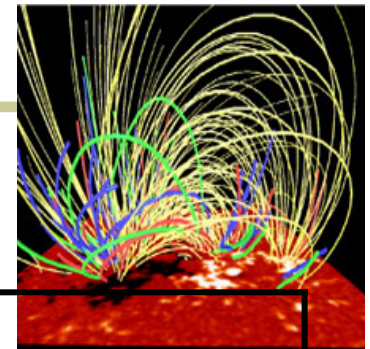


# ***Particle acceleration and radio emission from complex active regions***

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9/6/2004

# Main questions

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- Where are electrons accelerated?
- How are they accelerated?
- Can they all be accelerated at the loop top?

**For  $N = 10^{38}$  electrons in flare, where**

$$N = V n_e = L^3 n_e$$

- **For  $n_e = 10^{10} \text{cm}^{-3}$**
- **$L = 2 \cdot 10^9 \text{ cm}$  (0.5 arcminute!)**

# Model Questions

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- We need a **magnetic topology**
- We need an **energy release** mechanism.
- We need an **acceleration mechanism** with the ability to accelerate ions and electrons very efficiently
- We need **emission mechanism (s)**

## We can assume that

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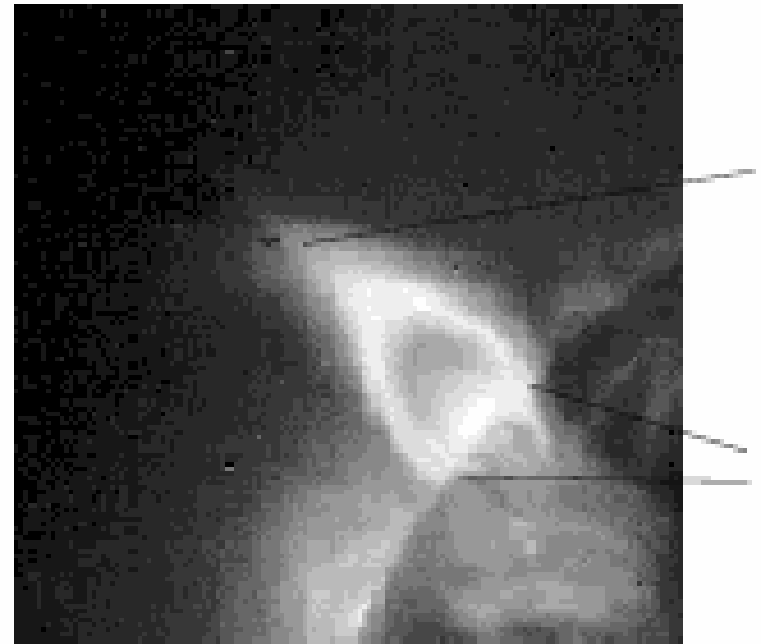
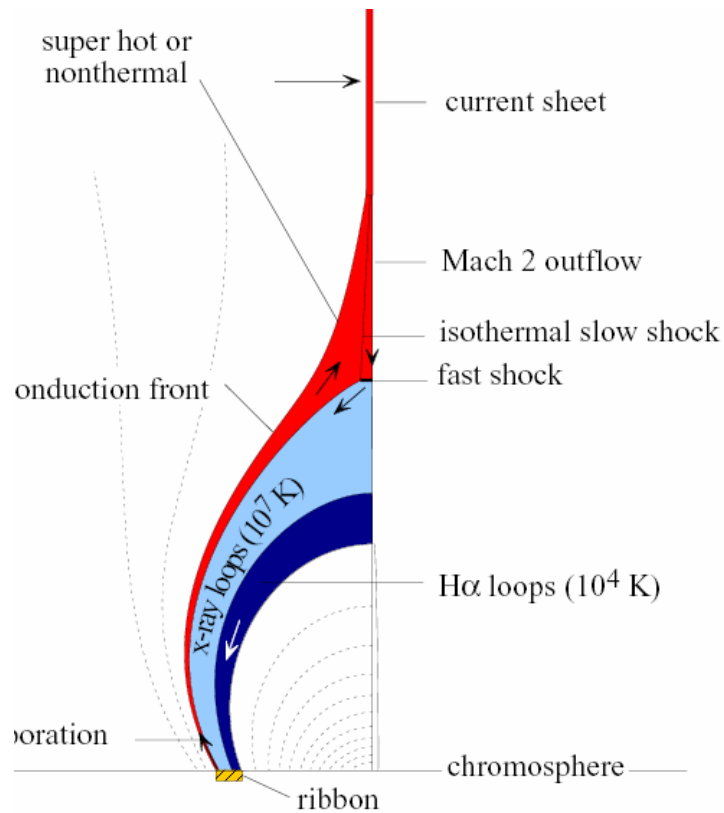
- The **magnetic topology** is simple or complex: a magnetic **loop or millions** of magnetic loops derived from the extrapolation of the magnetograms.  
(Unfortunately this is a free parameter)
- The **energy release** is magnetic dissipation in unstable current sheets
- The emission mechanism may be :  
**incoherent** ( **Bremsstrahlung** or **synchrotron** and/or **coherent** (plasma processes (beam plasma, loss cone))

# The acceleration mechanism?

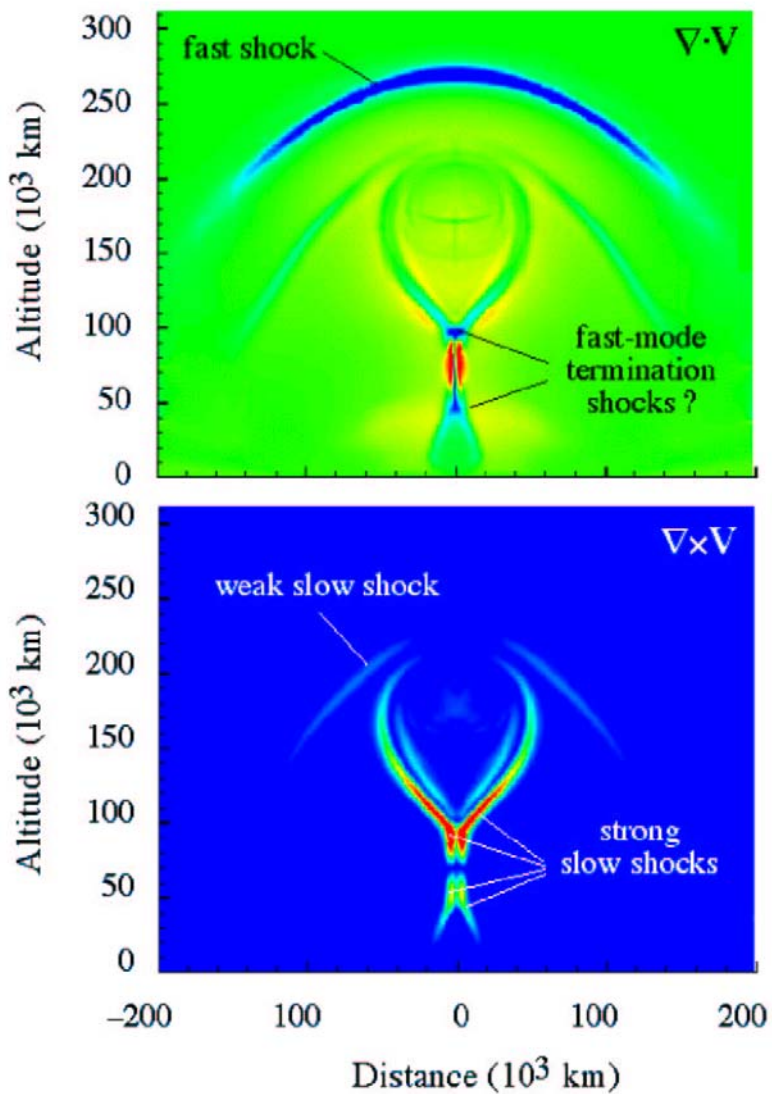
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- **Electric fields** inside the current sheet or directly driven (**Sub-Dreicer** or **super-Dreicer**)
- **Shocks** supported by the energy release processe(s)
- **MHD Turbulence** excited by the energy release

# The “standard(?)” solar flare model and the monolithic current sheet



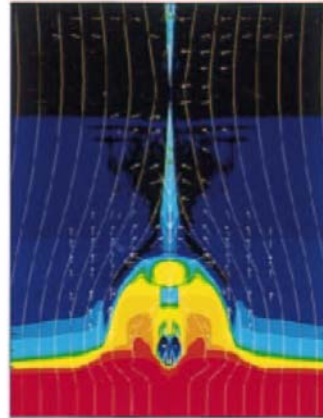
## Shock Formation



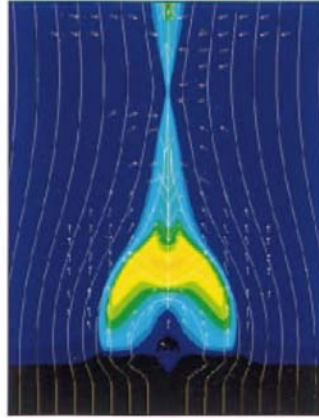
# Most Realistic Simulations of Flare Loops To Date

Yokoyama & Shibata (2001)

Density



Temperature

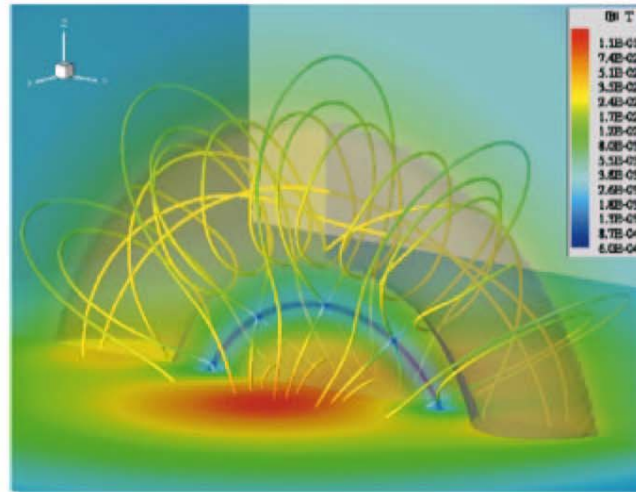




# 3D Flux Rope Simulation

Roussev et al. (2003)

