Nonlinear $r$-Modes in Rapidly Rotating Relativistic Stars

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The $r$-mode instability in rotating relativistic stars has been shown recently to have important astrophysical implications, provided that $r$-modes are not saturated at low amplitudes by nonlinear effects or by dissipative mechanisms. Here, we present the first study of nonlinear $r$-modes in isentropic, rapidly rotating relativistic stars, via 3D general-relativistic hydrodynamical evolutions. We find that (1) on dynamical time scales, there is no strong nonlinear coupling of $r$-modes to other modes at amplitudes of order one—the maximum $r$-mode amplitude is of order unity. (2) $r$-modes and inertial modes in isentropic stars are predominantly discrete modes. (3) The kinematical drift associated with $r$-modes appears to be present in our simulations, but confirmation requires more precise initial data.

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Considerable effort has been spent in the last two years in determining the properties of $r$-modes [1] in rotating compact stars, since the discovery [2,3] that these modes are unstable to the emission of gravitational radiation. This is motivated by the current understanding that the $r$-mode instability may have several important astrophysical consequences: it provides an explanation for the spin-down of accreting neutron stars, while being a strong source of de-...
configuration (apart from the influence of local truncation errors) in long-term evolutions. The rotational velocity profile ($v^i$ along the $y$ axis, where $v^i$ are the contravariant components of the equilibrium 3-velocity, as measured by a local zero-angular-momentum observer) remains nearly unchanged over more than 20 rotational periods (all simulations presented here are stopped due to computing time restrictions only and could be continued for longer times).

In order to excite $r$-modes during the time evolution, we perturb the initial stationary model by adding a specific perturbation $\delta v^i$ to the equilibrium 3-velocity $v^i$. As there is no exact linear eigenfunction available in the literature that would correspond to an $l = m = 2$ $r$-mode eigenfunction for rapidly rotating relativistic stars in the Cowling approximation, we are using an approximate eigenfunction, derived in spherical polar coordinates, valid in the slow-rotation $O(\Omega)$ limit to the first post-Newtonian order (see [15,16] for details). We apply the usual transformation between the spherical polar and Cartesian coordinate systems as the latter is used in the 3D code. We note that, in the Newtonian limit, an amplitude of $\alpha \approx 1.0$ corresponds to a maximum velocity perturbation (i.e., the maximum value of the velocity component $\delta v^\theta$ in the equatorial plane) equal to roughly 1/3 of the rotational velocity of the star at the equator. With this eigenfunction (multiplied by a chosen initial amplitude $\alpha$, which coincides, in the Newtonian limit, with the definition of the amplitude in [17]), we are able to excite mainly the $l = m = 2$ $r$-mode. Because of the approximate nature of the eigenfunction, additional modes are also excited, primarily $m = 2$ inertial modes.

Figure 2 displays the evolution of the axial velocity in the equatorial plane ($v^z$ along the $y$ axis) at a coordinate radius of $r = 0.75r_e$, where $r_e$ is the coordinate radius at the equator. It is clear that the evolution is a superposition of several modes and that one mode, the $l = m = 2$ $r$-mode, is the dominant component. The amplitude in this evolution is $\alpha \approx 1.0$. The perturbed star is evolved for more than 25 ms (26 $l = m = 2$ $r$-mode periods), during which the amplitude of the oscillation decreases due to numerical viscosity. Even at amplitudes larger than 1.0, the evolution is still similar to that in Fig. 2, with no sign of nonlinear saturation of the $r$-mode amplitude on a dynamical time scale (only when $\alpha$ exceeds a value much larger than 1.0, nonlinear hydrodynamical saturation sets in; however, the precise determination of the maximum saturation amplitude will require more accurate initial data and the simultaneous dynamical evolution of the gravitational field). Our result implies that, unless nonlinear saturation is set at time scales much longer than the dynamical one, nonlinear hydrodynamical couplings would not prohibit gravitational radiation reaction to drive unstable $r$-modes to large amplitudes (of order one) before saturation sets in. An absolute upper limit on the $r$-mode amplitude (relevant, perhaps, for the most rapidly rotating stars) is set by causality, requiring $\sqrt{v_i v^i} < c$ (where $c$ is the speed of light), or approximately, $\alpha_{\text{causal}} \approx 3c/\Omega R$, where $\Omega$ is the angular velocity of the equilibrium star (for slowly rotating stars, we expect other upper limits to be more relevant).

A Fourier transform of the time evolution shown in Fig. 2, as a function of the frequency in the inertial frame, shown in Fig. 3, reveals that the initial data we are using excite mainly the $l = m = 2$ $r$-mode ($r_2$) with a frequency of 1.03 kHz and, with smaller amplitudes, several inertial modes ($i_3$, $i_4$, $i_5$) and higher harmonics. The frequencies of the various modes are the same at any given point inside the star (the specific data shown in the figure correspond to radii $r = 0.5r_e$ and $r = 0.75r_e$); i.e., the evolution consists of a sum of discrete modes and possible associated continuous parts were not excited. This is in agreement with the de
with the conclusions in [16], regarding the $r$-mode spectrum in isentropic stars. The possible existence of a continuous part in the frequency spectrum of nonsentropic stars (see, e.g., [18]) will be examined in future work.

