

Modeling the Magnetosphere with the Multi-Fluid Lyon-Fedder-Mobarry Magnetosphere Model

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Lessons from Dennis

- Use your curiosity to find interesting questions to ask, but rely on physics to find the answers
- Simulations are only meaningful when compared to observations
- Make sure your students get the chance to present their results
- Work with good people
- Eat good food!

Outline

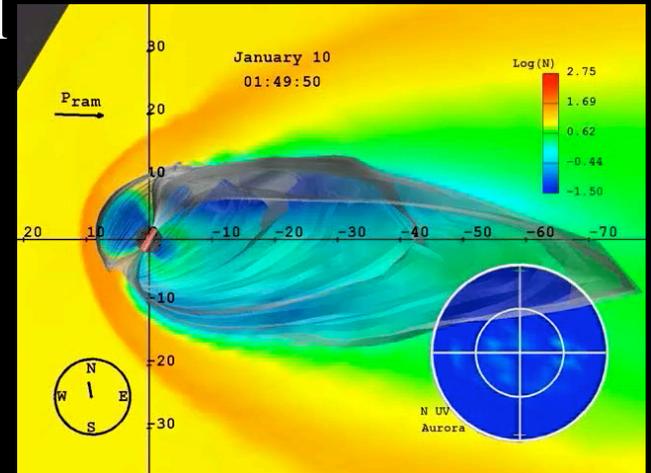
- Background and motivation
- The MultiFluid Lyon-Fedder-Mobarry Model
- Parameterized Cleft Outflow
 - Low Speed/High Flux
 - High Speed/High Flux
- Conclusions

Motivation for Including Ion Outflow

- In previous work with the LFM only electrodynamic coupling between the magnetosphere and ionosphere was considered
- Using Cluster observations *Kistler et al. (2005)* showed that O^+ can become the main component of the plasma sheet during strong magnetic storms
- *Shay and Swisdak (2004)* found that the rate of reconnection is controlled by the heavy ion species in multi-species simulations
 - They suggest this may have an impact on the size and occurrence frequency of substorms
- *Winglee et al. (2002)* found a significant reduction in the CPCP with their multifluid MHD simulation
 - The LFM is known to have a high values for the CPCP
- With the advent of the multifluid version of the LFM it is now possible to consider the effect of ionospheric O^+ on magnetospheric dynamics

Multi-Fluid LFM Magnetospheric Model

- Uses the multi-fluid MHD equations to model the interaction between the solar wind, magnetosphere, and ionosphere
 - **Computational domain**
 - $30 R_E < x < -300 R_E$ & $\pm 100 R_E$ for YZ
 - Inner radius at $2 R_E$ altitude
 - **Calculates**
 - Solves mass, parallel and perpendicular momentum, and plasma energy for each plasma species
 - Self consistent magnetic field configuration
 - **Requires**
 - Solar wind MHD state vector along outer boundary
 - Empirical model for determining energy flux of precipitating electrons
 - Cross polar cap potential pattern in high latitude region which is used to determine boundary condition on flow
 - Specification of plasma outflow from inner boundary



MFLFM Equations

$$\frac{\partial \rho_\alpha}{\partial t} = -\nabla \cdot \rho_\alpha \vec{u}_\alpha$$

$$\frac{\partial \vec{p}_\alpha}{\partial t} = -\nabla \cdot (\vec{p}_\alpha \vec{u}_\alpha + \vec{I}P_\alpha) + \vec{F}_\alpha^d + n_\alpha q_\alpha \vec{E}_\parallel$$

$$\frac{\partial \varepsilon_\alpha}{\partial t} = -\nabla \cdot \vec{u}_\alpha (\varepsilon_\alpha + P_\alpha) + \vec{u}_\alpha \cdot (\vec{F}_\alpha^d + n_\alpha q_\alpha \vec{E}_\parallel)$$

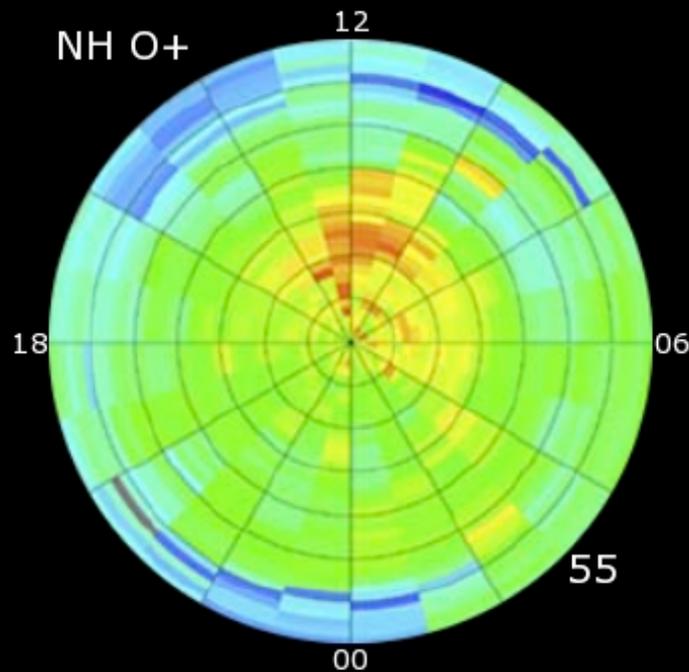
$$\frac{\partial \vec{B}}{\partial t} = -\nabla \times (\vec{u} \times \vec{B})$$

$$F_\alpha^d \equiv \hat{b} \times \left[\begin{array}{l} \vec{p}_\alpha \cdot \nabla \vec{u}_\alpha + \nabla P_\alpha + \frac{\rho_\alpha (\vec{u} - \vec{u}_\alpha) \cdot \hat{b}}{B} \frac{\partial \vec{B}}{\partial t} \\ -\frac{\rho_\alpha}{\rho} \left(\sum_\beta (\vec{p}_\beta \cdot \nabla \vec{u}_\beta + P_\beta) + \nabla P_e - \vec{j} \times \vec{B} \right) \end{array} \right] \times \hat{b}$$