before

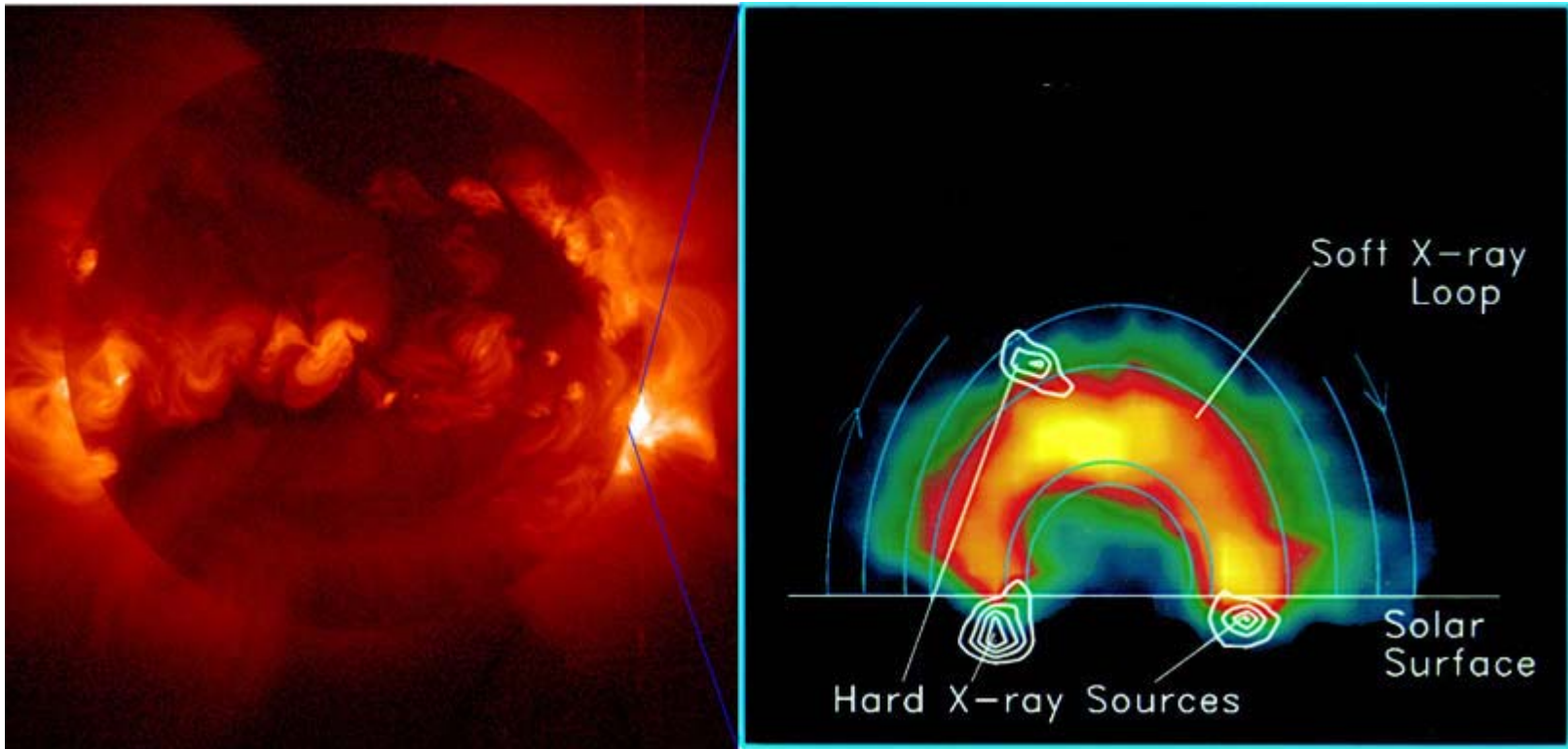


after



# Impulsive Flare Geometry

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# MHD Turbulence the “standard” acceleration model

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1. Alfvén and fast mode waves generated at large scales (assumption)
  2. Cascade to higher wave numbers (e.g., Zhou & Matthaeus 1989)
  3. Fast mode waves energize electrons via transit-time acceleration (e.g., Miller 1997)
  4. Alfvén waves energize ions via gyroresonant acceleration (e.g., Miller & Roberts 1995)
- ➔ Both species accelerated by MHD turbulence

# **We have a problem!**

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- The standard model for energy release is in favor of E-field and shock acceleration
- The standard acceleration mechanism is in favor of a loop filled with MHD turbulence
- Who is correct?

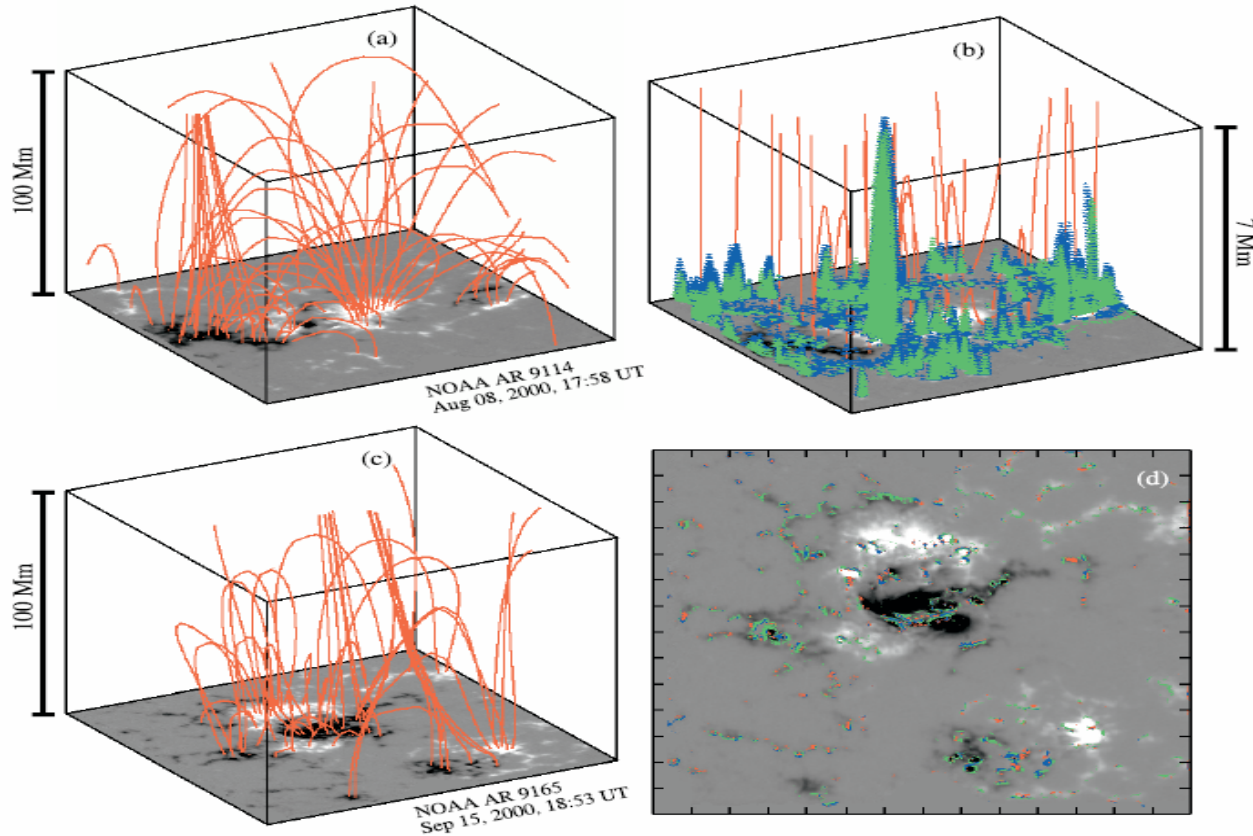
## For those looking for a standard model

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- "Solar Flares are highly individualistic"
  - Smith and Smith, 1963
- No need for the flare model

# Building up unstable discontinuities from the photosphere

(Vlahos+Georgoulis, ApJL, 2004)

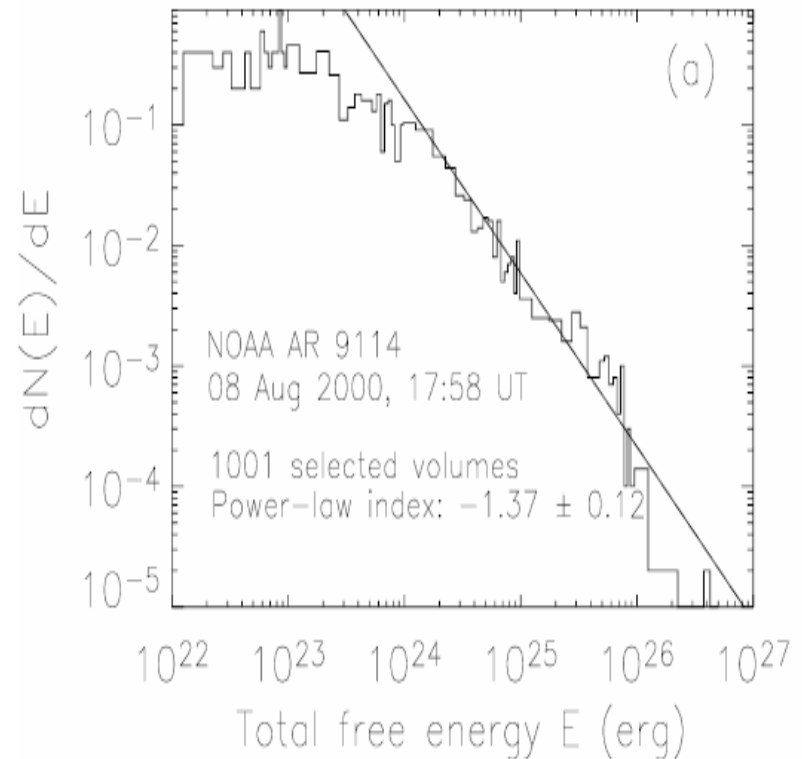


# How do you define an unstable discontinuity

- We mark the points where (Parker's criterion)

$$\vec{J}_c \sim \nabla \times \vec{B}$$

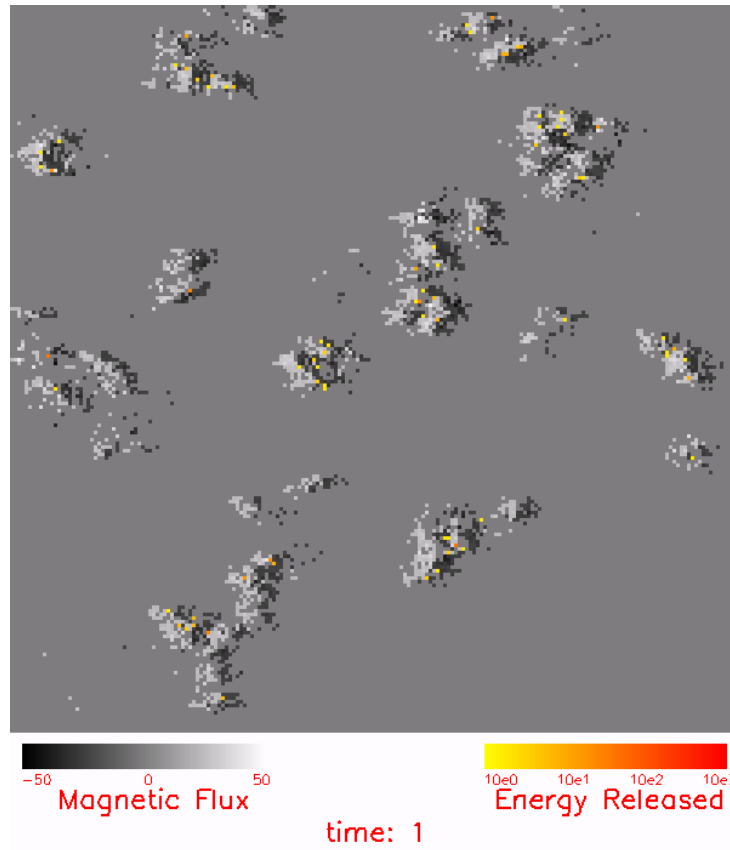
is satisfied and multiply this volume with the magnetic energy in excess the potential energy



# Evolving active regions build up constantly magnetic discontinuities....

(Fragos, Rantziou, Vlahos, AA, 2004)

