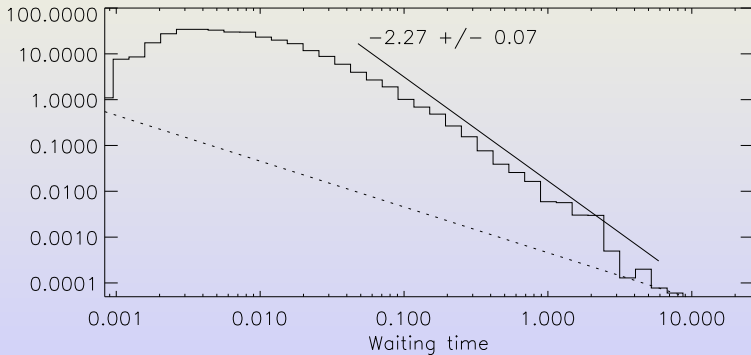
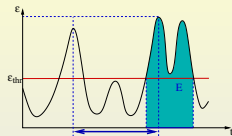
Distributions of events: energies

Distribution of energies of events
(defined by a threshold):



Distributions of events: waiting-times

Distribution of waiting times between events
(defined by a threshold):



Non-Poissonian, long-duration correlations

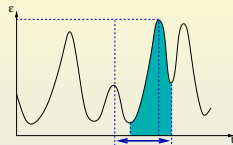
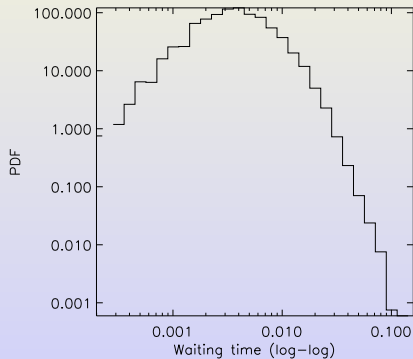


About the definition of events

Other definition:

Each peak can be considered as an event.

Distribution of waiting-times:

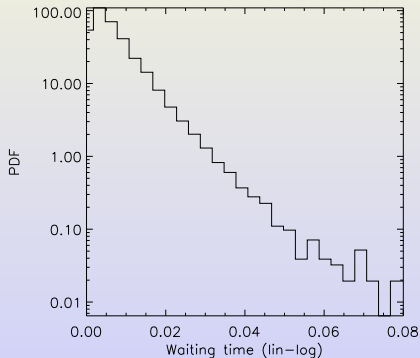
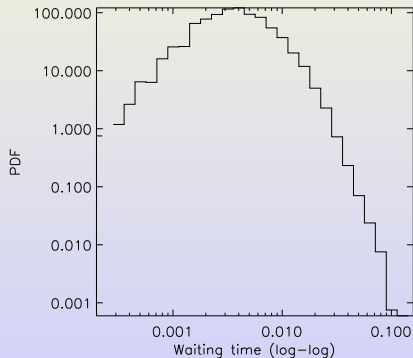
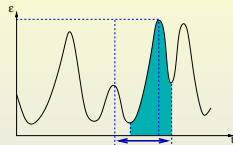


About the definition of events

Other definition:

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Distribution of waiting-times:



Still Poissonian?



About the definition of an event (2)

In the general case:

- ▶ Events statistics do depend on the definition of events (of course...)
Thus conclusions about the Poisson nature of the flaring process, or about the nano-flares hypothesis, **depend on the definition**
- ▶ **Higher sensibility** of statistics to the definition
when intermittency is low
- ▶ Are there really clear events in a time series / structures in a MHD field?
If not, it is better to use **statistics which do not need events to be defined** (structure functions, spectra...)

→ Buchlin et al. 2005 (submitted to A&A)



Conclusions

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- 2 Numerical models of a coronal loop
- 3 Coupled shell-models
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Conclusions

- ▶ Models designed to explore coronal heating at small scales created by MHD turbulence.
- ▶ Simplified so as to be able to produce statistics
- ▶ First models of this type with geometry of a loop and energy loading at footpoints.

Some results of the shell-model:

- ▶ Spectra of turbulence, large Reynolds numbers
- ▶ Intermittency (also as a function of parameters of the model)
- ▶ Heating function: dissipation power as a function of time and position
- ▶ Events statistics distributed as power-laws: -2 for energy, -2.3 for waiting-times.

