



Nonlinear Plasma Effects in Natural and Man-made Aurora



Evgeny Mishin and Todd Pedersen

Space Vehicles Directorate

Air Force Research Laboratory

Bedford, MA, USA



Modern Challenges in Nonlinear Plasma Physics
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OUTLINE

Natural and Artificial Auroras

- Introduction: “*Ordinary*“ and *Enhanced Aurora*
- Collisional vs. Collisionless (Beam-Plasma) Interaction
- SLT Aurora: *Plasma Turbulence Layer*

❖ Underlined text indicates Dennis's significant contributions to understanding these problems

HF-induced Airglow

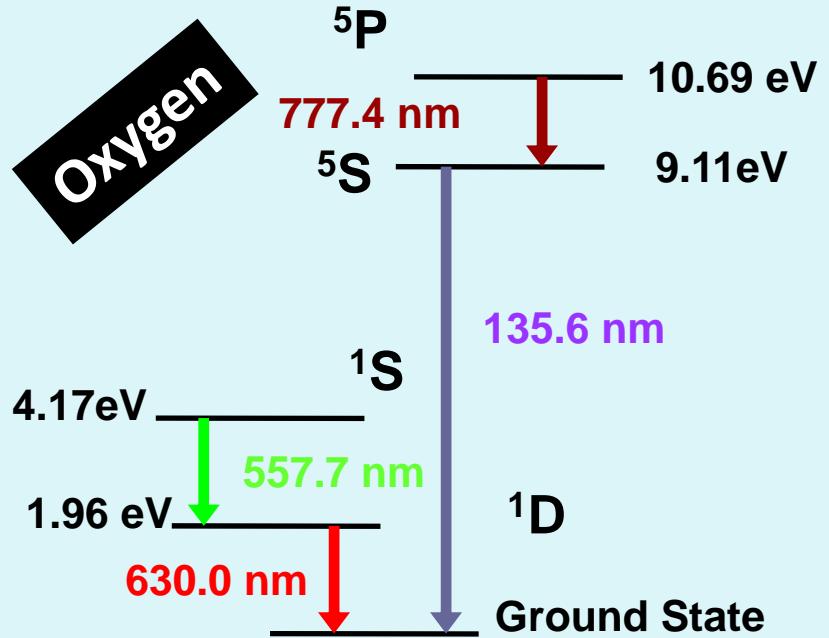
- ✓ HF Modification Experiments at HAARP
- ✓ Parametric instabilities
- ✓ Electron acceleration



Optical Emissions

Emissions stimulated by impact of energetic electrons on ambient species (N_2 , O , O_2)

➤ Excitation energy an indicator of electron energy spectrum



Green: O(1S)	557.7 nm	~5 eV
Blue: N ₂ (2PG)	~400.0 nm;	11 eV
	N ₂ +(1NG)	427.8 nm; 20 eV
Red: N ₂ (1PG)	676.5 nm;	7.4 eV
	O(1D)	630.0 nm; ~2 eV



Ordinary (collisional) Aurora

Precipitating (*primary*) electrons excite & ionize neutral particles via collisions

Energy dissipation rate
(Bethe's formula)

$$d\mathcal{E}_b / dz \cong -\mathcal{E}_b / l_b$$

$$l_b [\text{km}] \simeq 5 \frac{N_{[120 \text{ km}]} }{N} \sqrt{\mathcal{E}_b [\text{keV}]}$$

Flux of suprathermal
(secondary) electrons

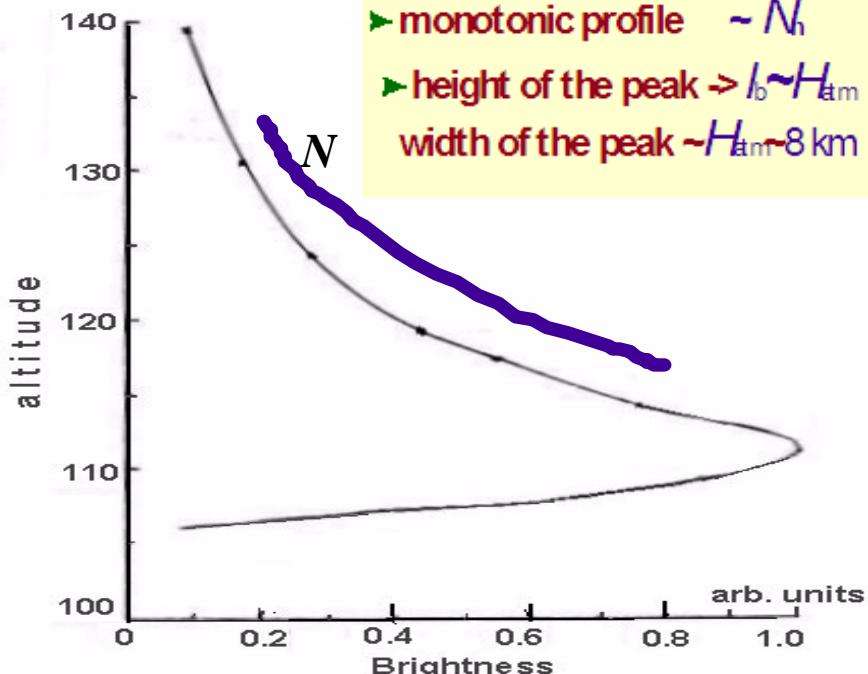
defines brightness and colors
of auroral glow

$$\Phi(\varepsilon) \propto \varepsilon^{-3.5}$$

~6 to ~300 eV



Auroral Ray Altitude Profile



$Q_\lambda = N_j \int \sigma_\lambda(\varepsilon) \Phi(\varepsilon, \vartheta) d\Omega d\varepsilon$

Excitation rate

neutral density

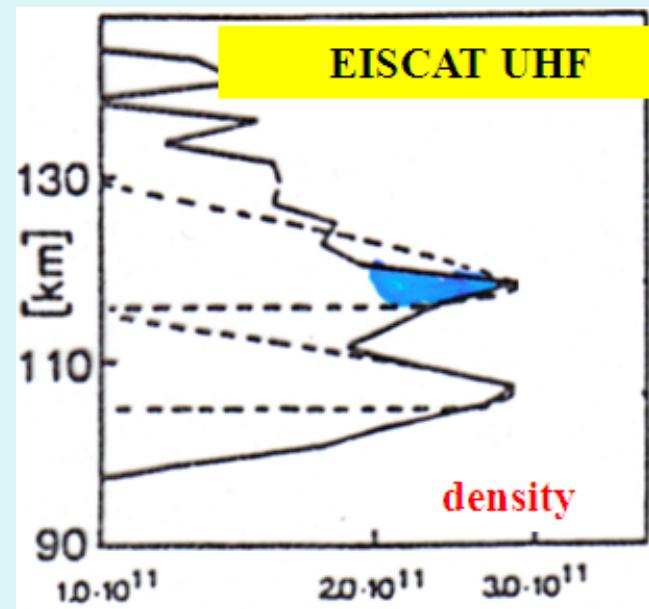
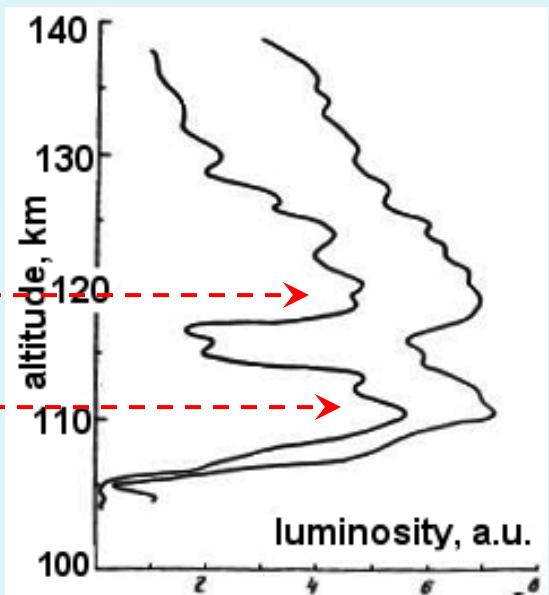
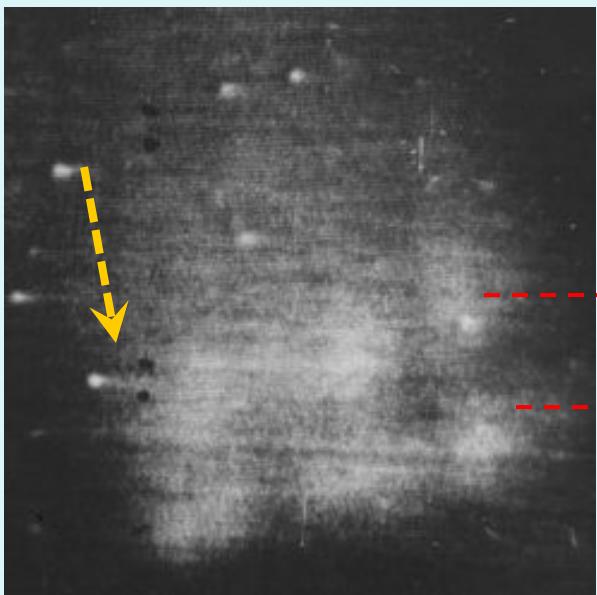
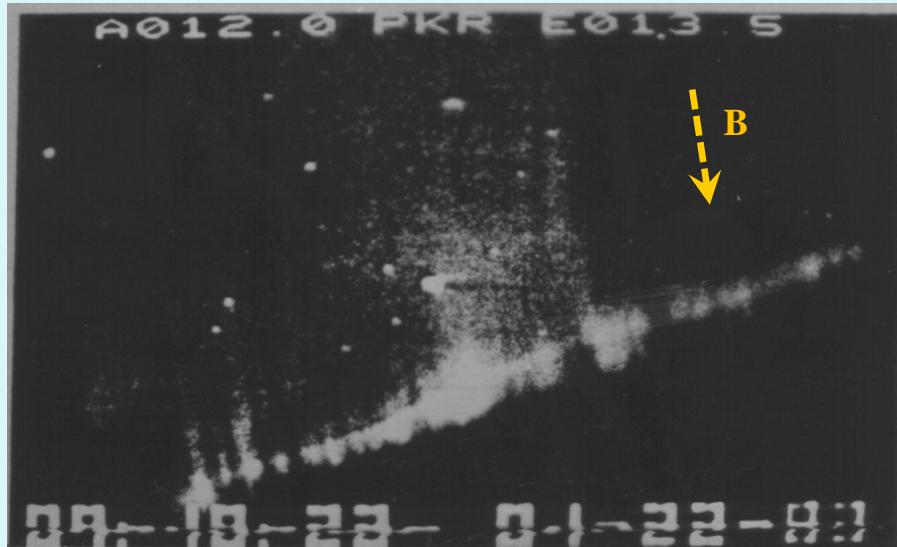
excitation cross-section