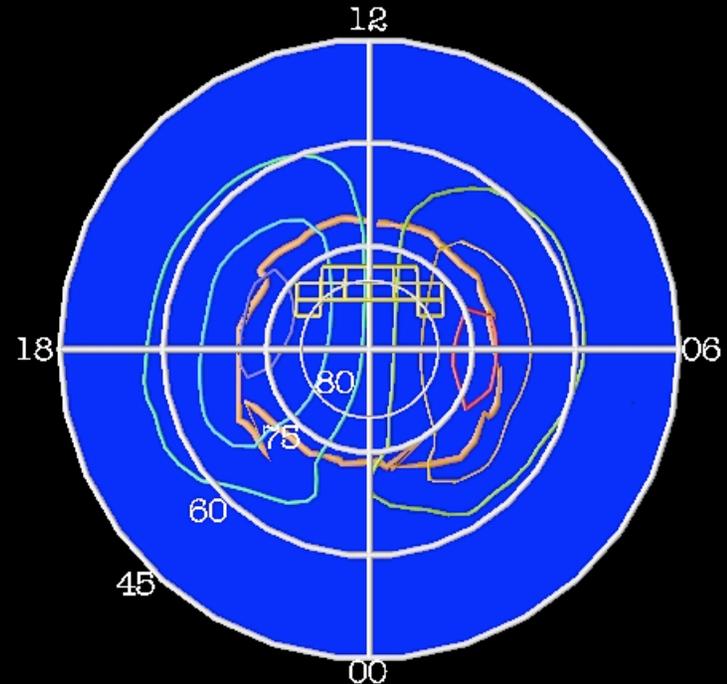
$$\vec{E}_\parallel = -\hat{b} \hat{b} \cdot \nabla P_e / ne$$

- These equations evolve the ion density, momentum, and plasma energy
 - The electrons are assumed to be inertialess
 - E_\parallel is the ambipolar electric field
- All ion species move with the $E \times B$ drift velocity in the perpendicular direction
 - The time rate of change of B in the drift force term transfers momentum between species when the magnetic field direction changes
- When the drift force is summed over species we recover the one fluid MHD equations less electron inertia

O⁺ Outflow Region



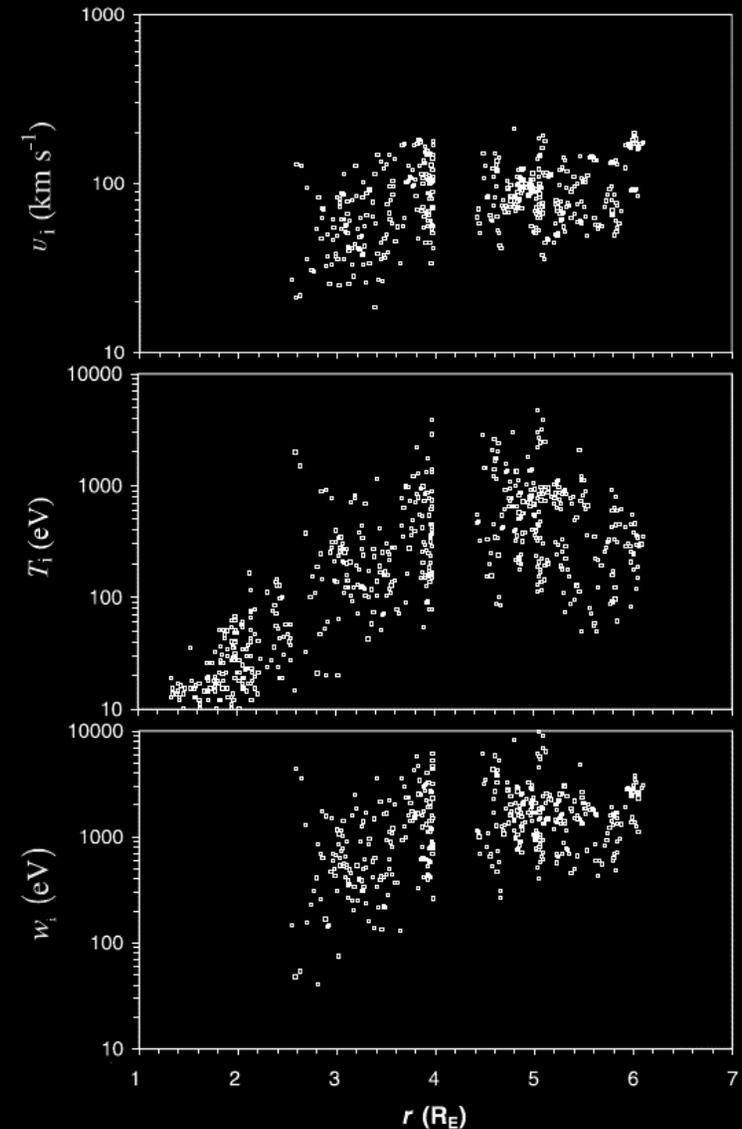
From Lennartsson et al. 2004



- *Lennartsson et al. 2004* and *Andersson et al. 2005* have statistically identified a cusp-like outflow region on the dayside
 - This region roughly corresponds to *Newell & Meng 1992* LLBL, cusp, and mantle electron precipitation region
- We turn on O⁺ in cells whose ionospheric foot points map to a similar region on the LFM ionospheric grid