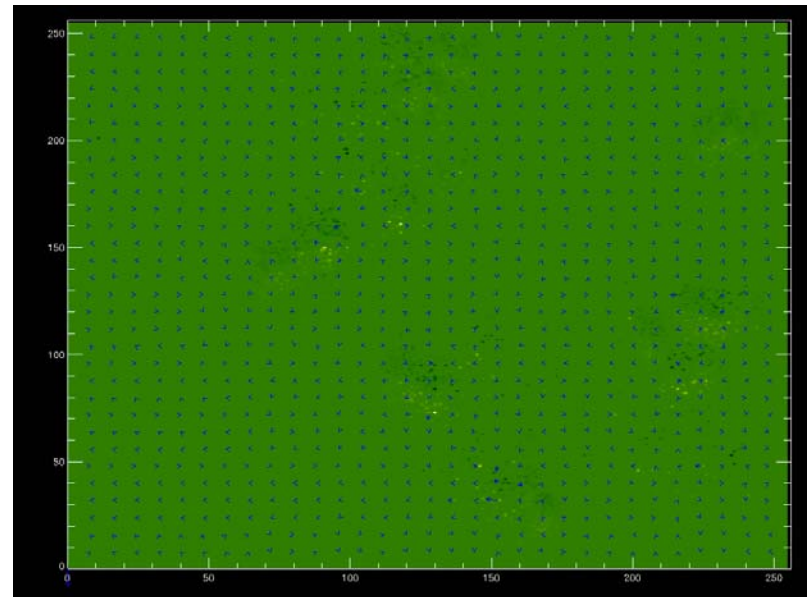
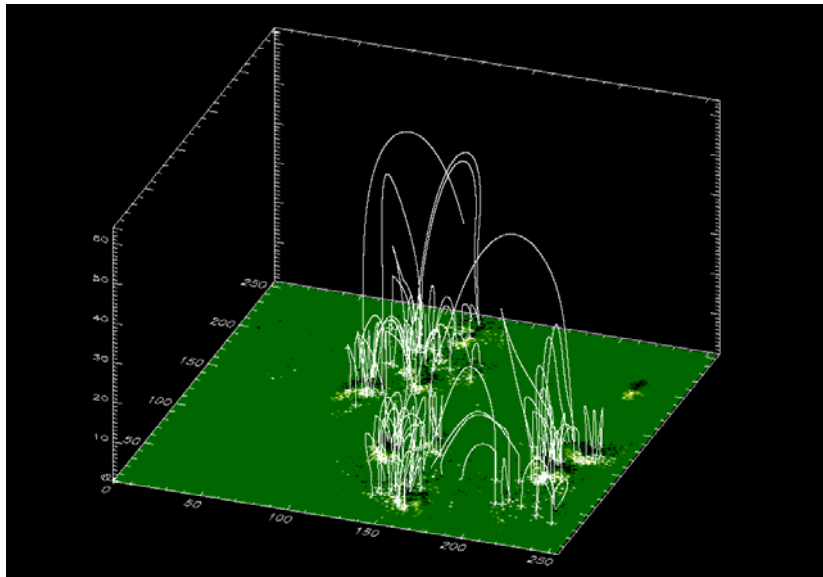
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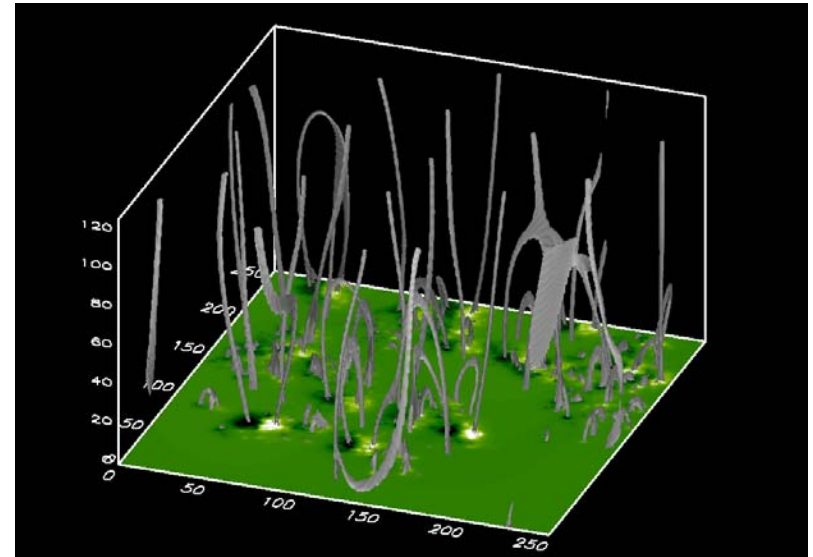
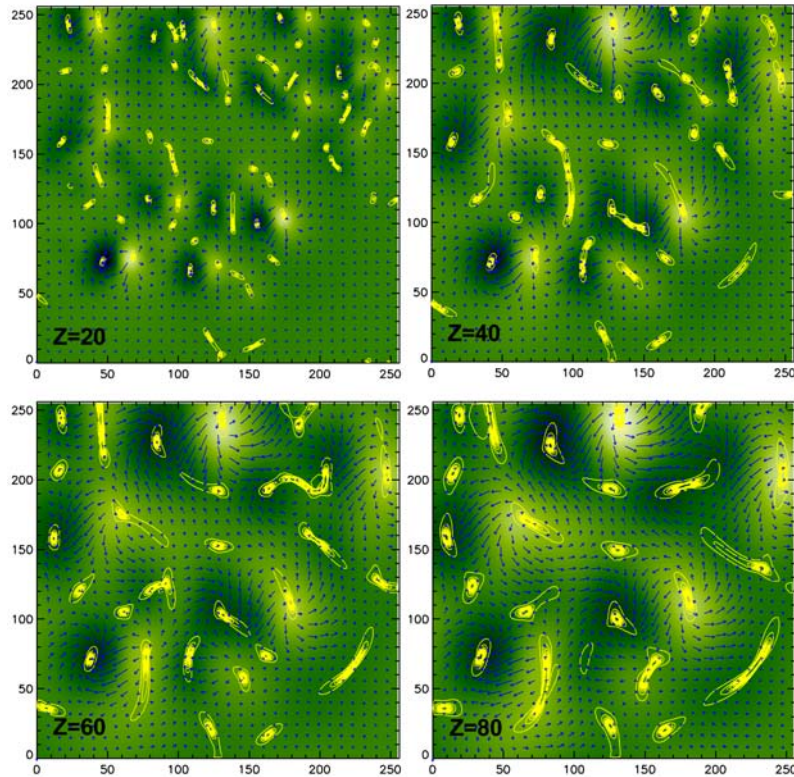


# Evolving active regions build up constantly magnetic discontinuities.... (Fragos, Rantziou, Vlahos, AA, 2004)

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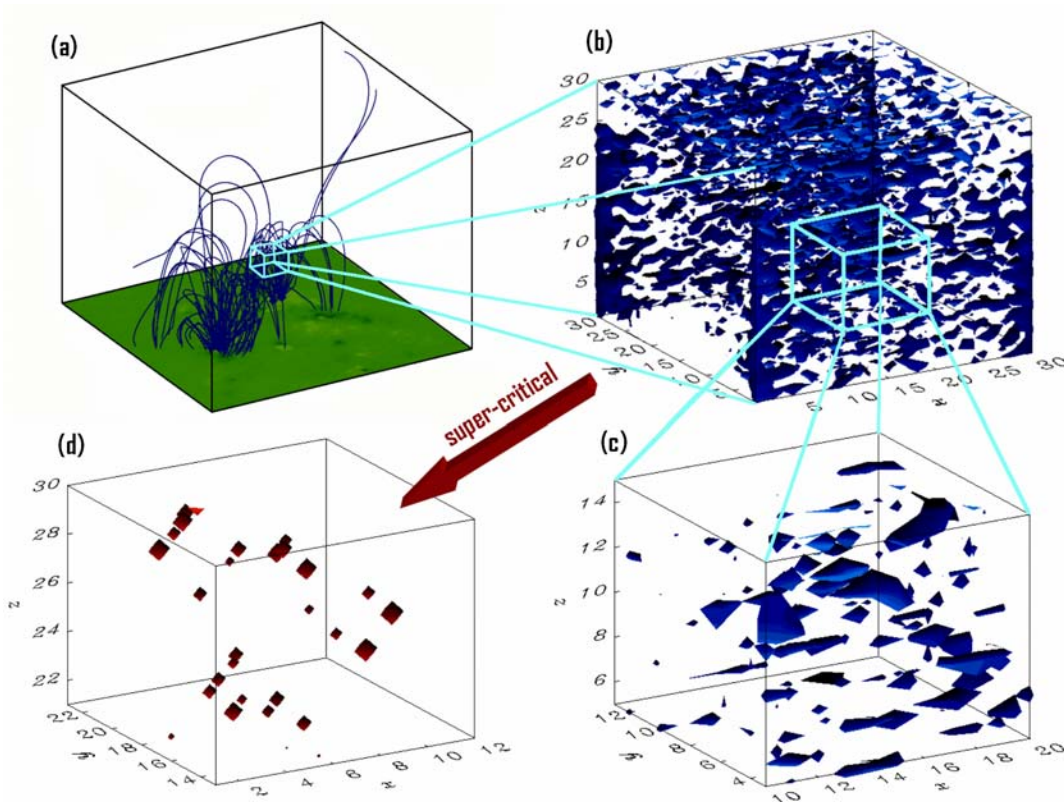