electron flux



“Enhanced” Aurora

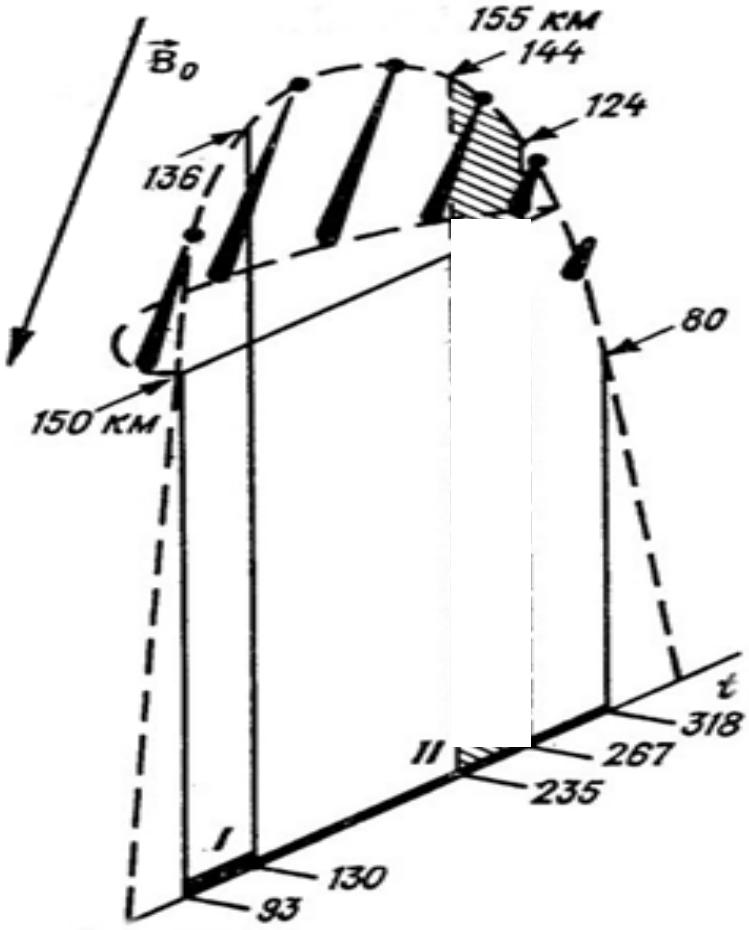
Sharp -gradient
or double-peaked
profiles



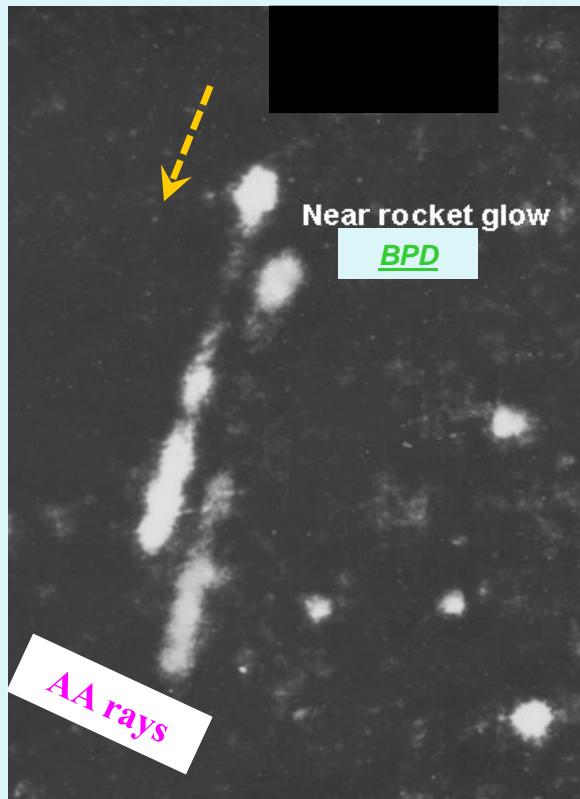


AA (active) Experiments

Zarnitsa-2 experiment (1975)

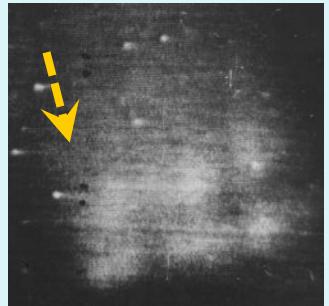
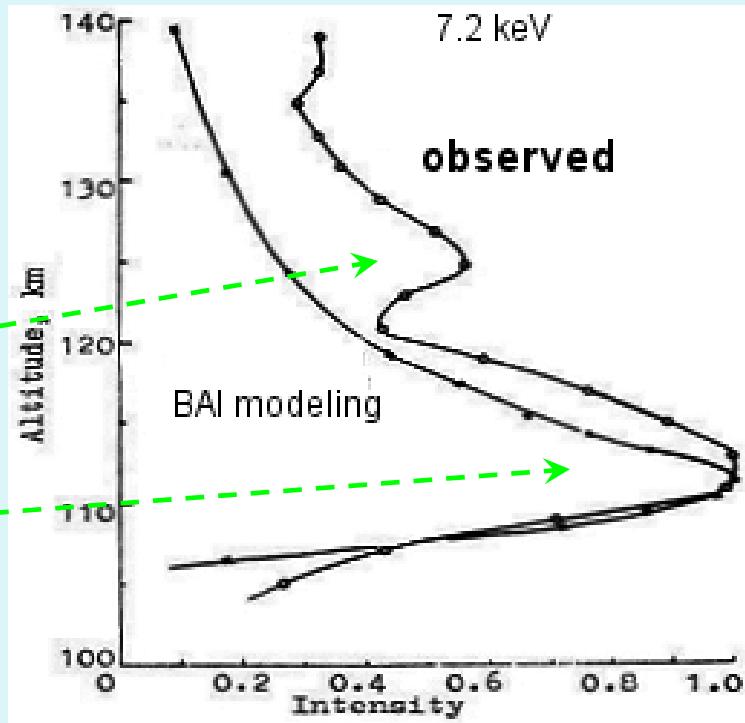
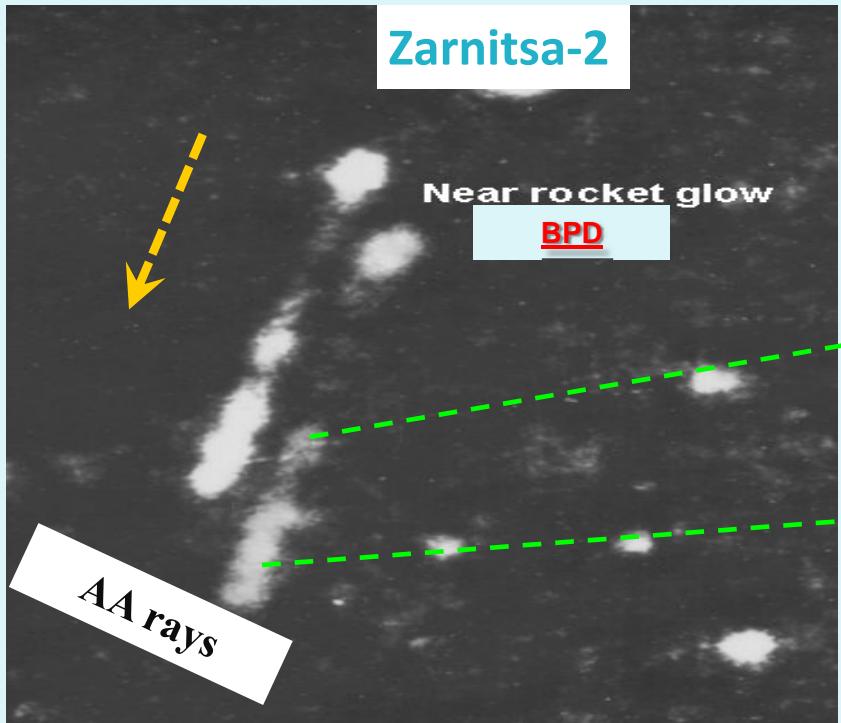