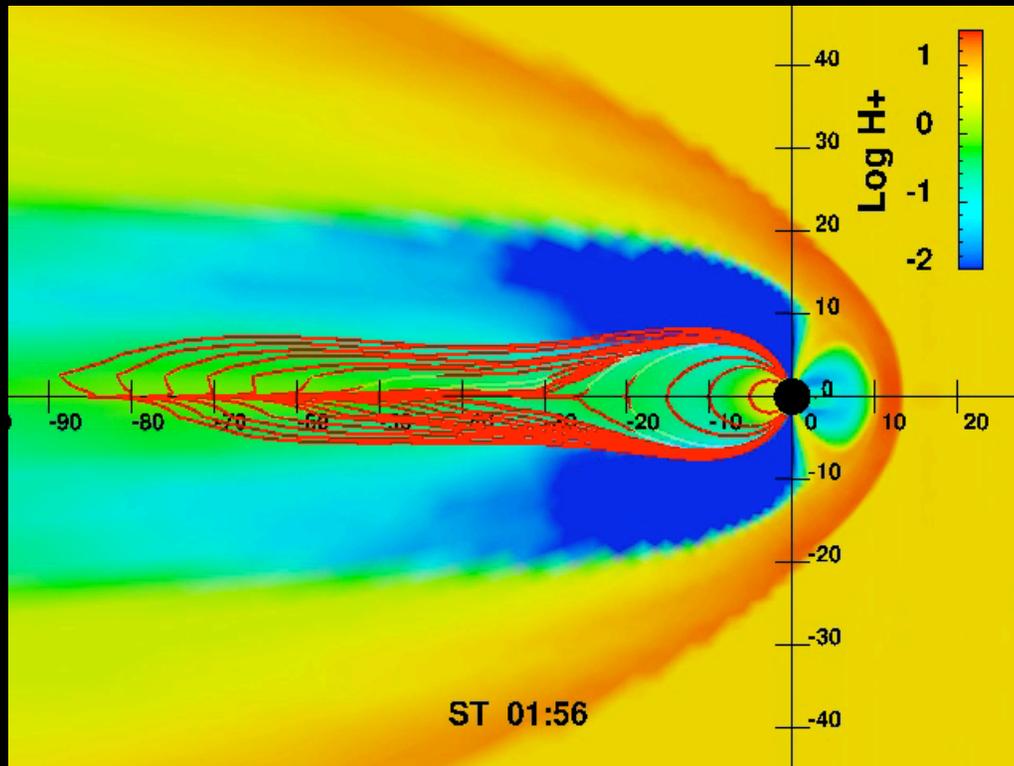
O⁺ Outflow Parameters

- *Bouhram et al. 2004* used Akebono, Interball, and Cluster observations to study the altitude profile O⁺ ions coming out of the cleft
 - Note the broad range in measurements
 - parallel velocities 20-100 [km s⁻¹]
 - Temperature 10-100 [eV]
 - Fluxes – 10⁶ [cm⁻² s⁻¹]
 - Flux and density have a solar cycle variance, but in parallel velocity and temperature it is less
- *Strangeway et al. 2005* report cusp region fluxes at FAST altitudes ranging from 10⁶ to 10⁹ [cm⁻² s⁻¹] during the Sep 24-25, 1998 geomagnetic storm
- *Nilsson et al. 2006* report fluxes ranging from 10⁷ to 10⁹ [cm⁻² s⁻¹] in a statistical study of Cluster data
- Given the broad range of measurements we have chosen fix parameters for the O⁺ coming from our cusp region
 - Flux - 10⁹ [cm⁻² s⁻¹]
 - Velocity - 20 [km s⁻¹]
 - Temperature - 10 [eV]