# Dynamic motion of the photosphere builds constantly magnetic discontinuities (Fragos, Rantziou, Vlahos, AA, 2004)



# A New approach to an old problem

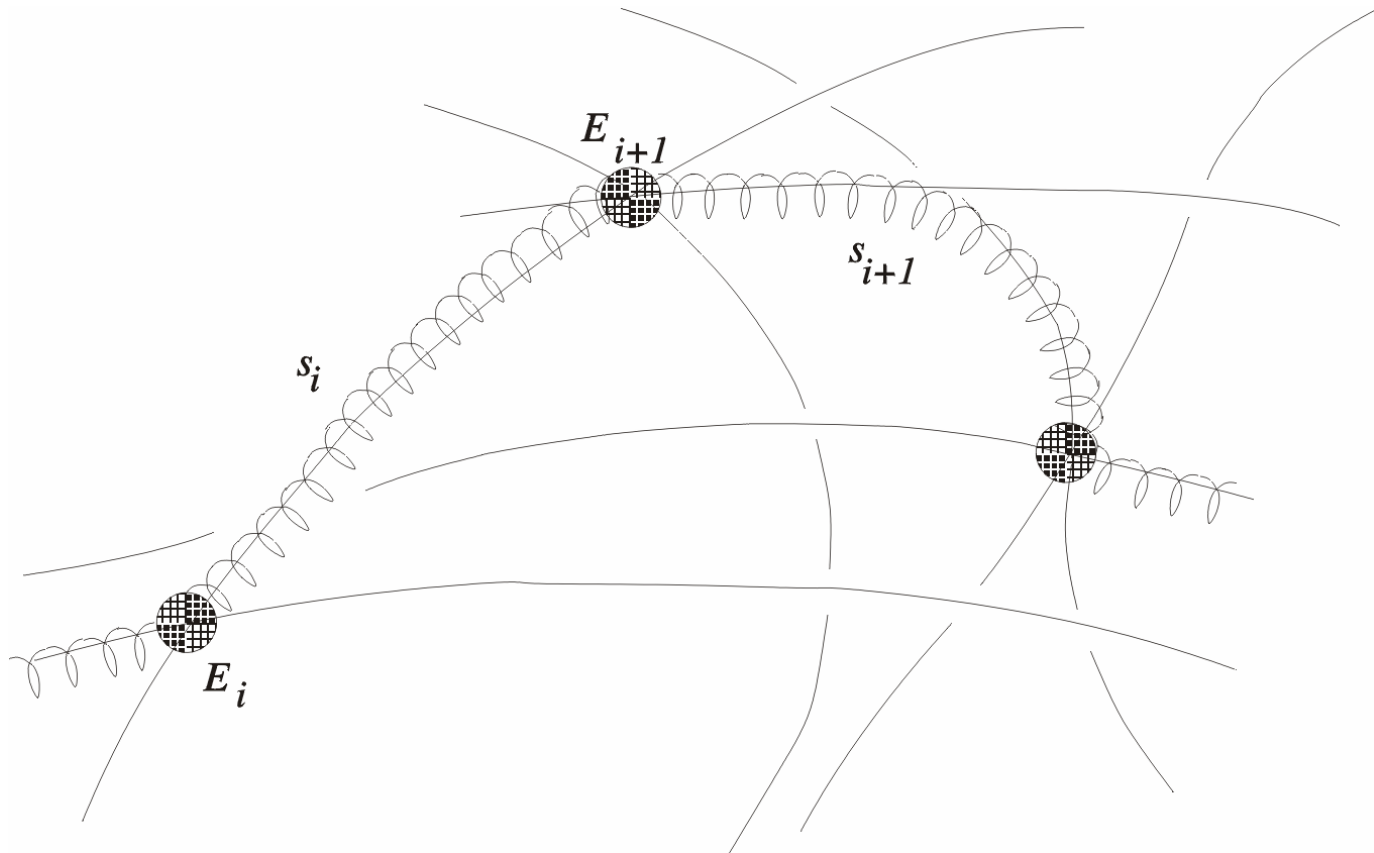
- From one current sheet to millions



# Sporadic formation of current sheets

Vlahos, Isliker and Lepreti (ApJ, June 10, 2004)

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# A 'Turbulent' Field Model (stochastic but not resonant accelerator) (Azner+Vlahos, APJL, 2004)

$$\mathbf{A} = \sum_{\mathbf{k}} \mathbf{a}_{\mathbf{k}} \cos(\mathbf{k} \cdot \mathbf{x} - \omega)$$

$$\langle |\mathbf{a}_{\mathbf{k}}|^2 \rangle \sim (1 + \mathbf{k}^T \mathbf{S} \mathbf{k})^{-\nu}$$

random  $\phi_{\mathbf{k}}$

$$\mathbf{B} = \nabla \times \mathbf{A}$$

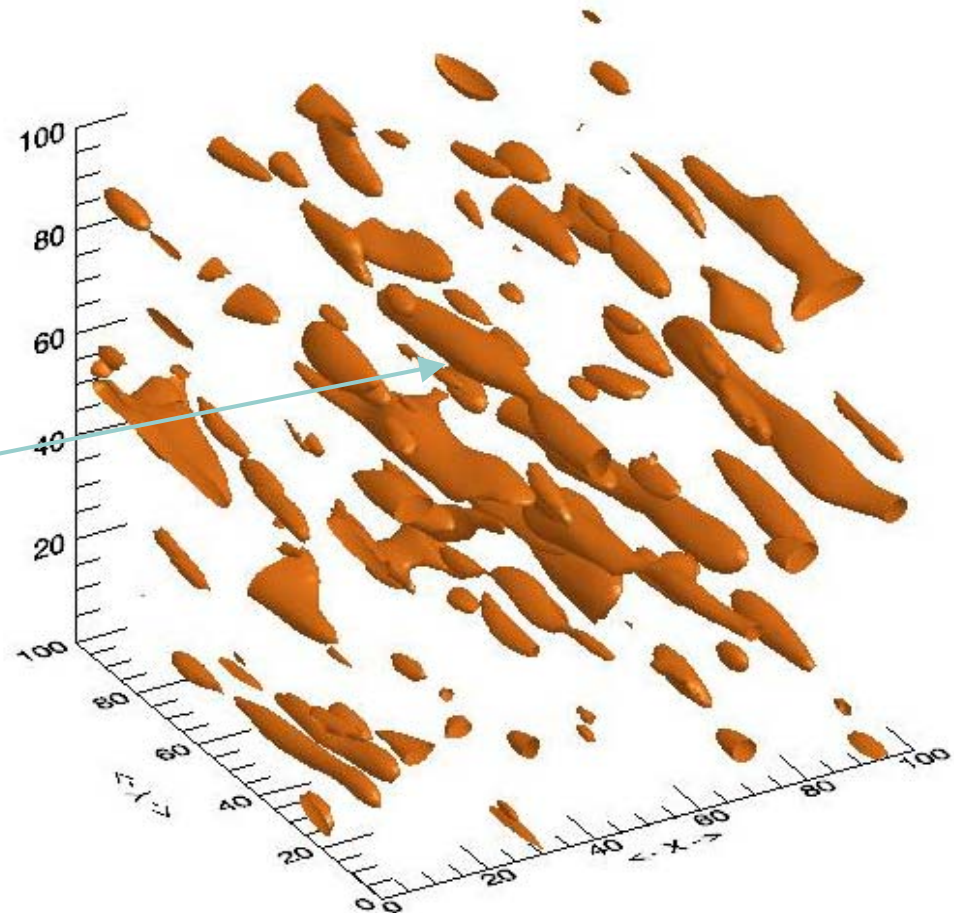
$$\vec{J} \sim \nabla \times \vec{B}$$

threshold  $j_c$

$$\mathbf{E} = -\partial_t \mathbf{A} + \eta(\mathbf{j}) \mathbf{j}$$

$$\partial_y \ll \partial_x, \partial_z$$

$$v_{ph} \ll |v_{ptcl}|$$





# Relativistic Particle Dynamics

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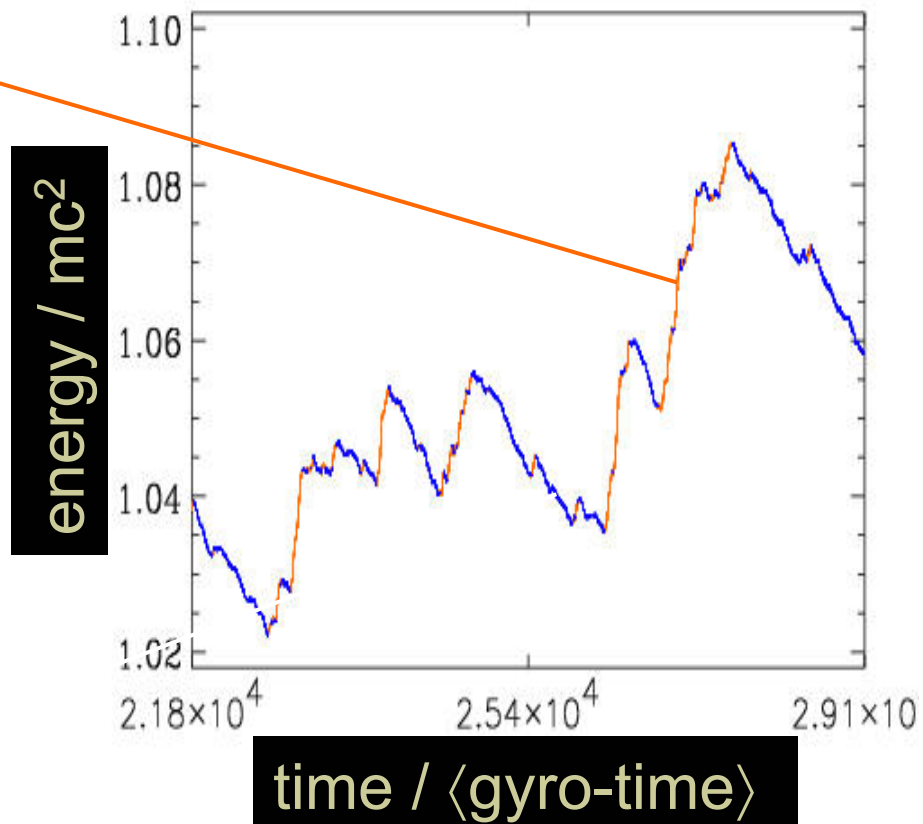
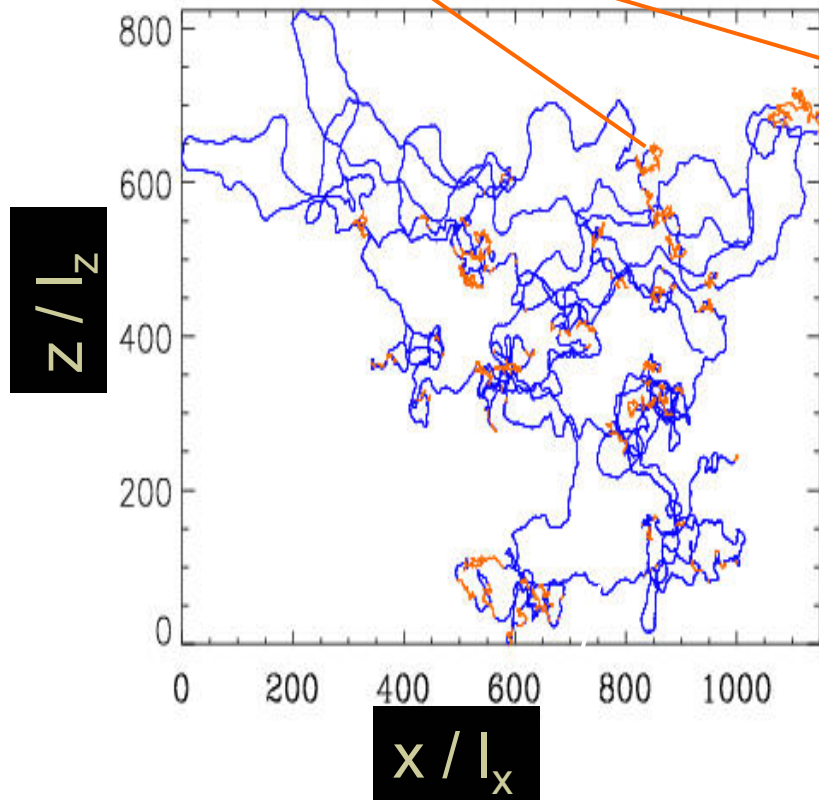
$$\frac{d\vec{r}}{dt} = \vec{v}$$

$$\frac{d\vec{p}}{dt} = e\vec{E} + \frac{e}{c}\vec{v} \times \vec{B}$$

- Relativistic equations of motion are solved numerically with adaptive step-size scheme
- Magnetic and electric fields are interpolated to provide field values in between grid-points