Low-light TV Observations

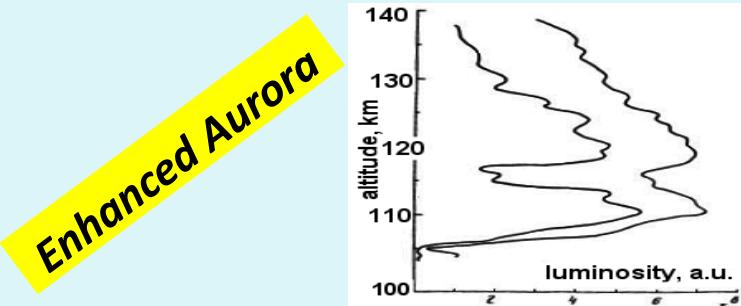




Artificial Aurora Rays



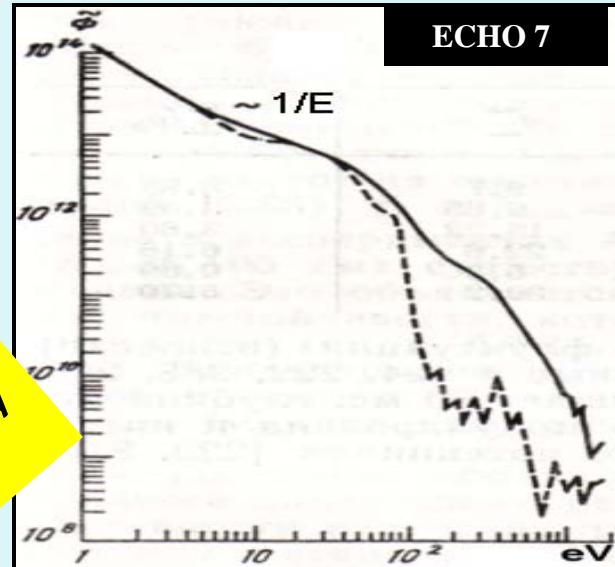
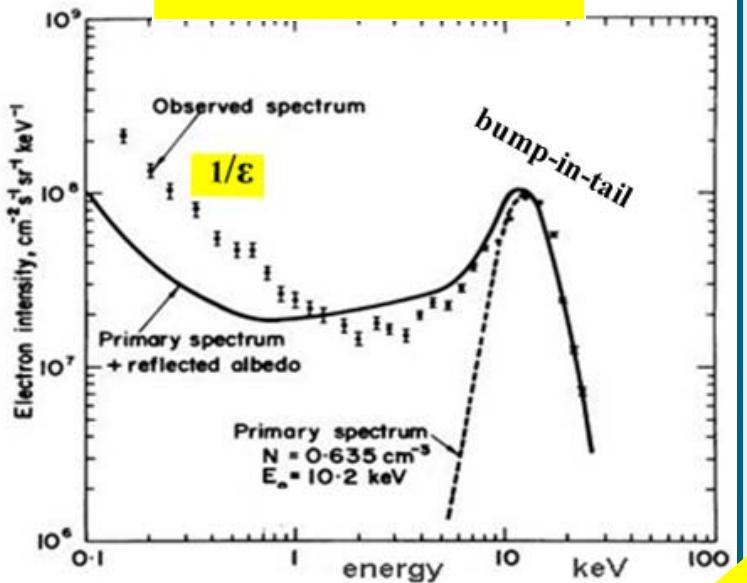
Mishin et al., 1981



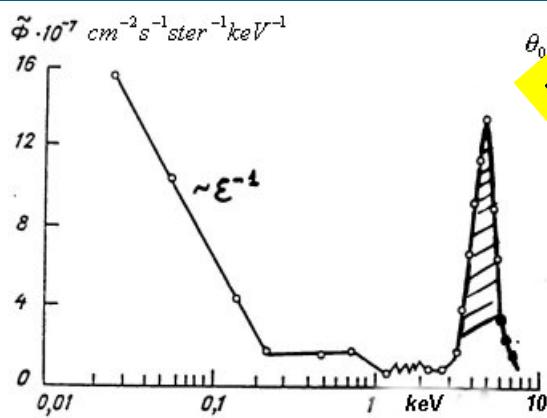


Suprothermal Electrons

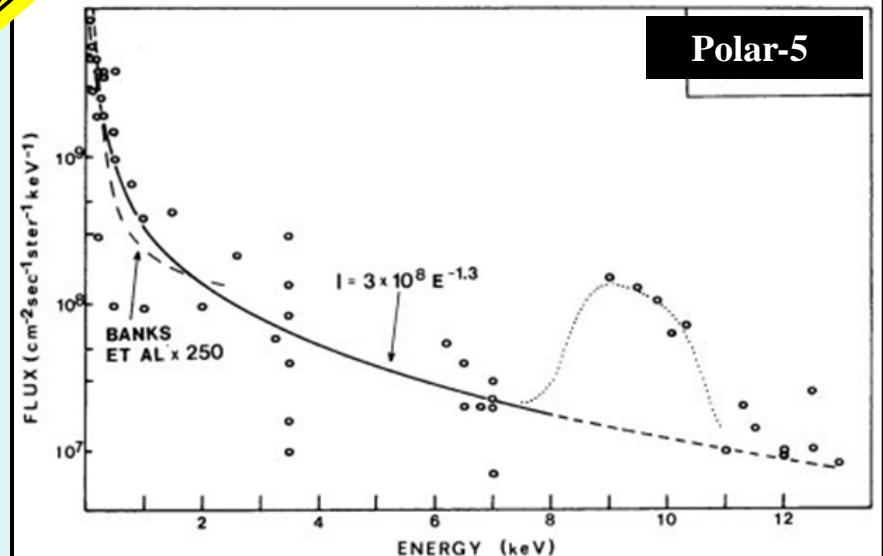
Natural Auroras



Flat suprothermal spectra indicate electron acceleration by Strong Langmuir Turbulence

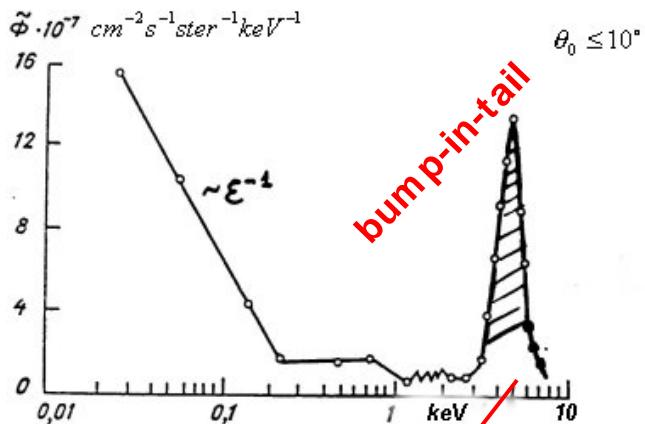


Beam of field-aligned electrons over the II class arc
[Arnoldy et al., 1974]





Beam-Plasma Instability

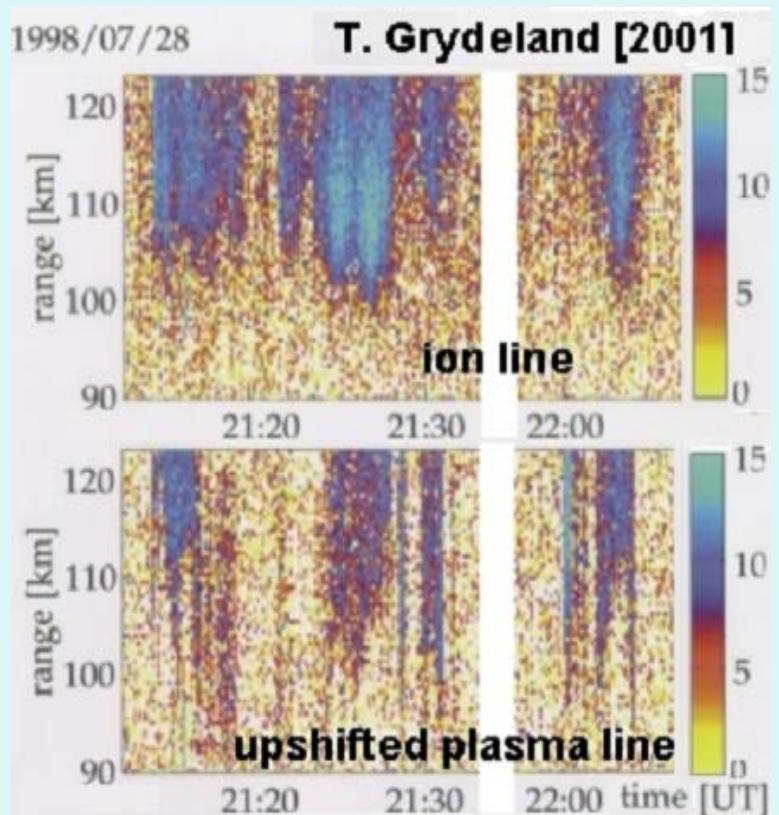


Beam of field-aligned electrons over the II class arc
[Arnoldy et al., 1974]

Inverse Landau damping

$$\Gamma_b \sim \omega_p \frac{\pi n_b}{n_e} \left(\frac{\varepsilon_b}{\Delta \varepsilon_{||}} \right)^2 - \nu_e(T_e)$$

EISCAT UHF ISR



Langmuir waves grow at altitudes $> \sim 110 \text{ km}$



Strong LT in auroral plasma



$$n_b^{(th)} \approx 10^{-6} \left(1 + \frac{2\nu_e \varepsilon_b}{3\omega_p T_e} \right) n_e$$