From Bouhram et al. 2004

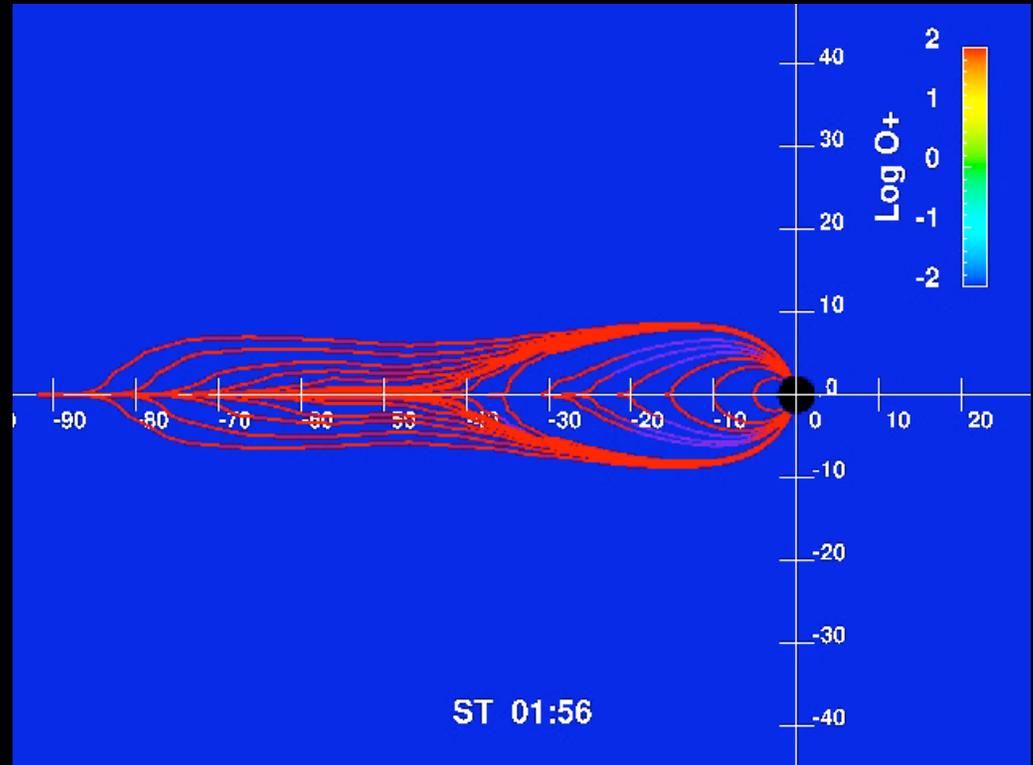
Baseline



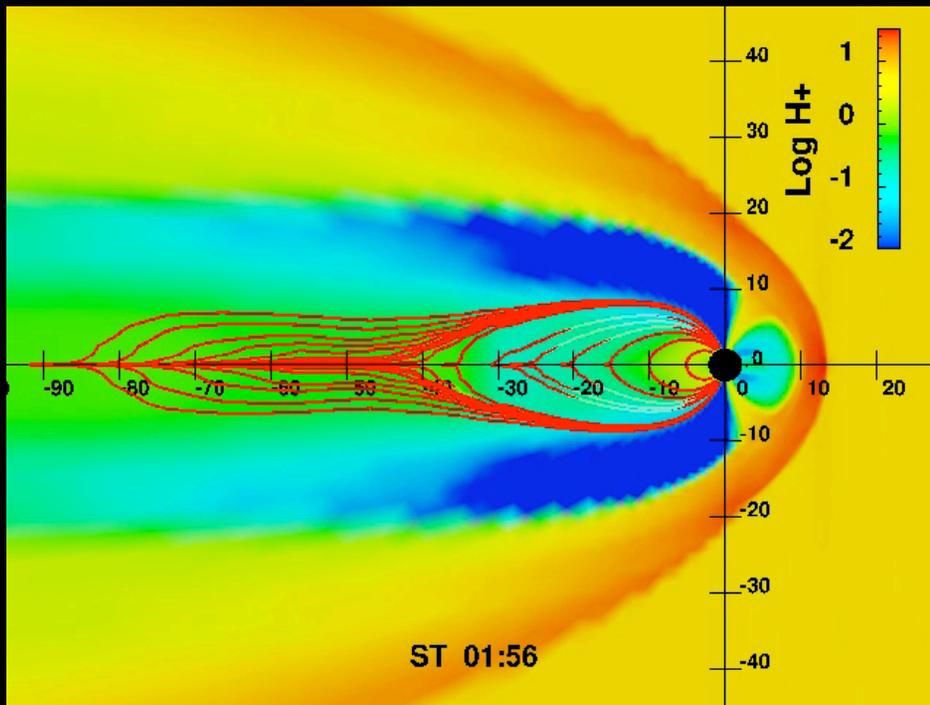
- A run of the MFLFM was completed using 99.99% H^+ in the solar wind
 - Solar wind speed is 400 km/s with a density of 5 #/cc
 - The IMF changes from +5nT at 02 ST to -5nT
- At approximately 0245 ST the simulation produces a plasmoid with reconnection site occurring at $-20 R_E$
- The simulation then enters a quasi-steady equilibrium with reconnection occurring near $-30 R_E$

Cusp-Cleft Outflow

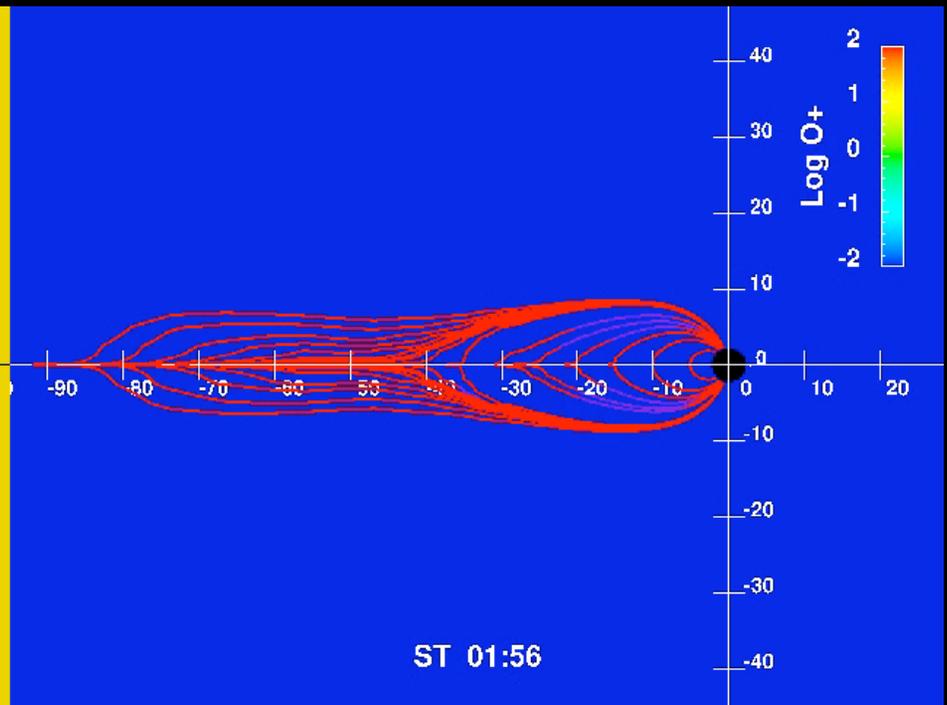
- The initial substorm at is not dramatically affected by the outflow since it does not reach the X line
- A second substorm occurs with an onset at 0530 ST with the X line occurring earthward of $-20 R_E$
- Simulation then appears to enter a period of a quasi-steady reconnection with tailward of $-20 R_E$