# Finding: intermittent particle orbits

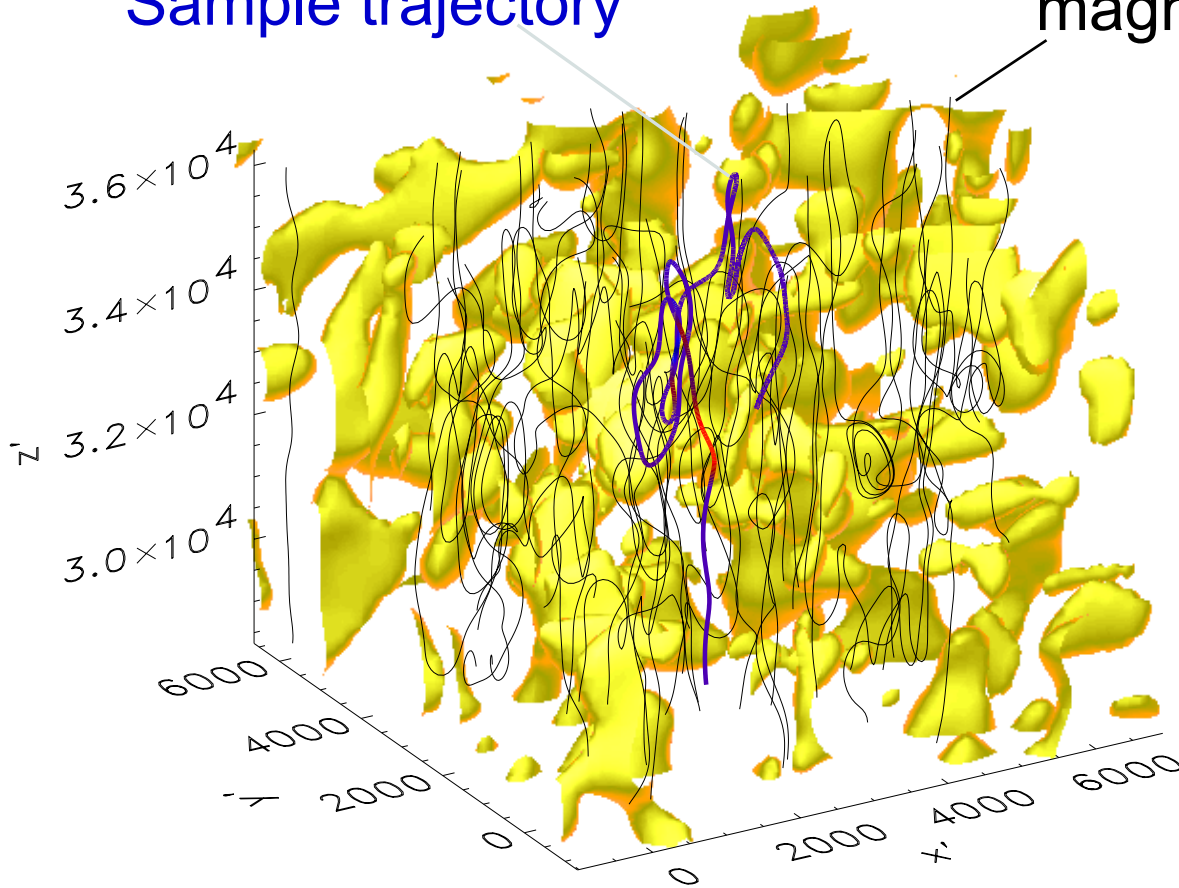
acceleration **within** local dissipation regions



# Dissipation Regions

Sample trajectory

magnetic field lines



- $B_0 = 20\text{G}$ ,  $\delta B_x \sim \delta B_y \sim 20\text{G}$ ,  $\delta B_z \sim 100\text{G}$

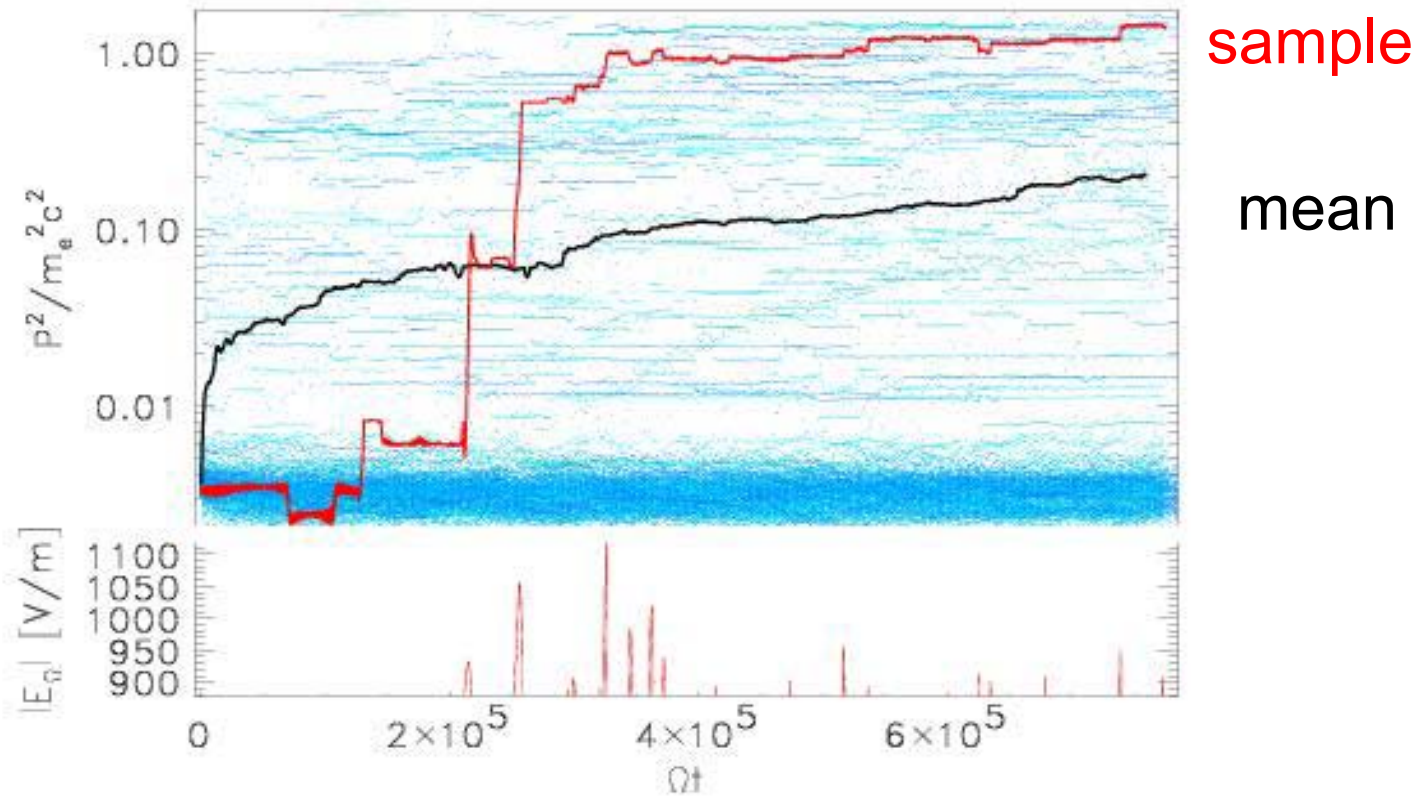
- Outer scales  
 $l_x=l_y=1\text{km}$ ,  $l_z=20\text{km}$

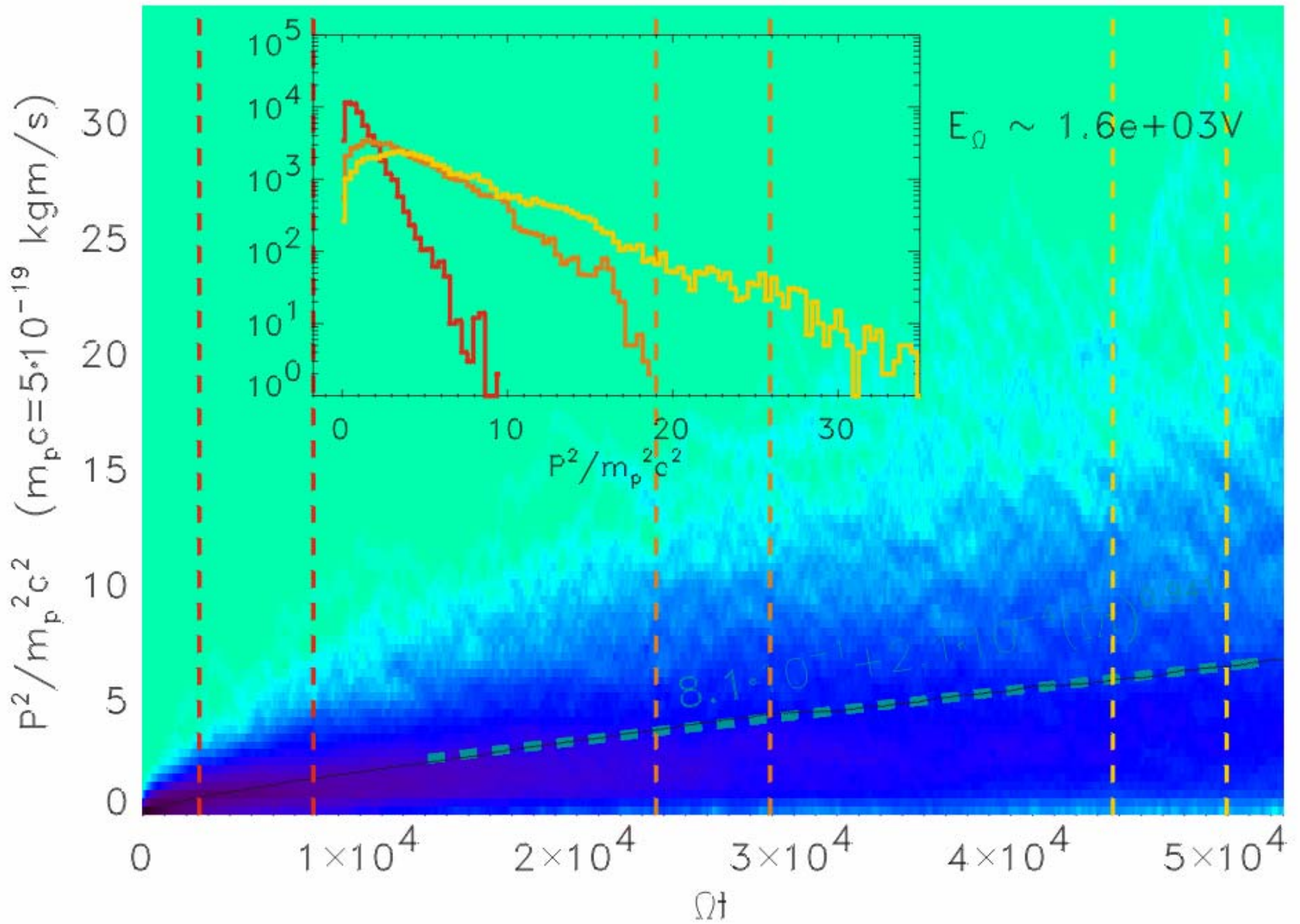
- Matern class PSD  
with index  $\alpha=1.5$