SLT

Flat accelerated electron spectrum

$$F_a(\varepsilon_{||}) \simeq \frac{2p_a - 1}{v_{\min}} n_a \left(\frac{\varepsilon_{\min}}{\varepsilon_{||}} \right)^{p_a}$$

($p_a \simeq 0.8-1$)

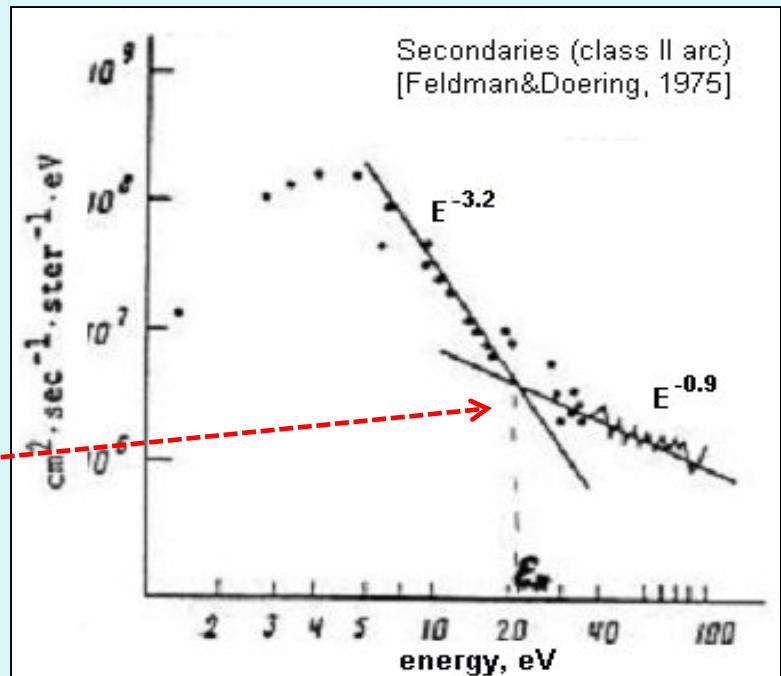
Joining condition

$$F_a(\varepsilon_{\min}) = F_0(\varepsilon_{\min})$$

➤ Acceleration of **secondary** electrons

$$\varepsilon_{\min} \simeq 30(W_L/n_s T_e)^{-2/5} [\text{eV}]$$

n_s is the density of secondary electrons





➤ Effects of collisions on SLT

$$\Gamma_b \gg \nu_e > \frac{m}{M} \omega_p$$

As the collapse rate is smaller than Γ_b , the beam can excite waves *but* the trapped waves are damped faster than collapsing. As nonlinear transfer is reduced, the Langmuir wave energy grow until collapse will be possible.

the limiting collision frequency

$$\nu_* = \omega_p \left(\frac{m}{M} \frac{\Gamma_b}{\omega_p} \right)^{1/2}$$

Wave energy density

$$W_L/n_e T_e \simeq \frac{3M}{m} \left(\frac{\nu_e}{\omega_p} \right)^2$$

ionization by
accelerated electrons

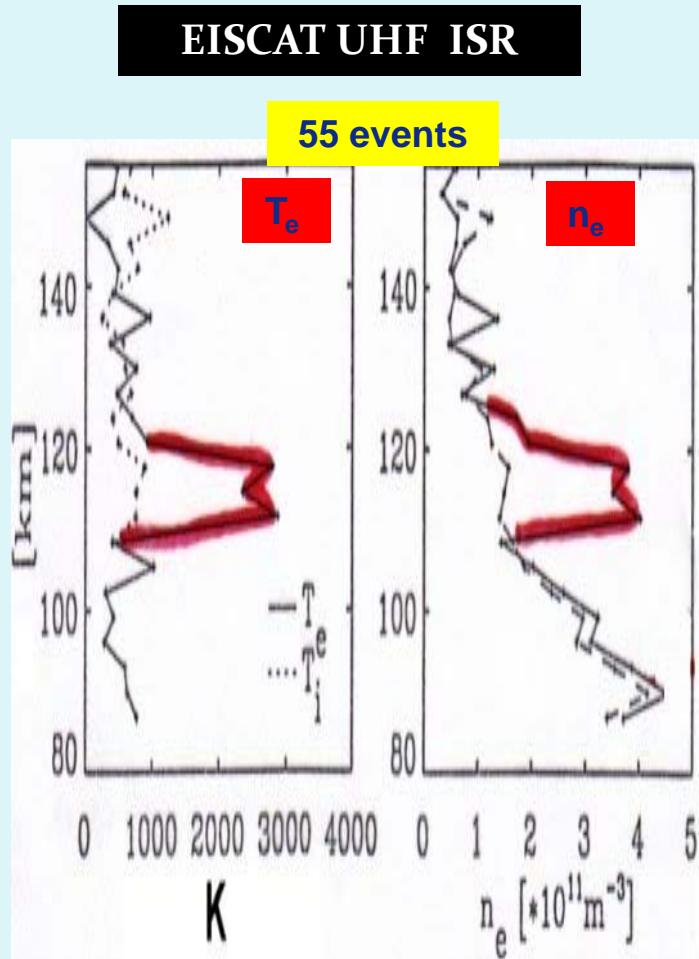
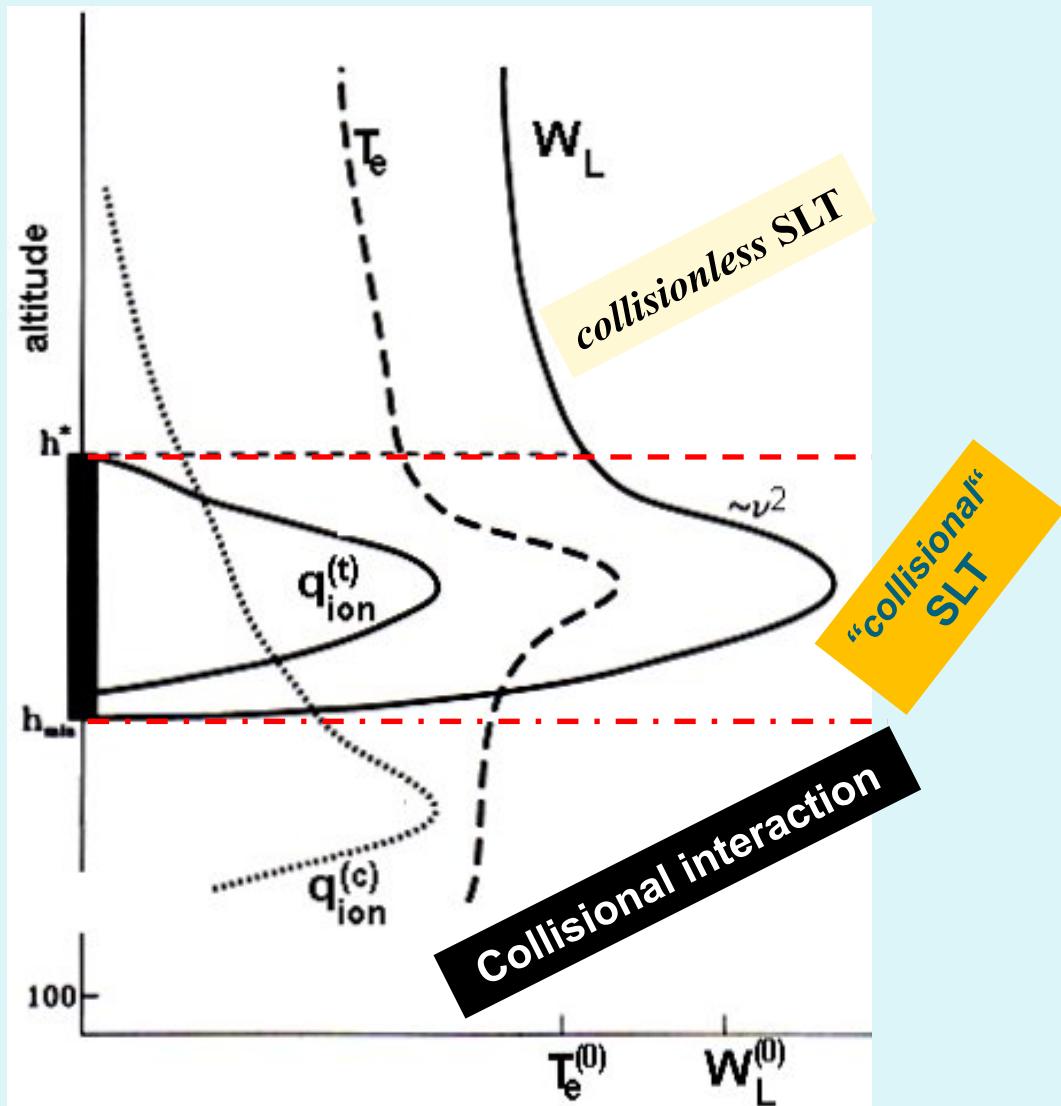
$$T_e \approx \sqrt{3\varepsilon_{ion} \frac{W_L}{n_e}}$$

$$q_{ion}^{(t)} \simeq 10 \nu_e(T_e) n_b \frac{T_e}{\varepsilon_{ion}} \left(\frac{\varepsilon_b}{\Delta\varepsilon_{||}} \right)^2$$



Plasma Turbulence Layer

Schematic of altitude-profiles





High Frequency Active Auroral Research Program

HAARP Research Station, Gakona (62.4 N, 145 W)