Comparison

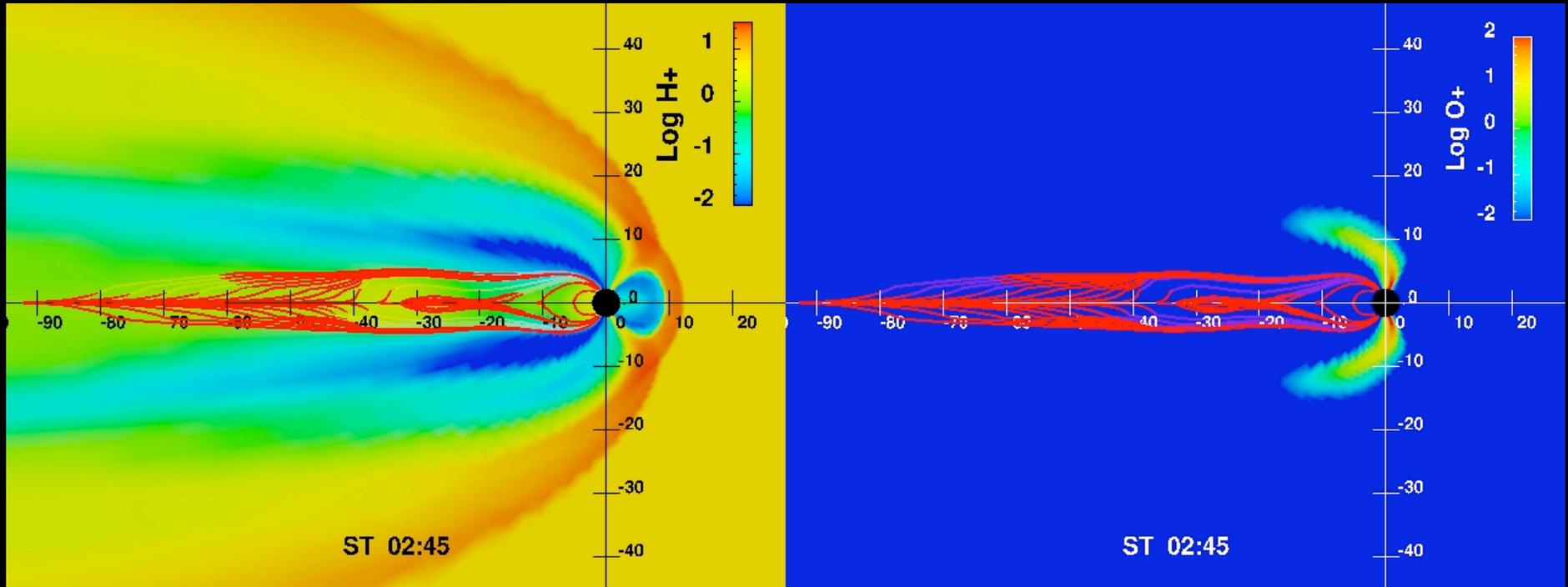


No Outflow



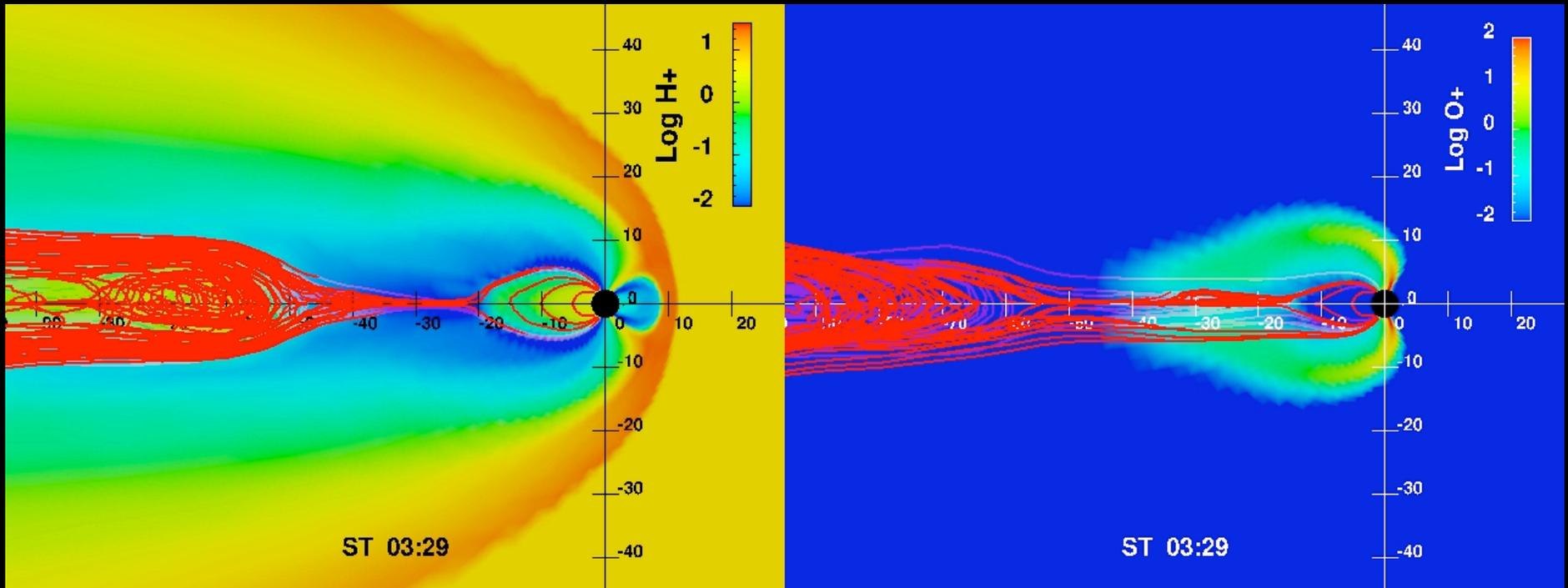
Outflow

02:45 ST



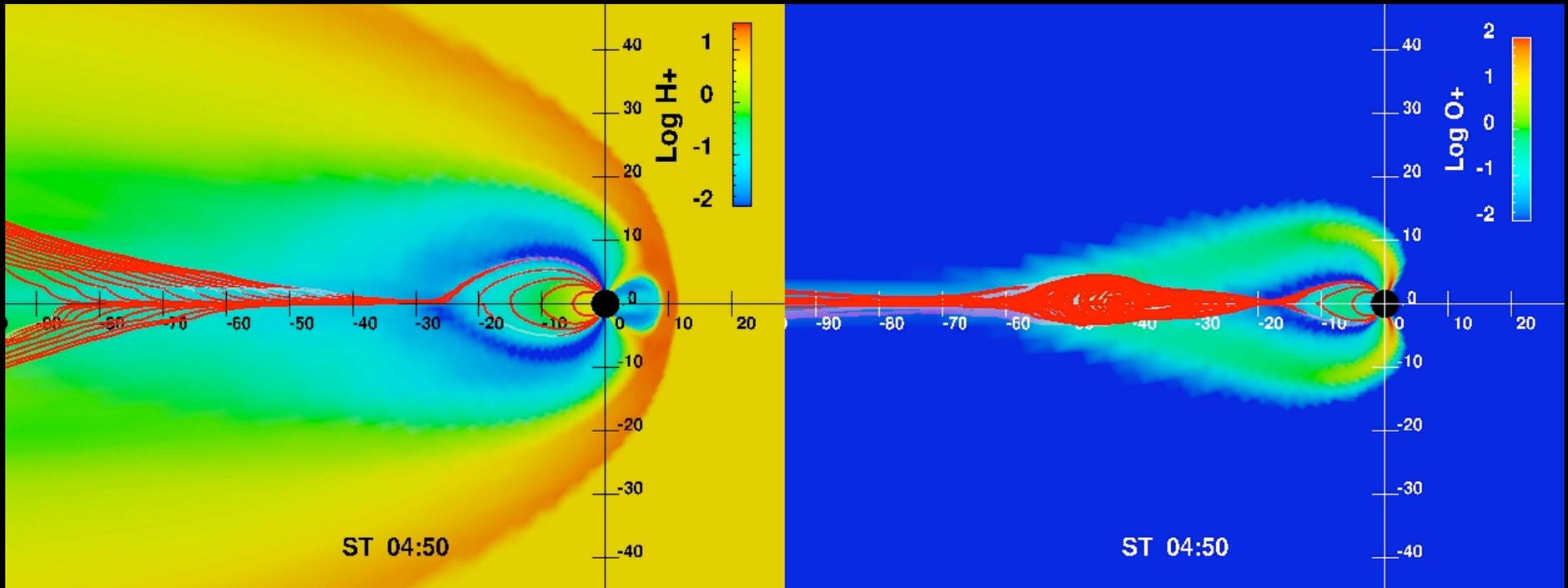
- During the initial formation of the first plasmoid the two simulations appear quite similar because the O^+ flow from the ionosphere has not propagated very far into the magnetotail