- Threshold current  
 $j_c = en_c v$  exceeded in  
5% of volume

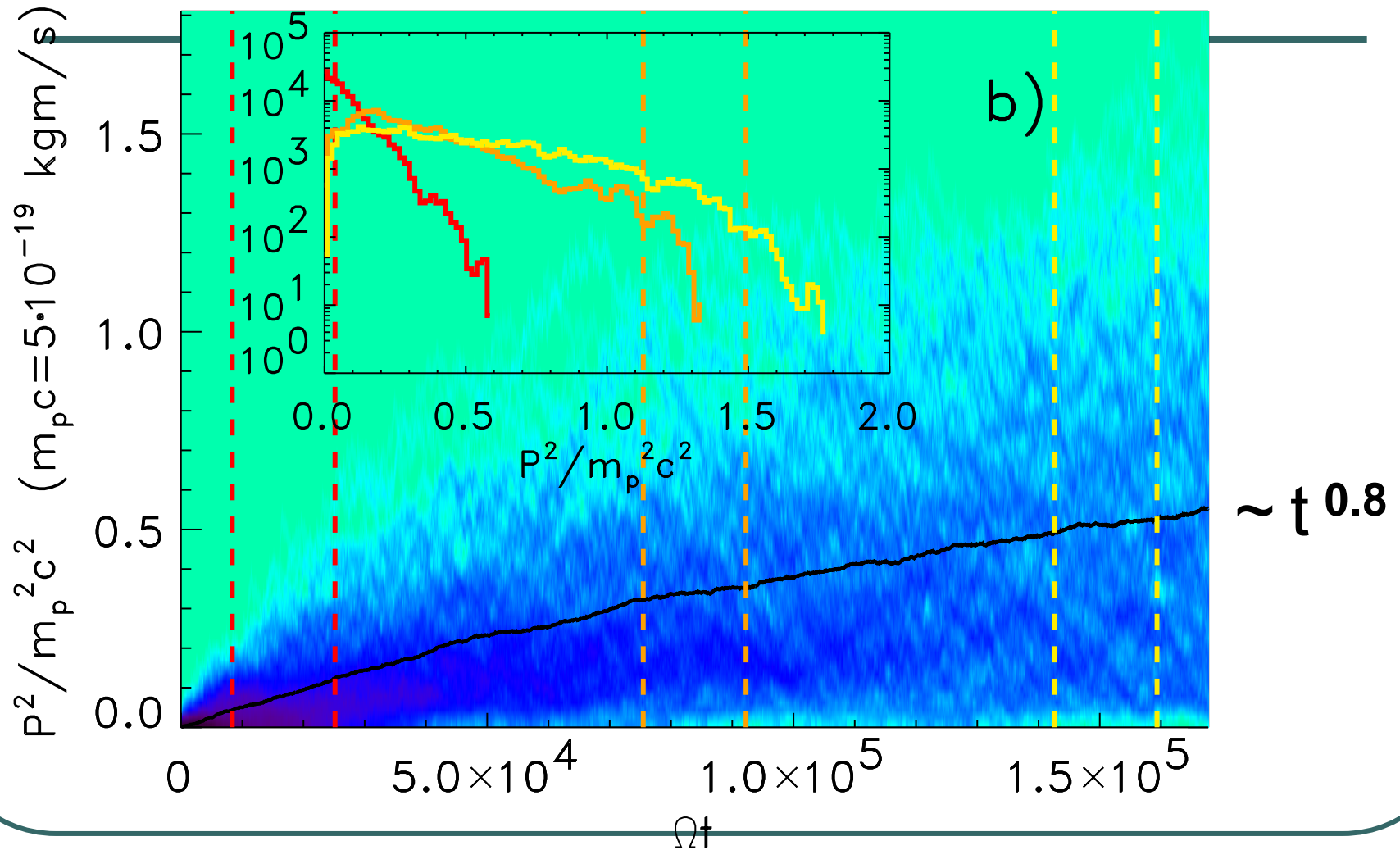


# Electron Acceleration





$$E_{\Omega} / E_D = 10^5$$



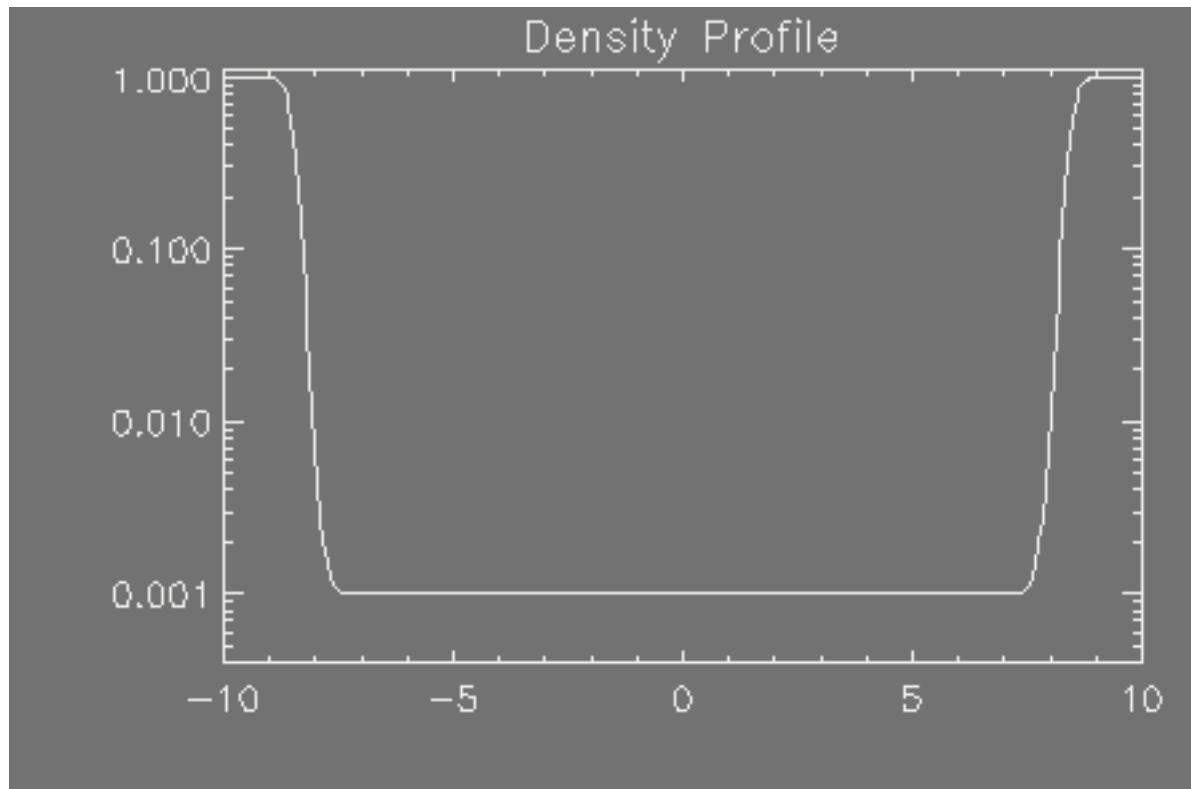
# Solve the MHD equations inside a simple loop atmosphere (*Galsgaard*)

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$$\begin{aligned}\frac{\partial \rho}{\partial t} &= -\nabla \cdot \rho \mathbf{u}, \\ \frac{\partial \rho \mathbf{u}}{\partial t} &= -\nabla \cdot (\rho \mathbf{u} \mathbf{u} + \underline{\underline{\tau}}) - \nabla P + \mathbf{J} \times \mathbf{B} + \mathbf{F}_e, \\ \frac{\partial e}{\partial t} &= -\nabla \cdot (e \mathbf{u}) - P \nabla \cdot \mathbf{u} + Q_{\text{Joule}} + Q_{\text{visc}}, \\ \frac{\partial \mathbf{B}}{\partial t} &= -\nabla \times \mathbf{E}, \\ \mathbf{E} &= -(\mathbf{u} \times \mathbf{B}) + \eta \mathbf{J}, \\ \mathbf{J} &= \nabla \times \mathbf{B}\end{aligned}$$

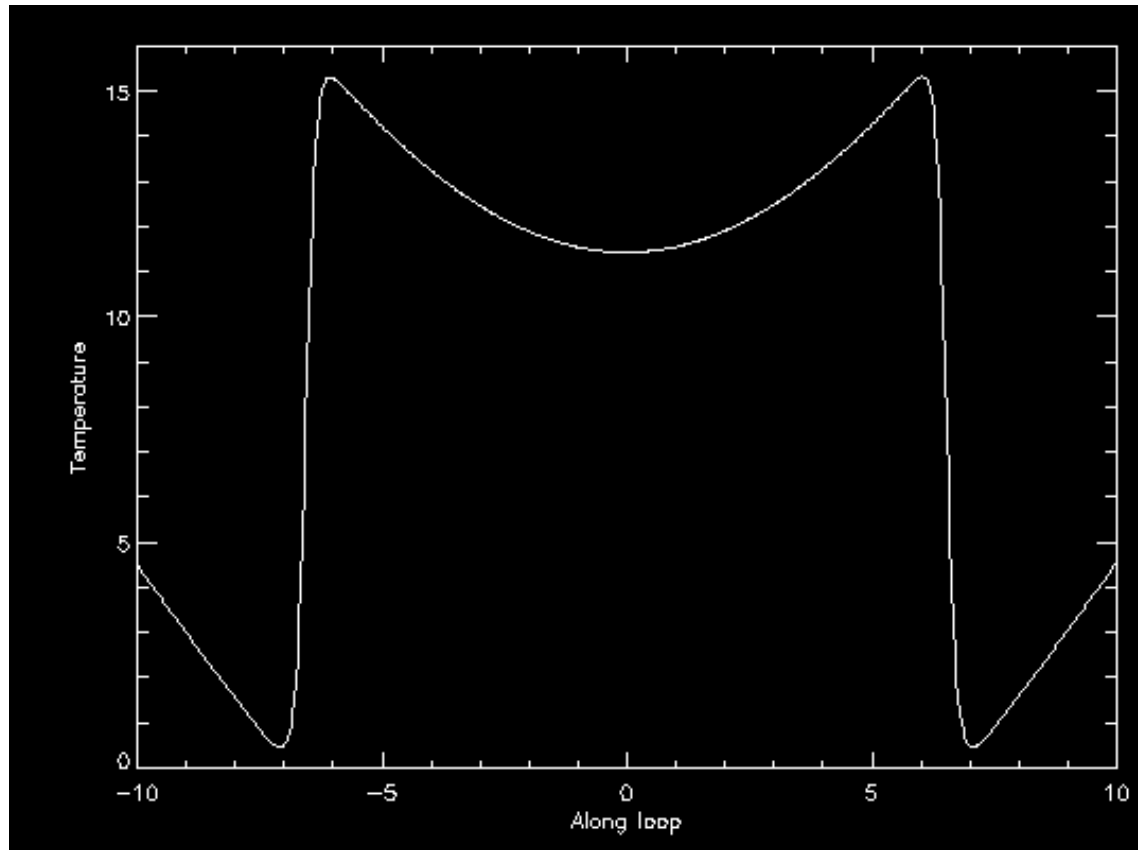
# Density profile along the loop

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# Temperature along the loop