<http://haarp.alaska.edu>

- 180-element phased HF antenna array
- 3.6 MW radiated power → 1.3 GW ERP
- beam pointing +/-30 deg off vertical

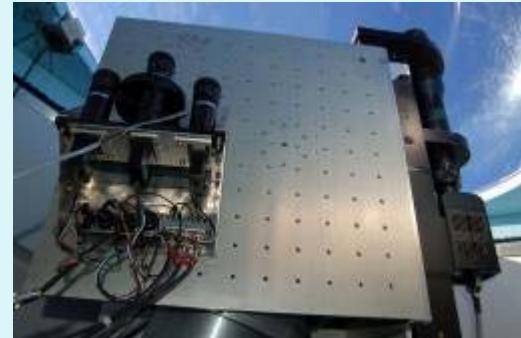
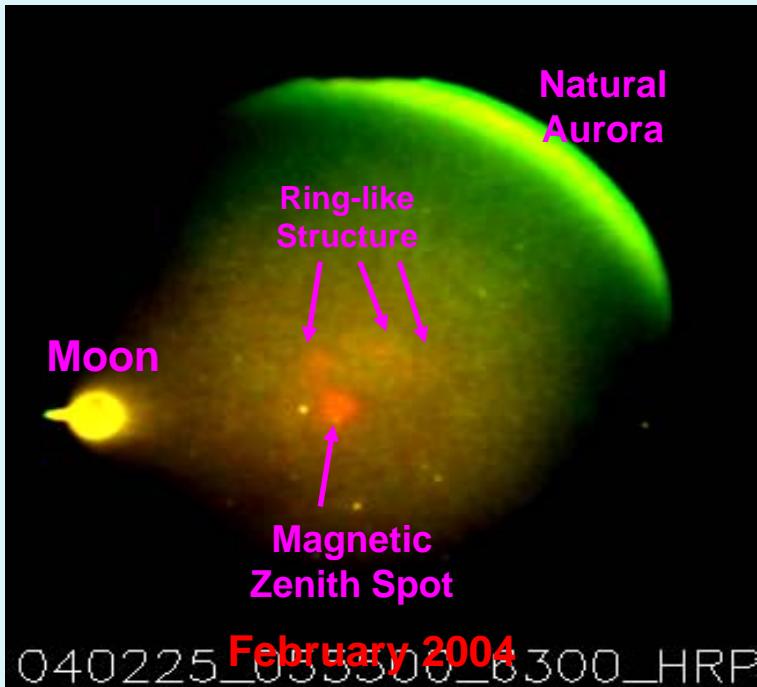


frequency : 446 MHz peak power : 512 kW
aperture : 219.8 m² (16 panels)
beam width : 4° (north-south), 2.5° (east-west)



HAARP Optical Diagnostics

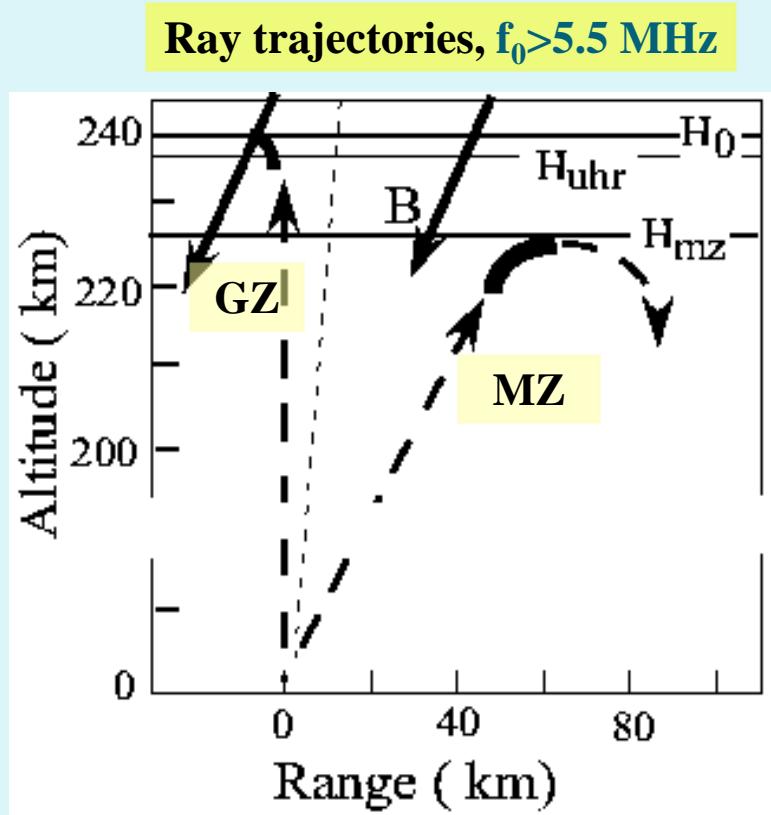
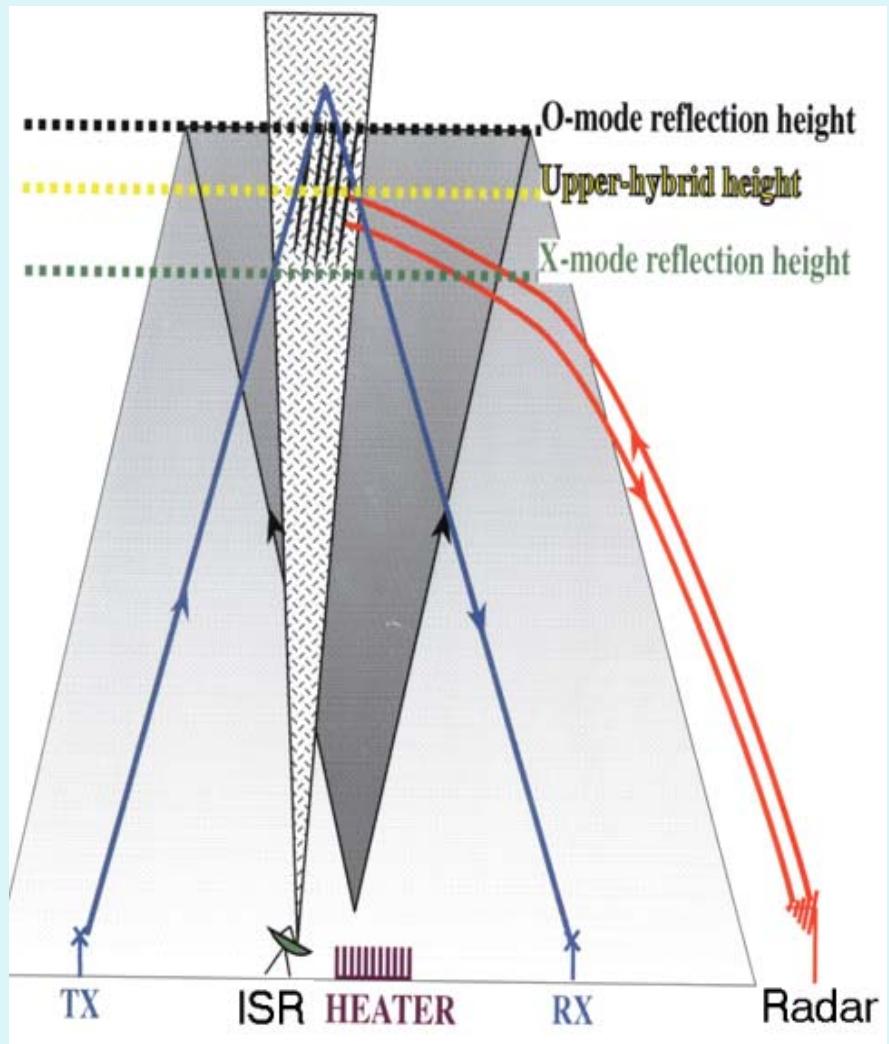
- 2 optical shelters
- 6" telescope up and running in open-air 14' dome; utilizing same CCD cameras
- Imager and photometer mounted on hydraulic lift and motorized stage under 5' clear dome
- Photometer electronics upgraded: 3 channels with one filter wheel
- 3.5" Optical imager (bare CCD)
- 4-channel all-sky low-light webcam operates year-round (except June)