03:29 ST



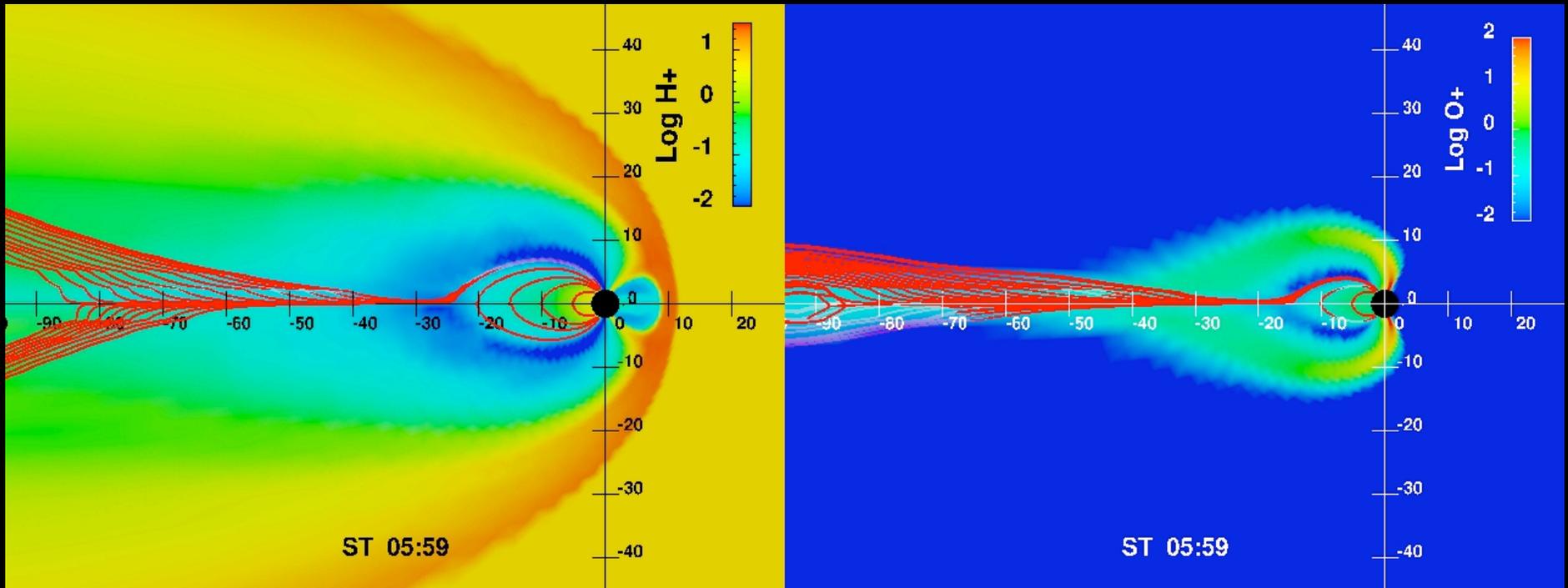
- After the release of the first plasmoid O^+ from the outflow region begins to directly enter the plasma sheet in the near Earth region
 - O^+ impact on the reconnection process has altered the shape and timing of the plasmoid