But these distributions depend on the definition of an event



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Current and future developments

To allow a better comparison between model output and observations:

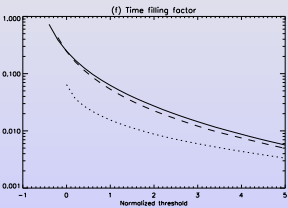
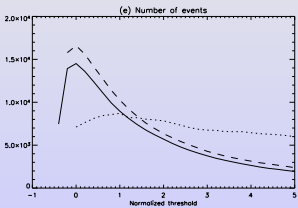
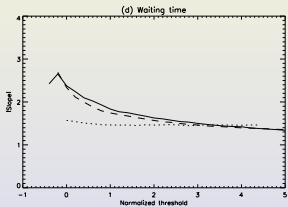
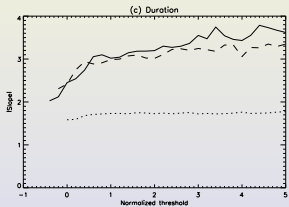
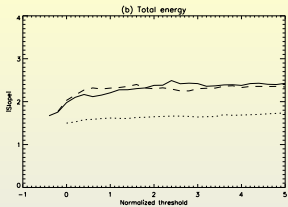
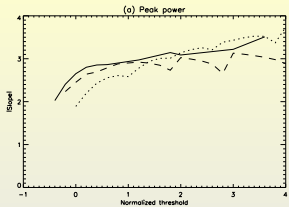
- ▶ Computation of radiative output
(via thermodynamics and atomic models)
- ▶ More realistic:
 - Stratification: $B_0(z)$...
 - Parameters and energy input deduced from observations
- ▶ Produce images of the luminosity of the loop
(geometry from extrapolations)

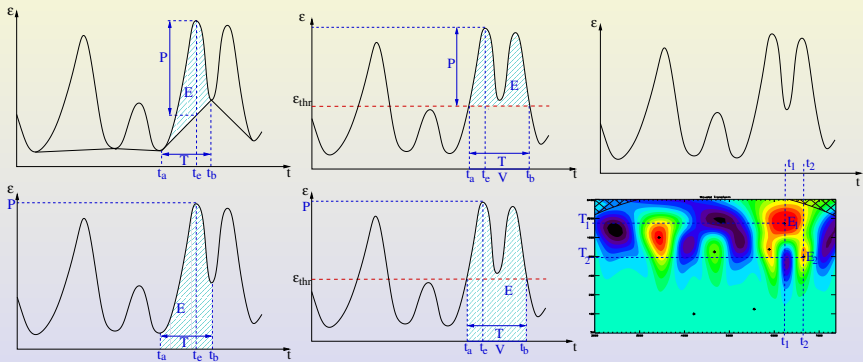


Appendix

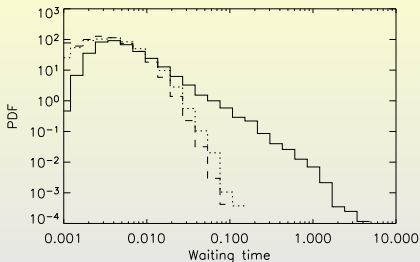
- ▶ Variation of the slope of the distributions as a function of the threshold
- ▶ Definitions of events
- ▶ Waiting-time statistics, definitions, and intermittency



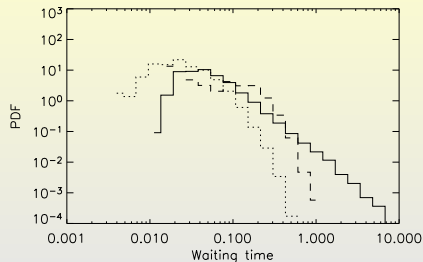




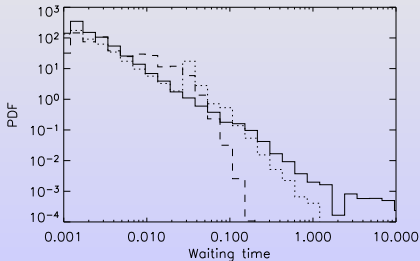
Time series < 1 >:



Time series < 2 >:



Time series < 3 >:



—: threshold

- - -: peaks

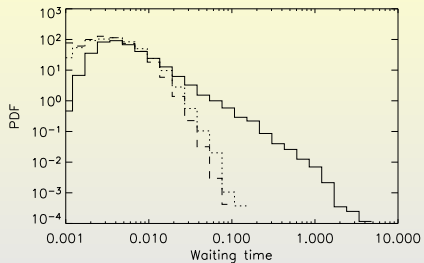
. . . : wavelets

More intermittency

→ less sensitivity



Time series $\langle 1 \rangle$:



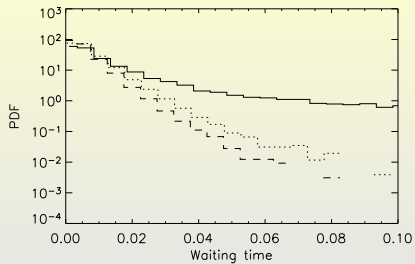
—: threshold

- - -: peaks

. . .: wavelets



Time series $\langle 1 \rangle$:



—: threshold

- - -: peaks

. . .: wavelets