In order to identify some of the peaks in the Fourier transform with specific inertial modes, we compare the ratio of their frequencies over the frequency of the $l = m = 2$ $r$-mode to the corresponding ratios derived from the normal mode linear eigenfrequencies, for a Newtonian $N = 1.0$ polytrope [19]. This comparison shows good agreement and allows us to identify, in Fig. 3, the $m = 2$ inertial modes with $l_0 = 3, 4, 5$ [19], having corresponding frequencies $f = 0.68, 1.16$, and $0.38$ kHz. In the Newtonian limit, one would expect the above inertial modes to have frequencies $f = 0.70, 1.14$, and $0.37$ kHz, for an $l = m = 2$ $r$-mode frequency of $1.03$ kHz, which shows that, even though the effects of relativity and rapid rotation on the individual frequencies are not negligible, the ratios of the $r$-mode frequency to the frequencies of the inertial modes remain close to their Newtonian, slow-rotation values.

We investigate the influence of the numerical viscosity of the code (due to finite differenting) on the evolution, by comparing evolutions at different grid spacings. Figure 4 shows the late-time evolution of the axial velocity in the equatorial plane at $\alpha = 1.0$, for $r = 0.75$ and $0.59$ km, for grid spacings of $\Delta x = 0.59$ km and $\Delta x = 0.31$ km. The decrease in the amplitude of the oscillation (when compared to the initial amplitude) scales as roughly first order with grid spacing. This indicates that the numerical viscosity damping the oscillations is dominated by the truncation error at the surface of the star, which is only first order (compared to the second-order accuracy in the interior; see related discussion in [13]).

In a recent paper [8] it is indicated that the hydrodynamical oscillatory motion of an $r$-mode in a rotating star may be accompanied by a differential drift (of kinematical origin). In the equatorial plane, the drift velocity has the opposite sign, compared to the rotational velocity of the unperturbed star. In a rotating star with a poloidal magnetic field, this kinematical drift may wind up the magnetic field lines and limit the effect of the $r$-mode instability [7,8]. The differential drift reported in [8] is second order in the $r$-mode amplitude, but it is derived from the first order in amplitude definition of the $r$-mode velocity. Still, in the restricted case of a spherical shell [20], the full nonlinear hydrodynamical equations yield the same kinematical drift as in [8], independent of the amplitude of the oscillations.

Using our code we attempted to investigate the presence of such a kinematical drift in the numerical evolutions. Our numerical results do indicate that the perturbed star is rotating slower near the surface (compared to the rotational velocity of the unperturbed star) and that this drift scales roughly as $\alpha^2$ for amplitudes of order $\alpha \sim 1.0$, although its magnitude is significantly smaller than estimated in [8]. We cannot test the presence of the drift for amplitudes much less than $\alpha \sim 1.0$, since it becomes smaller than numerical truncation errors. Even at $\alpha \sim 1.0$, it is not completely unambiguous that the drift we find is the same as that predicted by [8], as it could be due to the fact that our initial data do not correspond exactly to a single $l = m = 2$ $r$-mode and thus also contain other modes and violate (to some extent, which is still acceptable for the main purpose of this work) the relativistic Hamiltonian and momentum constraints. However, the fact that our drift scales as $\alpha^2$ points to an association with the results reported in [8]. Our preliminary findings on the differential drift need to be confirmed by simulations with more precise eigenfunctions and higher resolution before a definite conclusion can be drawn.

In summary, our hydrodynamical evolutions of nonlinear $r$-modes show that (1) nonlinear hydrodynamical couplings would not prohibit $r$-modes in isentropic,
rapidly rotating relativistic stars from attaining a large amplitude, of order one, when driven unstable by gravitational radiation reaction—it remains to be investigated whether nonlinear saturation could set in at time scales longer than the dynamical one, (2) $r$-modes and inertial modes, in isentropic stars, are discrete, and (3) a kinematical drift associated with $r$-mode oscillations appears to be present in our simulations, although an unambiguous confirmation will require more precise initial data.

Our finding that gravitational radiation could drive unstable $r$-modes to a large amplitude implies that $r$-modes can easily melt the crust in newly born neutron stars [21], leaving the initial conclusions about the $r$-mode instability being a strong source of gravitational waves (see [5,6]) essentially unchanged. Our results also imply that $r$-modes could be excited to large amplitudes during the violent formation of a protoneutron star, after a supernova core collapse [22]. Our finding that $r$-modes oscillate at a predominantly discrete frequency (at least in isentropic stars) significantly simplifies the expected gravitational wave signal compared to a signal emitted by a source with a significant continuous part in its frequency spectrum. The indications for the existence of the kinematical drift call for a more detailed consideration of the interaction between $r$-modes and magnetic fields.

We will present details and tests of our numerical method, as well as extensive higher-resolution results for various equations of state and rotation rates, in a forthcoming paper [23]. In future work, we plan to study $r$-modes in nonisentropic stars and implement an accelerated gravitational-radiation reaction force (see [24]).

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