04:50 ST



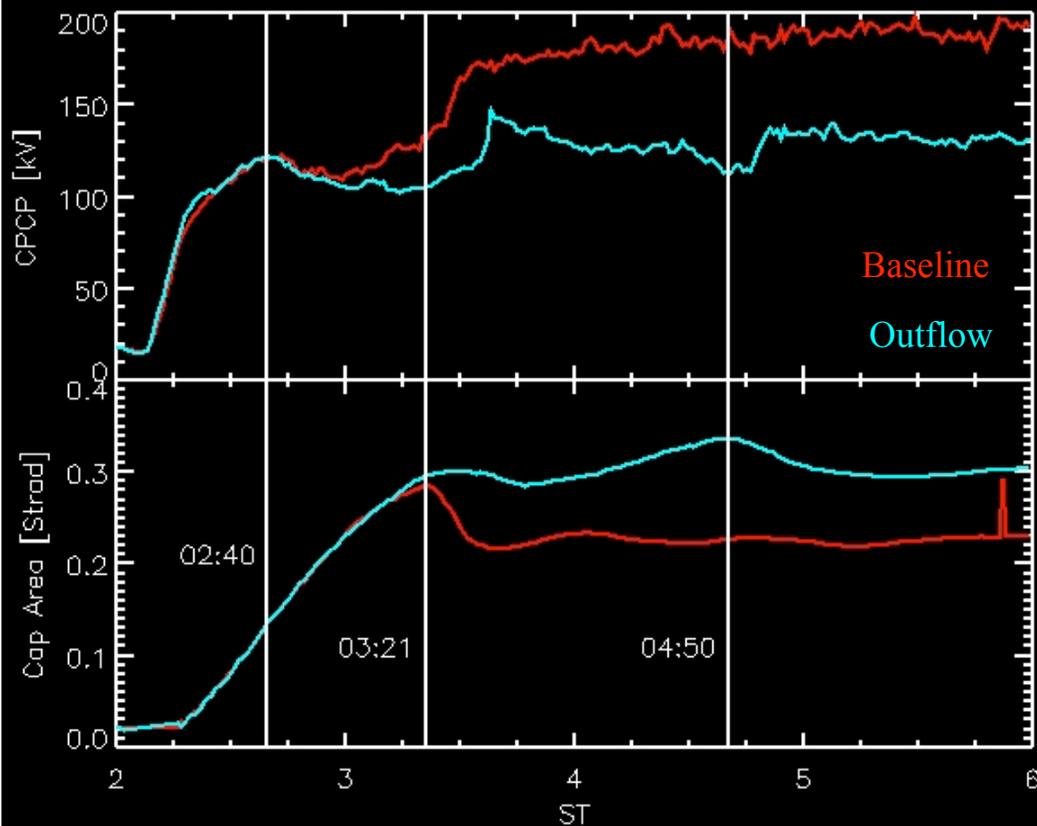
- Here we can see the release of the second plasmoid
 - Most of the O^+ outflow is now landing tailward of the reconnection region
 - O^+ is significant in the near Earth plasma sheet

05:59 ST



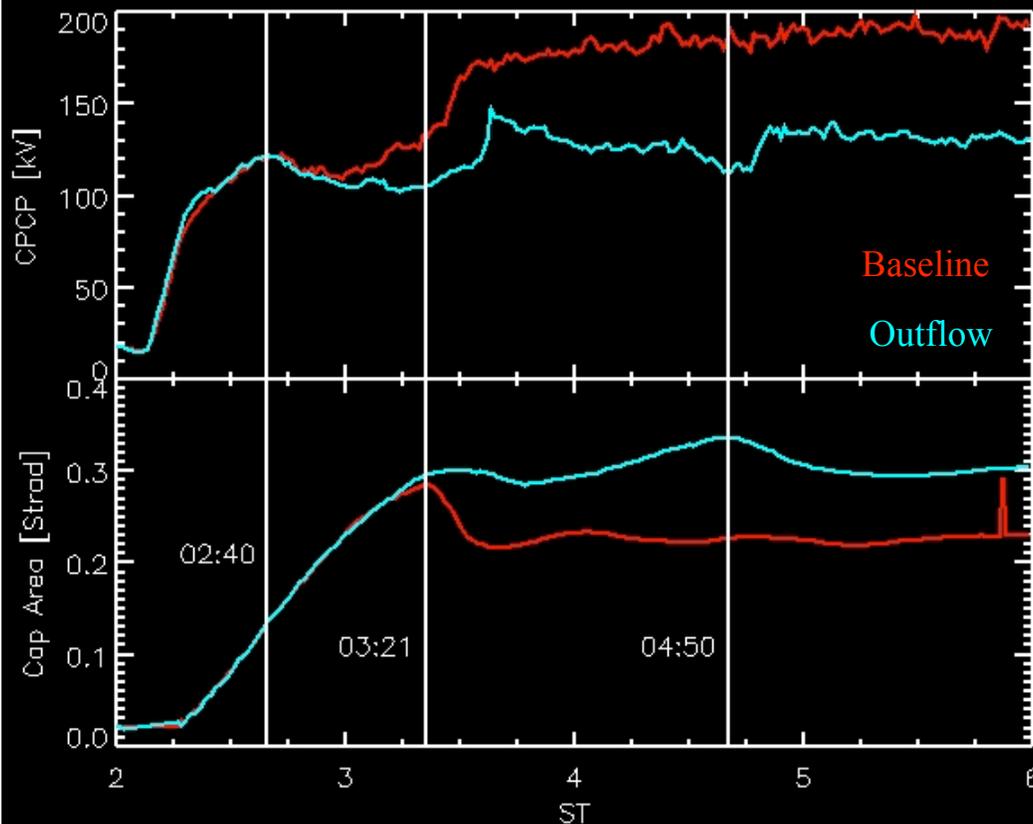
- The run with outflow has established a new quasteady equilibrium with the reconnection line near $x=-15R_E$
 - Contrast that with the reconnection point in the baseline simulation of $X=-20 R_E$

Cross Polar Cap Potential



- Shortly after the arrival of the southward IMF both simulations show a similar response in both the polar cap potential and area
 - An initial peak is seen at 02:40 ST near the time of initial plasmoid formation
- The simulation with outflow saturates a significantly lower potential
 - A global reduction in R1 currents is also seen
 - Linked to both changes in conductance and mass loading