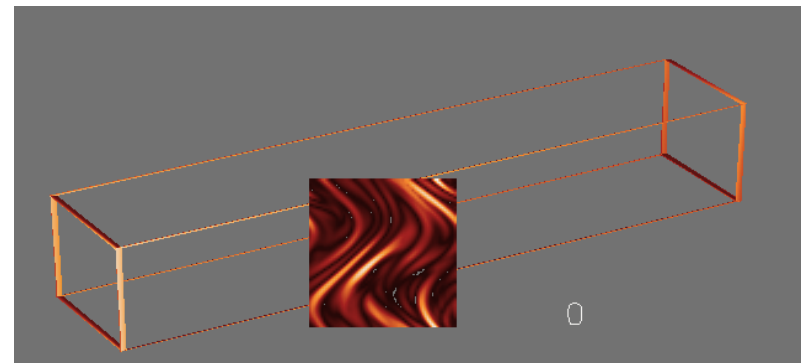
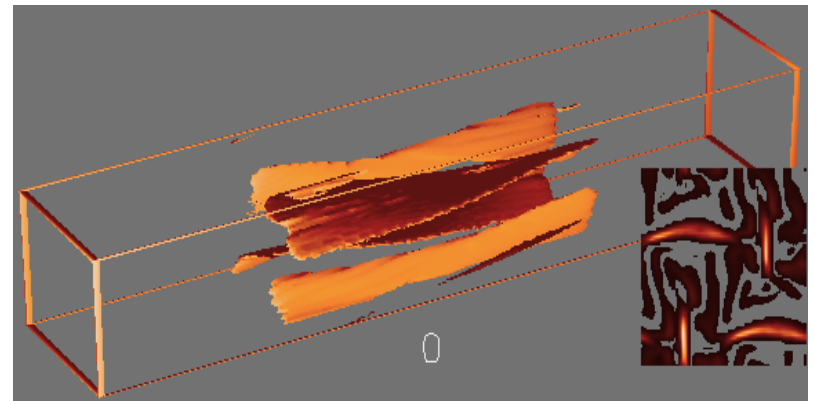
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# The stochastic loop model (Galsgaard)

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- 3D MHD experiment of photospherically driven slender magnetic flux tubes
- Continued random driving of the foot points (incompressible sinusoidal large scale shear motions )
- Reconnection jets generate secondary perturbations in  $B$
- Formation of stochastic current sheets

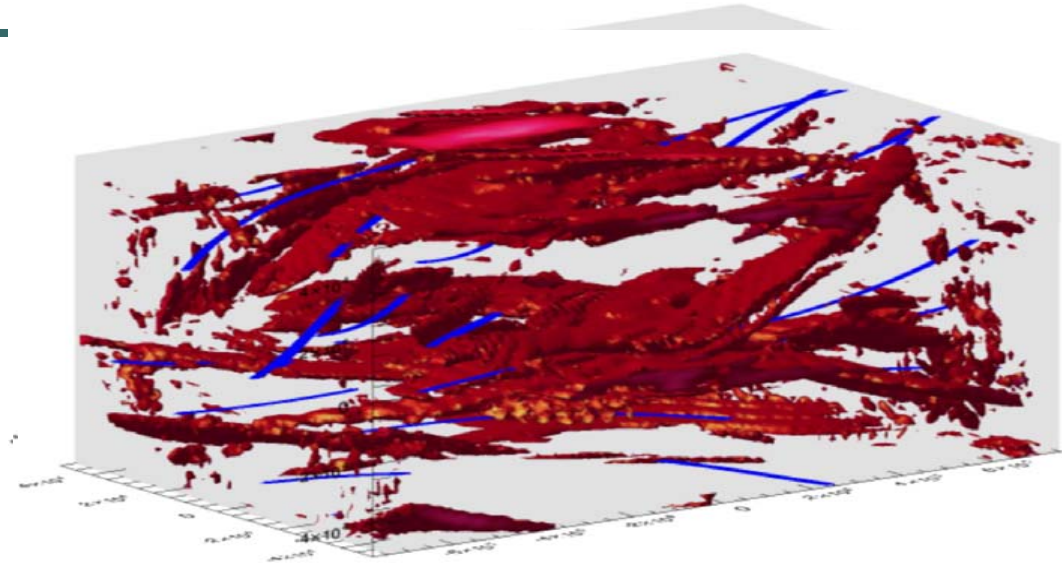




# Particle acceleration in stochastic current sheets

(Rim Turkmani et al)

- Particles injected at random positions within an MHD box
  - Protons 0.027 keV
  - Electron 1.16 keV
- Initial velocity fixed in amplitude, random in direction

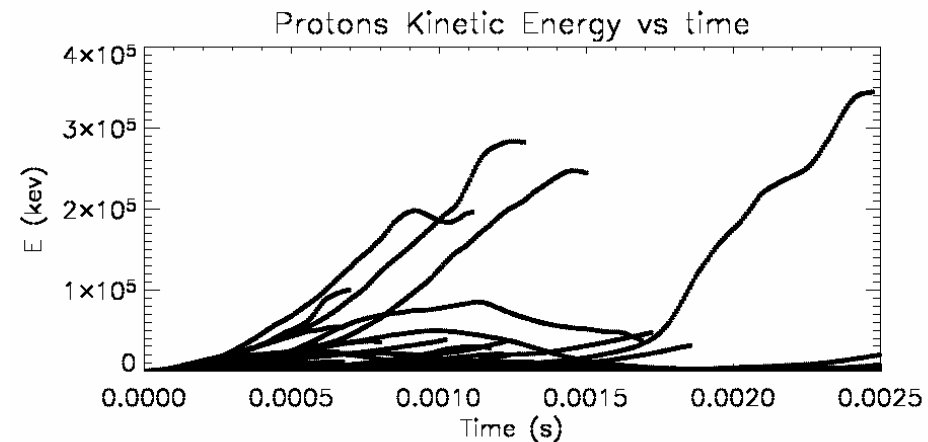
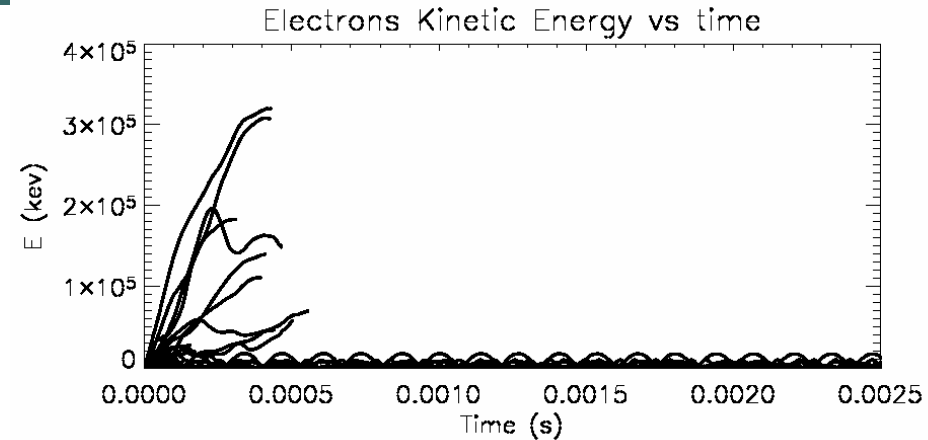


- Acceleration time scale much shorter than MHD time scale
- B and E are scaled; initial values:
  - B: Mean  $\sim 1.0$  (0.89 – 1.08)
  - E: Mean  $\sim 7e-4$  ( $e-5$  –  $e-2$ )

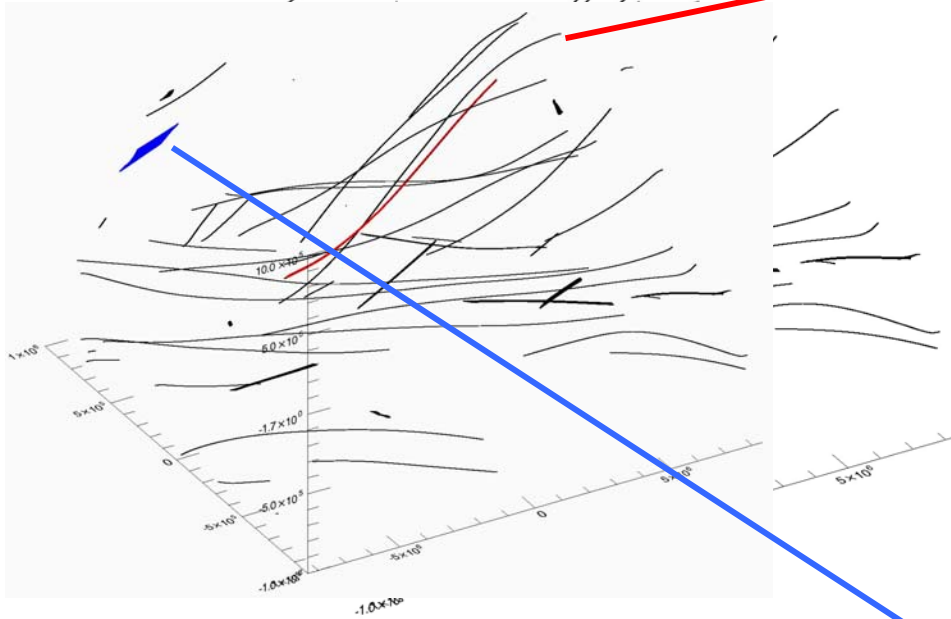


# Electrons versus Protons

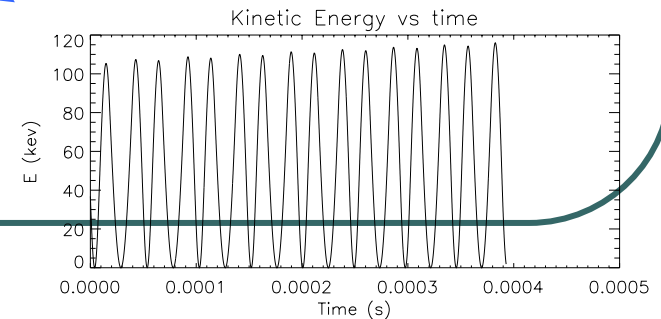
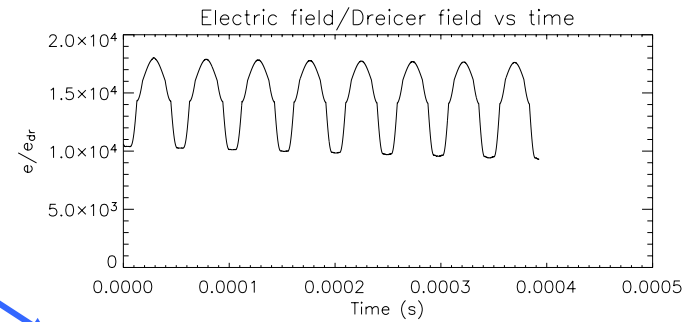
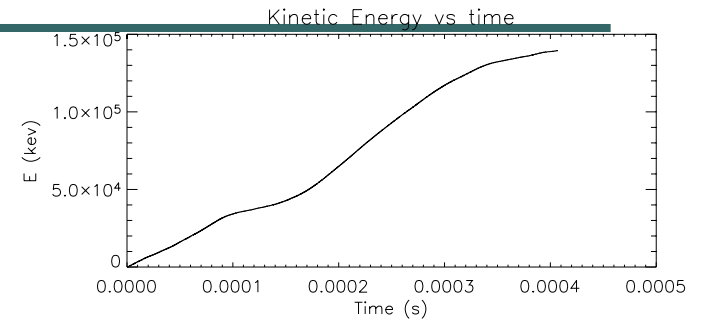
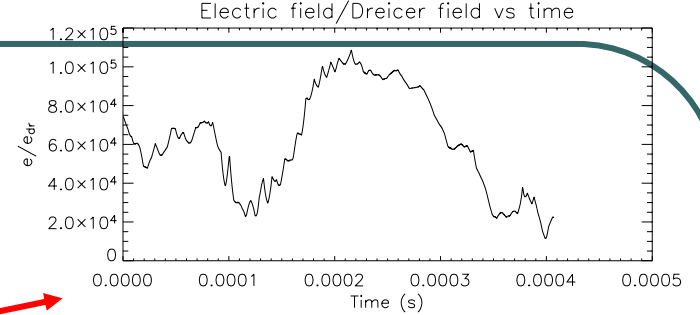
- Electrons and protons reach  $\sim 300$  Mev in a  $1.5e7$  cm long loop with  $B=100$  G within:
  - 0.5 ms for Electrons
  - 2.5 ms for Protons
- In general, electrons are accelerated  $\sim 5-6$  times faster than protons