Pedersen et al., 2008



Dependence on HF-beam Pointing Angle



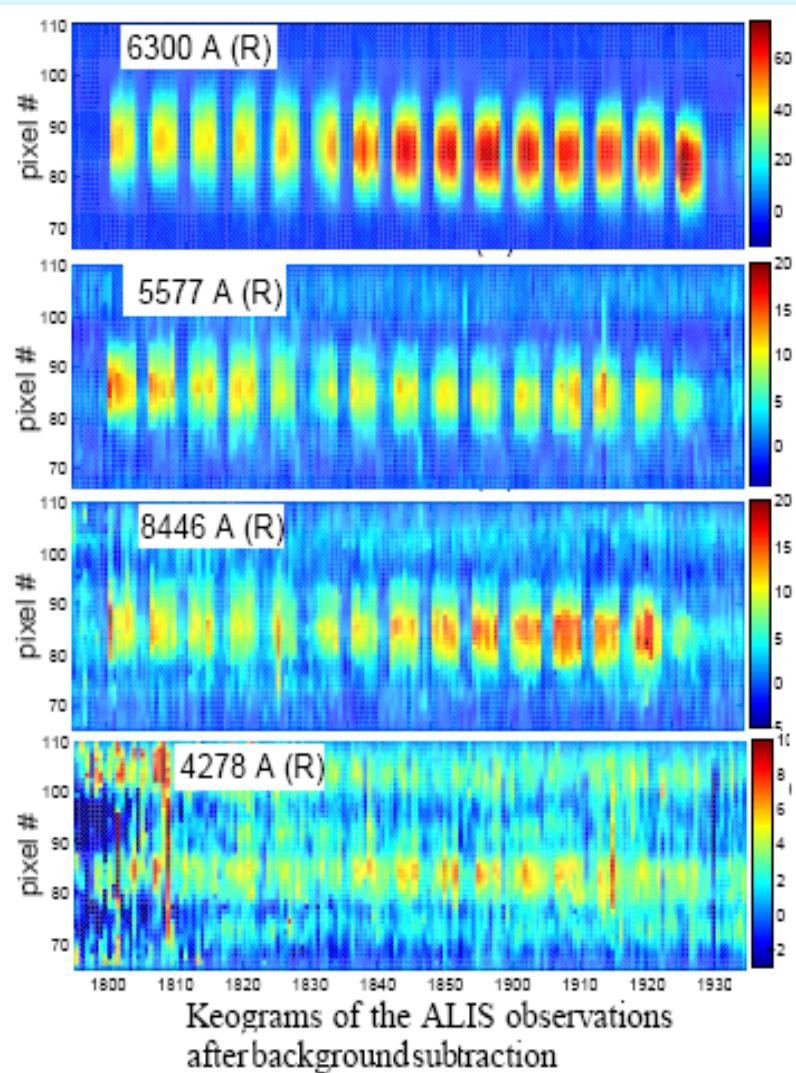
$$f_P(H_0) = f_0$$

$$f_{\text{UH}}(H_{\text{uhr}}) = f_0$$

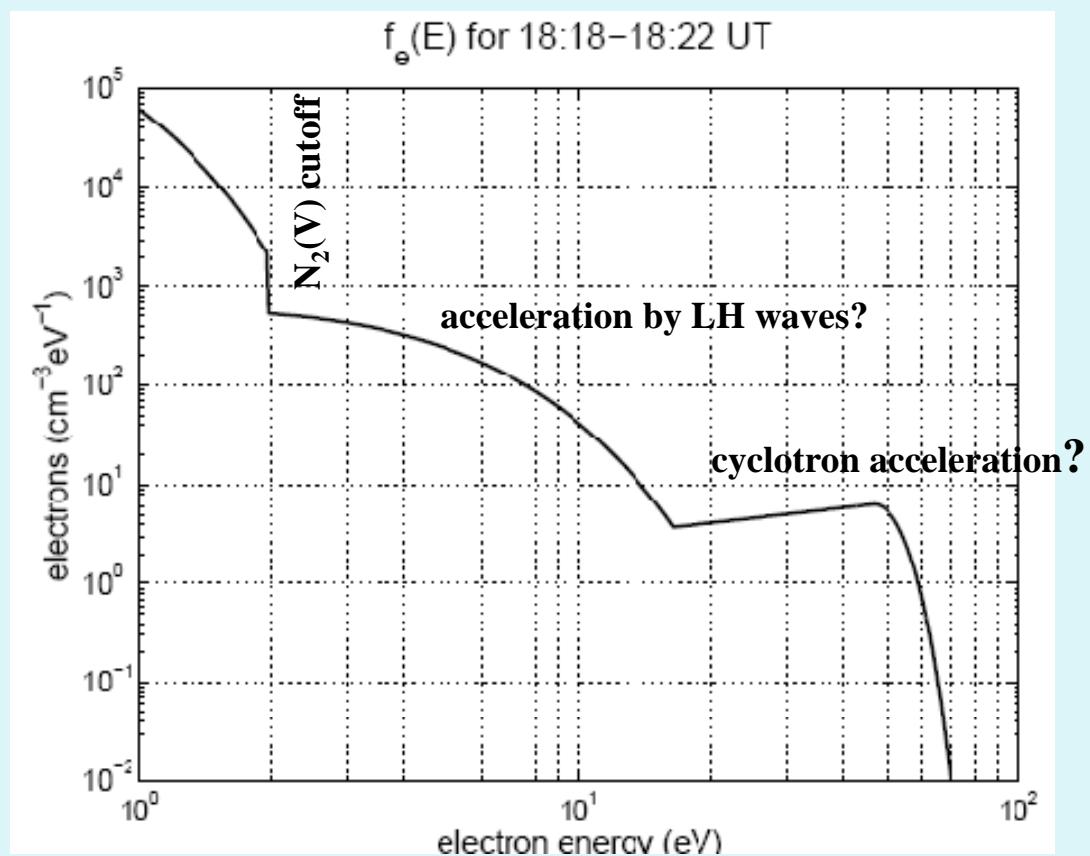
$$f_P(H_{\text{mz}}) = f_0 \cos \chi$$



Accelerated Electrons



- Tromsø: 5.423 MHz, 375 MW, MZ, 4/2-min on/off.
- Te enhancements 3000-3500 K

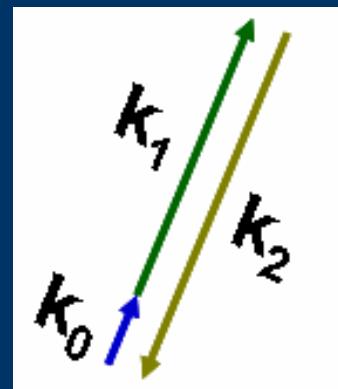




Parametric Decay Instability

$$\mathbf{EM}_0 \rightarrow \mathbf{EP}_1 + \mathbf{EP}_2$$

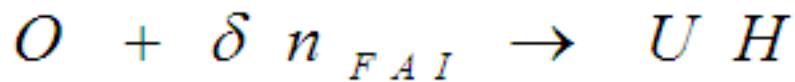
- Electromagnetic Pump Wave \mathbf{EM}_0
- Daughter HF (\mathbf{EP}_1) and LF (\mathbf{EP}_2) waves
 - ✓ $\mathbf{EP}_1 = \text{Langmuir and Upper Hybrid/Electron Bernstein}$
 - ✓ $\mathbf{EP}_2 = \text{Ion Acoustic and Lower Hybrid}$
- Matching Conditions
 - $\mathbf{k}_0 = \mathbf{k}_1 + \mathbf{k}_2$
 - $\omega_0 = \omega_1 + \omega_2$





Thermal Parametric Instability

Conversion of the pump wave on
field-aligned density irregularities



For $\Delta n \rightarrow 0$ and $\left| \rho_{uh}^2 \frac{d^2}{dx^2} \right| \ll |\epsilon_{\perp 0}|$

near H_{UHR}

$$\delta E_{uh} \simeq \frac{p \cdot E_0}{\epsilon_{\perp 0}} \cdot \frac{\Delta n(x)}{n_0}$$

Here $\rho_{uh} = r_D \sqrt{\frac{3}{1 - (2\Omega_c/\omega_o)^2}}$

$$\epsilon_{\perp 0} = 1 - \omega_p^2 / (\omega_0^2 - \Omega_c^2) \leq 0 \text{ above/below } H_{uhr}$$

$\epsilon_{\perp 0} < 0$ above H_{uhr} and hence $\delta E_{uh} > 0$ if $\Delta n < 0 \rightarrow$
excess heating in density rarefactions and deficit heating in compressions.

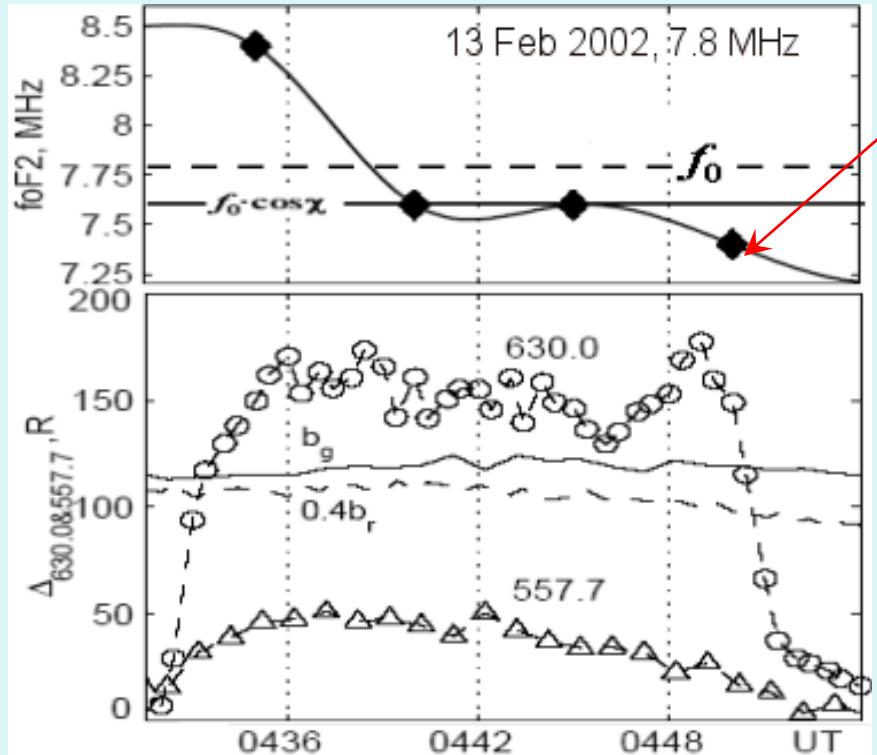
Positive feedback:

UH -trapping *inside striations* → enhanced heating → further depletion

The growth time for the TPI is of order seconds

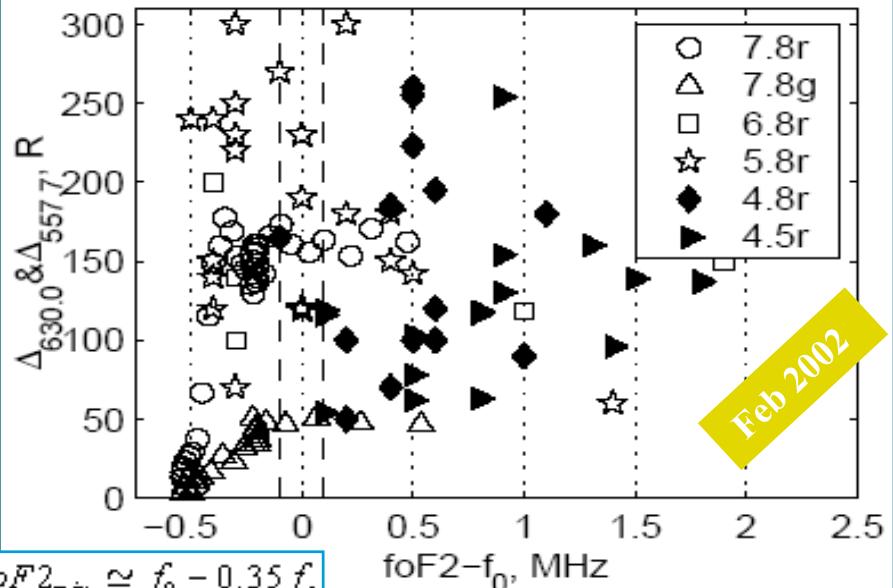


Airglow at MZ in an underdense ionosphere



continuous injection

Fig. 6. Summary plot of the variation of the HF-induced airglow at MZ with $f_{0F2} - f_0$. Here 7.8, ... 4.5 stand for $f_0 = 7.8, \dots, 4.5$ MHz;



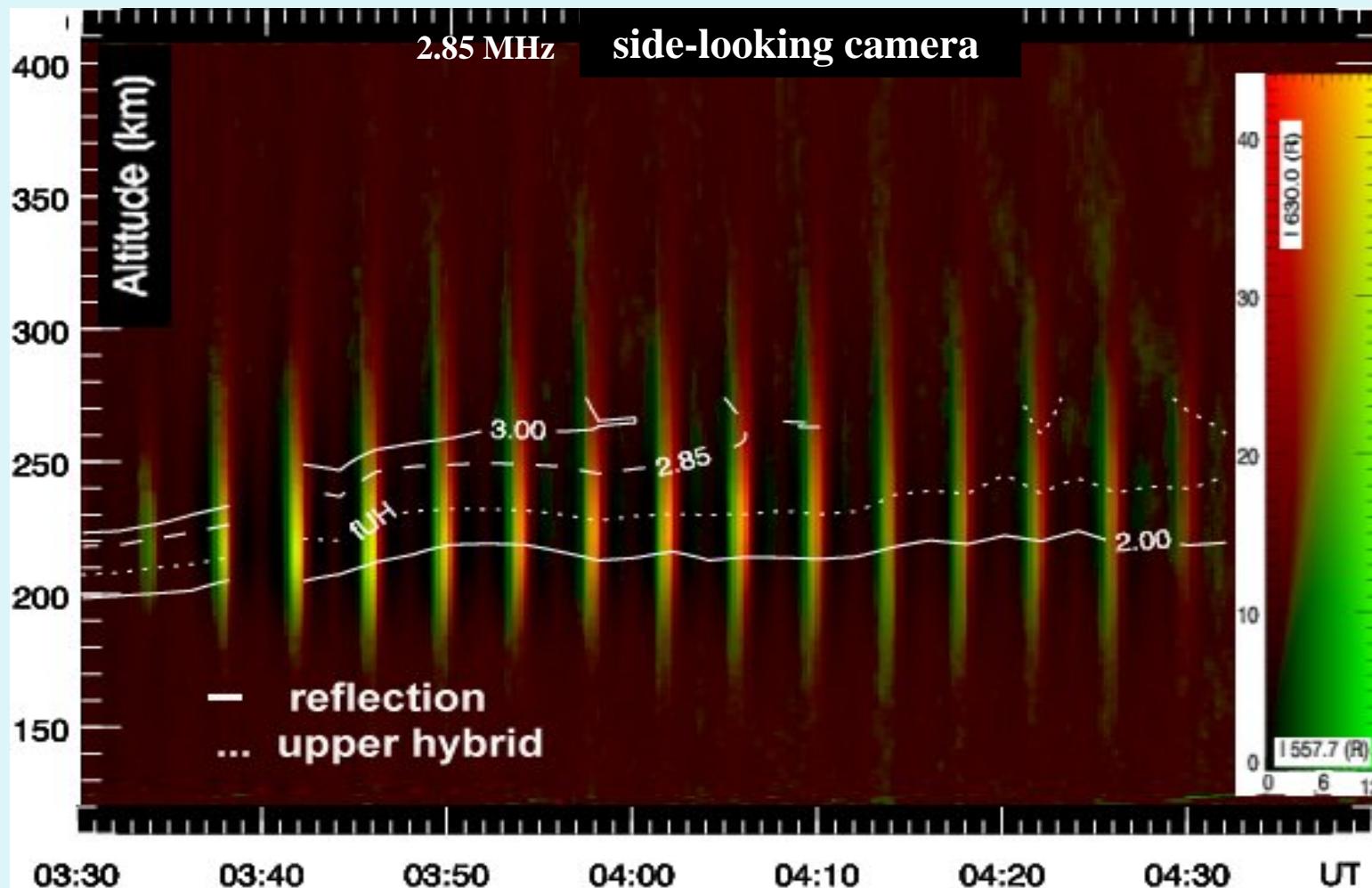
$$O \rightarrow UH + LH$$

$$x_{uh}^*(H, f_0) \simeq \frac{2}{3} \frac{s(s^2 - 4)}{s^2 - 1} \frac{f_0 - f_{uh}(H)}{f_c(H)}$$

Mishin et al., 2005

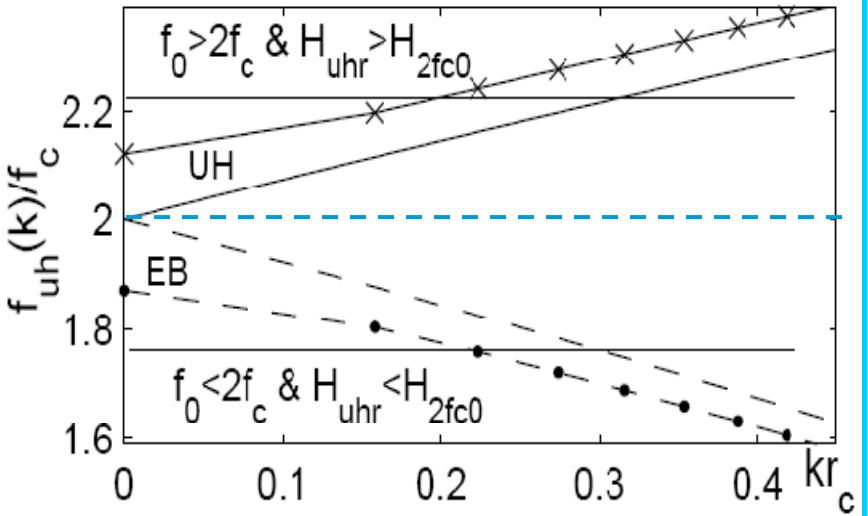


Underdense F-region





Second electron gyroharmonic



O-mode reflection at and above $H_{2fc0} \rightarrow$
TPI & PDI_{EB} above & below H_{2fc0}

Suppression of PDI_L by striations

O-mode reflection below $H_{2fc0} \rightarrow$
PDI_{EB} & PDI_L

The TPI_{UH} threshold near the second GH strongly depends on the frequency mismatch $\delta_{\text{UHR}} = \frac{f_{\text{UHR}} - 2f_c}{2f_c}$

The TPI threshold

Grach, 1979

$$\frac{E_o^2}{4\pi n T_e} > \frac{3\nu_e}{\omega_o} \left(\lambda_{\parallel}^2 l_e^2 + \lambda_{\perp}^2 r_c^2 \right) \frac{\partial(\omega^2 \epsilon_{\perp}(\omega))}{\partial \omega^2}$$

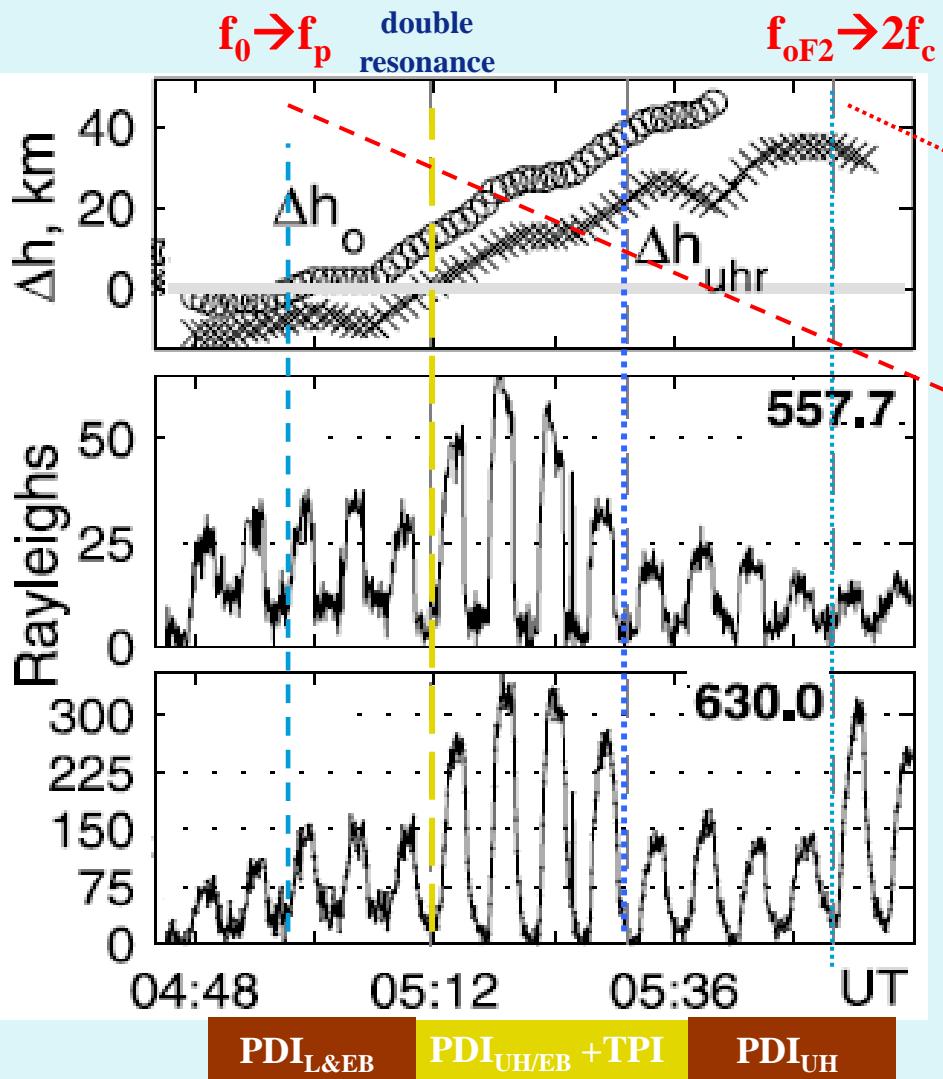
$\approx 16/3$ for $|\delta_{\text{UHR}}| < \frac{3}{4} \kappa_{\perp} r_c$

$$\delta_{\text{UHR}} < \delta_{\text{UHR}}^{(m)} \simeq 0.01$$

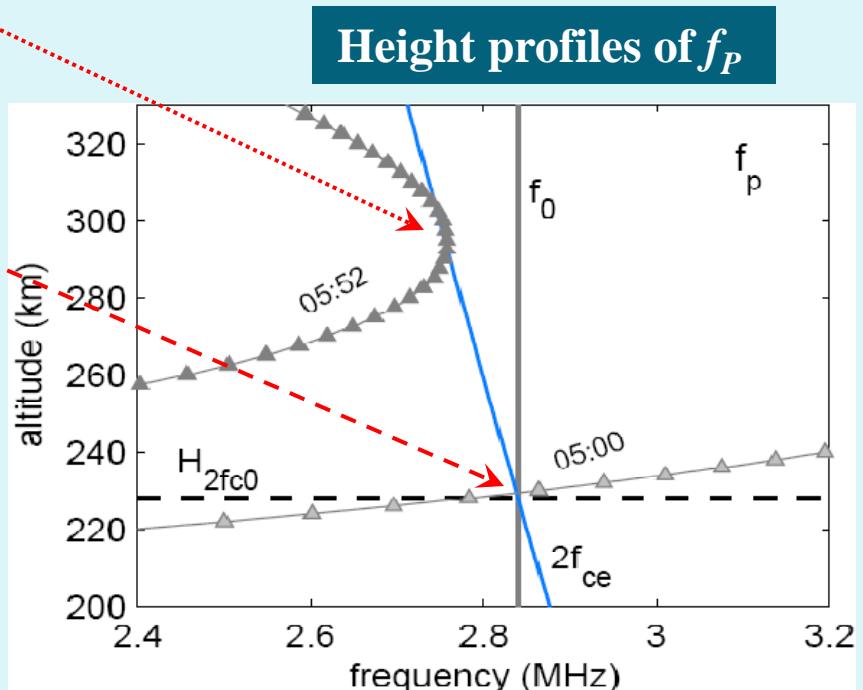
>20, otherwise



2GH, 25-Feb-2004 (2.5/2.5 min on/off, 10 MW at MZ)



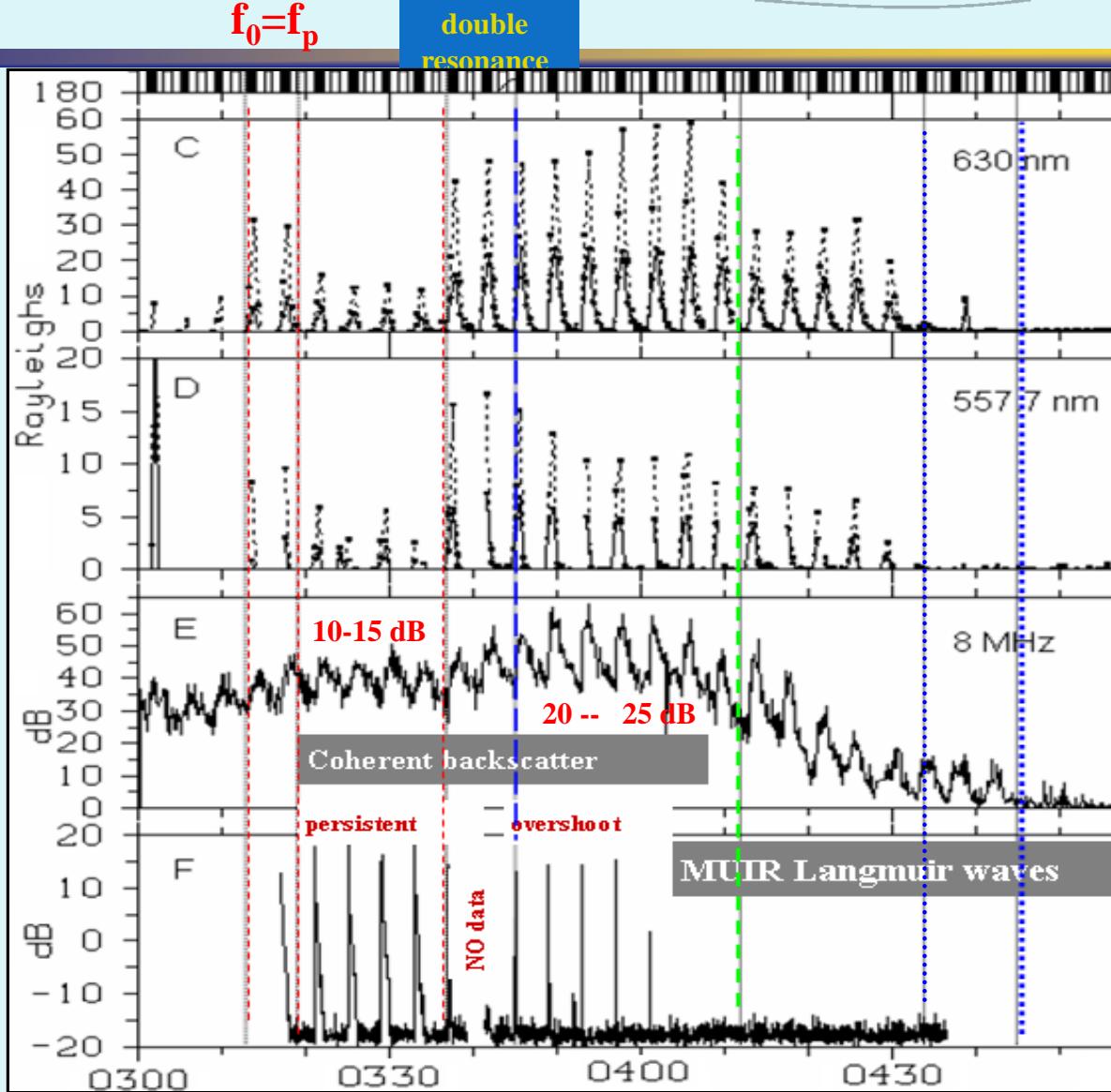
$$\delta_{uhr} < \delta_{uhr}^{(m)} \simeq 0.01$$



- Green/Red ratio > 0.1
- No change when $f_{oF2} < f_0$



2GH, 4-Feb-2005 (10 MW, 1/2 min on/off)



Kosch et al., 2007

PDI_{L&EB}

PDI_{EB&UH} + TPI

PDI_{UH}



Summary



- Common characteristics of natural and man-made auroras are flat suprathermal spectra and altitude profiles consisting of two narrow peaks displaced by ~10 km.
- These features can be explained by accounting for strong Langmuir turbulence excited by precipitating/injected electron beams in weakly-collisional ionospheric plasma.
- Up to three parametric instabilities (PDI_L , $\text{PDI}_{\text{UH}/\text{EB}}$, and TPI) act simultaneously during HF heating at the magnetic zenith.
- Optical and radar observations during a frequency pass through the second GH show the coexistence of PDI_L and $\text{PDI}_{\text{UH}/\text{EB}}$ below 2GH and of the parametric decay and thermal parametric instabilities just above 2GH.
- Airglow at MZ persists after the critical F-layer frequency drops below the pump frequency by ~0.5 MHz, in agreement with the development of the PDI_{UH} .