# Particles motion

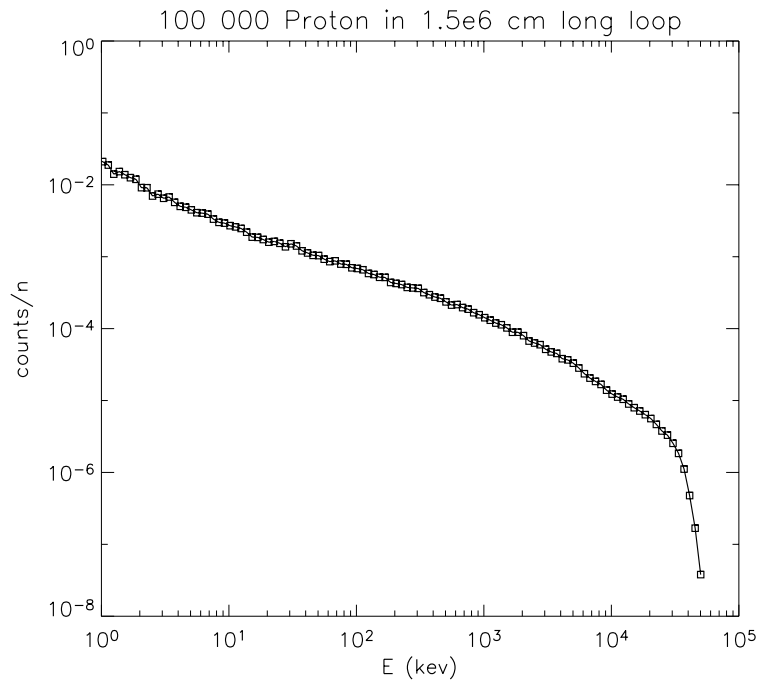


- 50 Electron run for 0.6 ms

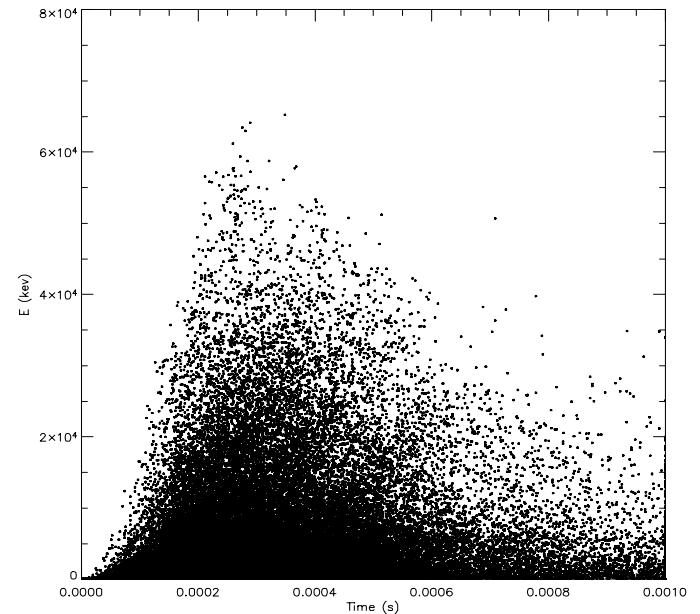


# Distribution Functions

- 100,000 proton in 100 G magnetic field run for 1 ms

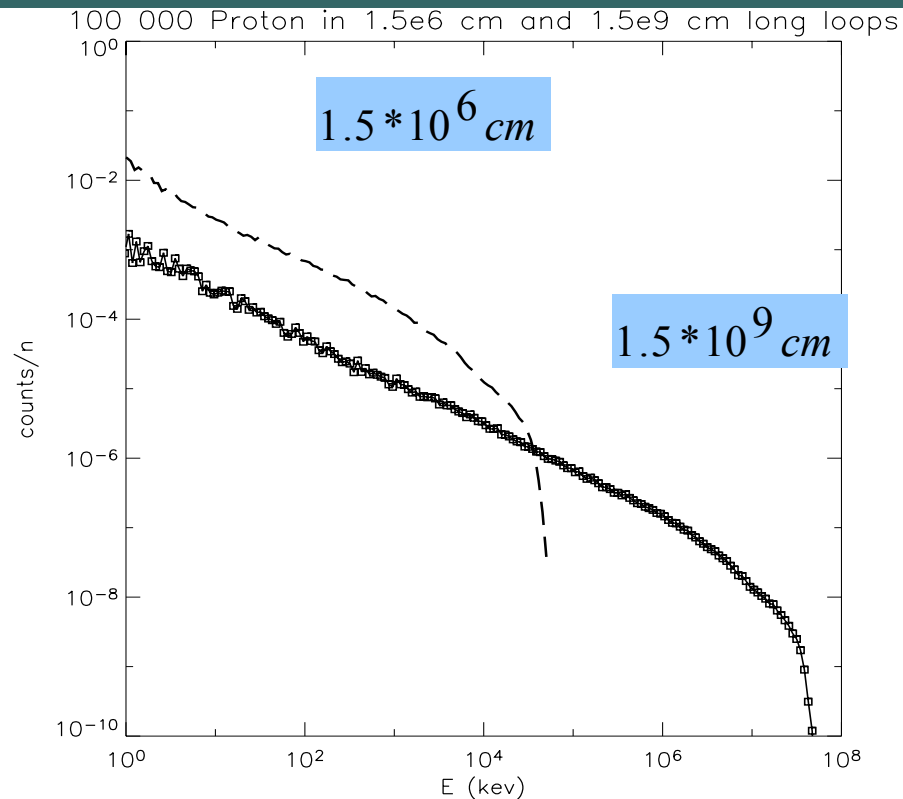


- Two parts power law



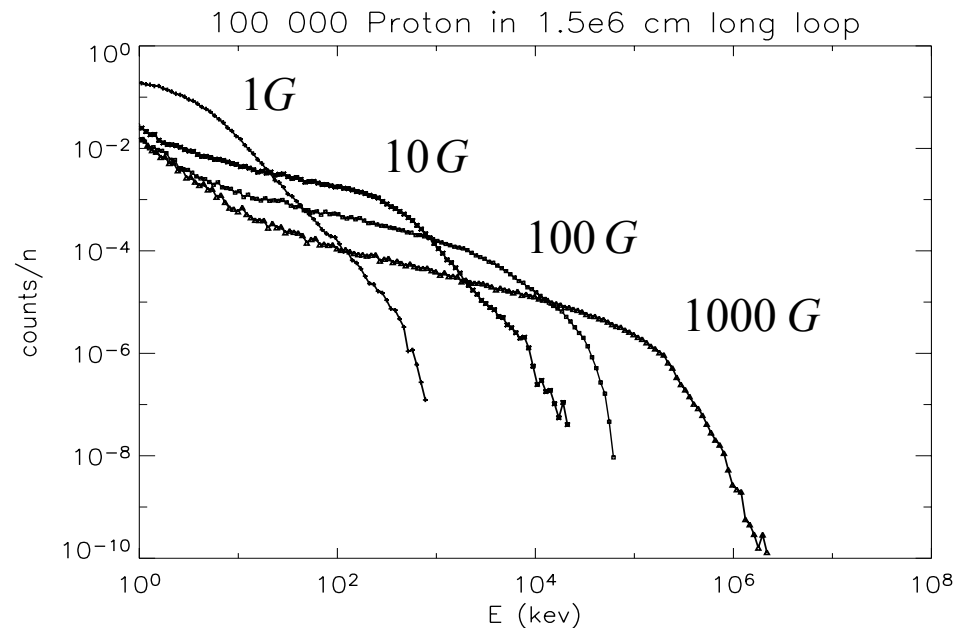
- $\sim 60$  Mev in 0.3 ms

# Scaling with loop dimensions



Acceleration scales linearly with the spatial scale of the loop

# Scaling with E and B



Acceleration scales almost linearly with the values  
of the magnetic and electric fields

# Quit time acceleration: an example of MHD Turbulence

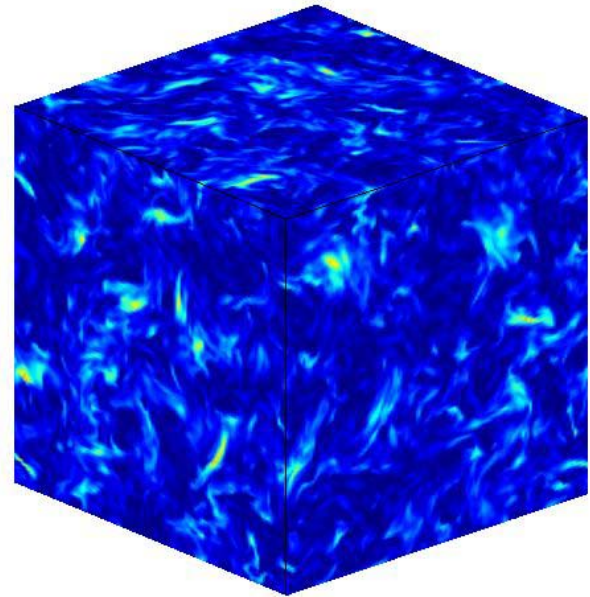
(Lepreti, Isliker, Vlahos, submitted)

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- The presence of electric fields in a driven magnetic field forming continuously magnetic discontinuities

$$\vec{j}(r,t) = \frac{c}{4\pi} \nabla \times \vec{B}(r,t)$$

$$\vec{E}(r,t) = -\frac{\vec{v} \times \vec{B}(r,t)}{c} + \eta \vec{j}(r,t)$$



# My summary

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- The photospheric motions drive the formation of unstable discontinuities
- Fast, slow, organized and random flows are all part of the photospheric activity
- New emerging flux adds complexity to this picture and enhances the concentration of magnetic discontinuities
- The extrapolated force free magnetogram holds important information for the activity of the complex AR.

# My summary

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- In the model presented the formation of critical and sub-critical discontinuities is crucial.
- The flare may start from any critical discontinuity and spread to large volumes creating an avalanche
- The net result is the formation of distributed electric fields of all scales inside the magnetic field created by the extrapolated topology.



# How can we observe all this?

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- My proposal is that decimetric spikes are the crucial link to this model. We may be looking the accelerators in action.
- The loop like structures are going to remain and are now have a filling factor of current sheets inside.
- Particles are transported inside an electrically active environment.
- The E-fields are not aligned, so no need for return current problems. The formation of million current sheets is natural and the acceleration is stochastic.

# Why loop tops

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- The particles are accelerated perpendicular to magnetic fields and since they propagate in stochastic field lines are very efficiently trapped.
- Foot points are due to precipitating particles which are constantly accelerated as they propagate towards the photosphere.
- There is acceleration outside the loop (Type III)
- There is acceleration before the flare and there is acceleration long after the flare....due to the large enhancement of discontinuities
- The acceleration is going hand and hand with the formation of discontinuities so it can be present in all scales
- The foot point sources are not unique, they have internally structure but are currently unresolved