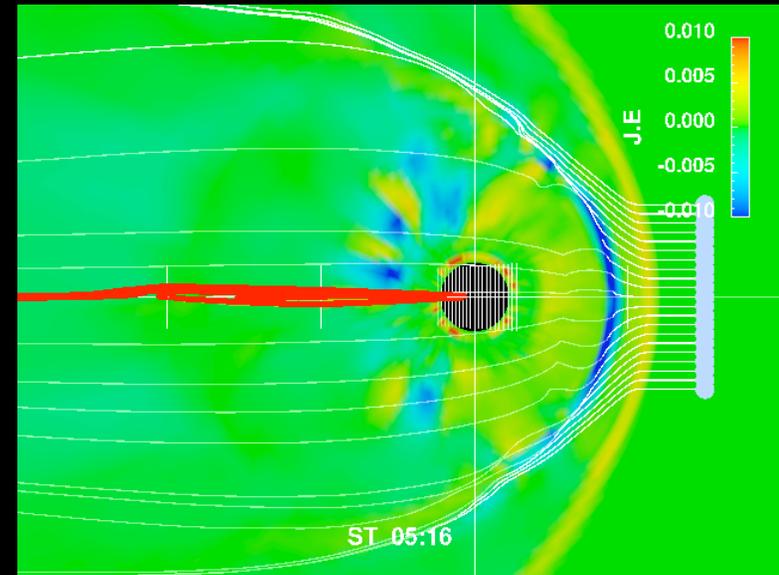
Polar Cap Area



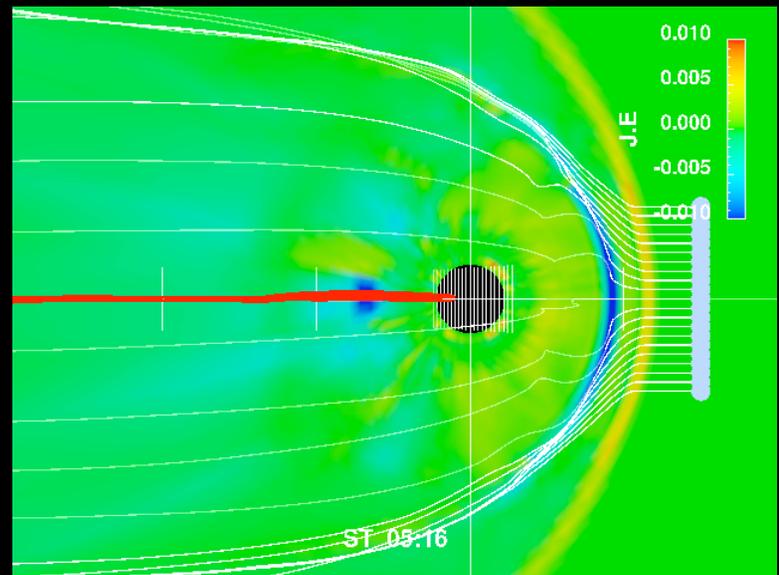
- The peak in polar cap area for the baseline run is seen at 03:21 ST, but in the outflow run it is seen 3 minutes later
 - This is a clear indication that the O+ is affecting the reconnection rate
- After the release of the first plasmoid the baseline simulation stabilizes while the polar cap area increases until a second plasmoid is released at 04:50 ST
 - A significant contributing factor is the repopulation of the plasma sheet

Hints of Nonlinearity

- Answering the question of how CPCP is lower with a lower polar cap area shows indications of nonlinear response
 - The outflow is impacting the shape of the magnetopause leading to a blunter object in the flow
 - This shape maps into a smaller region the solar wind and therefore a smaller polar cap potential



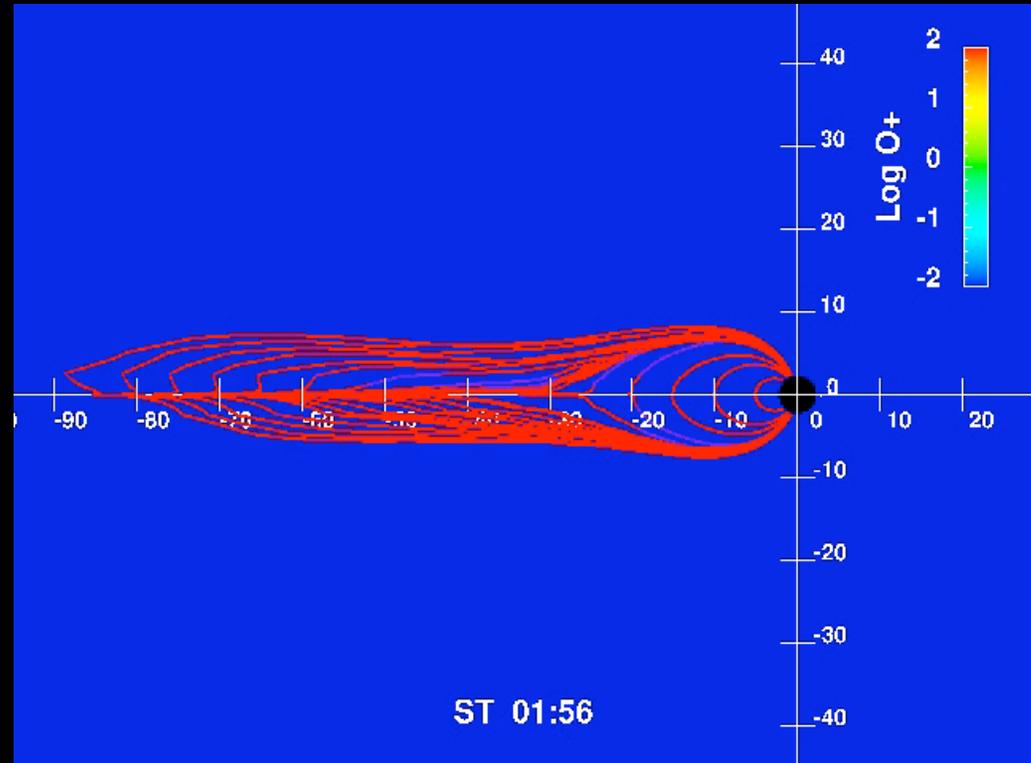
Baseline



Outflow

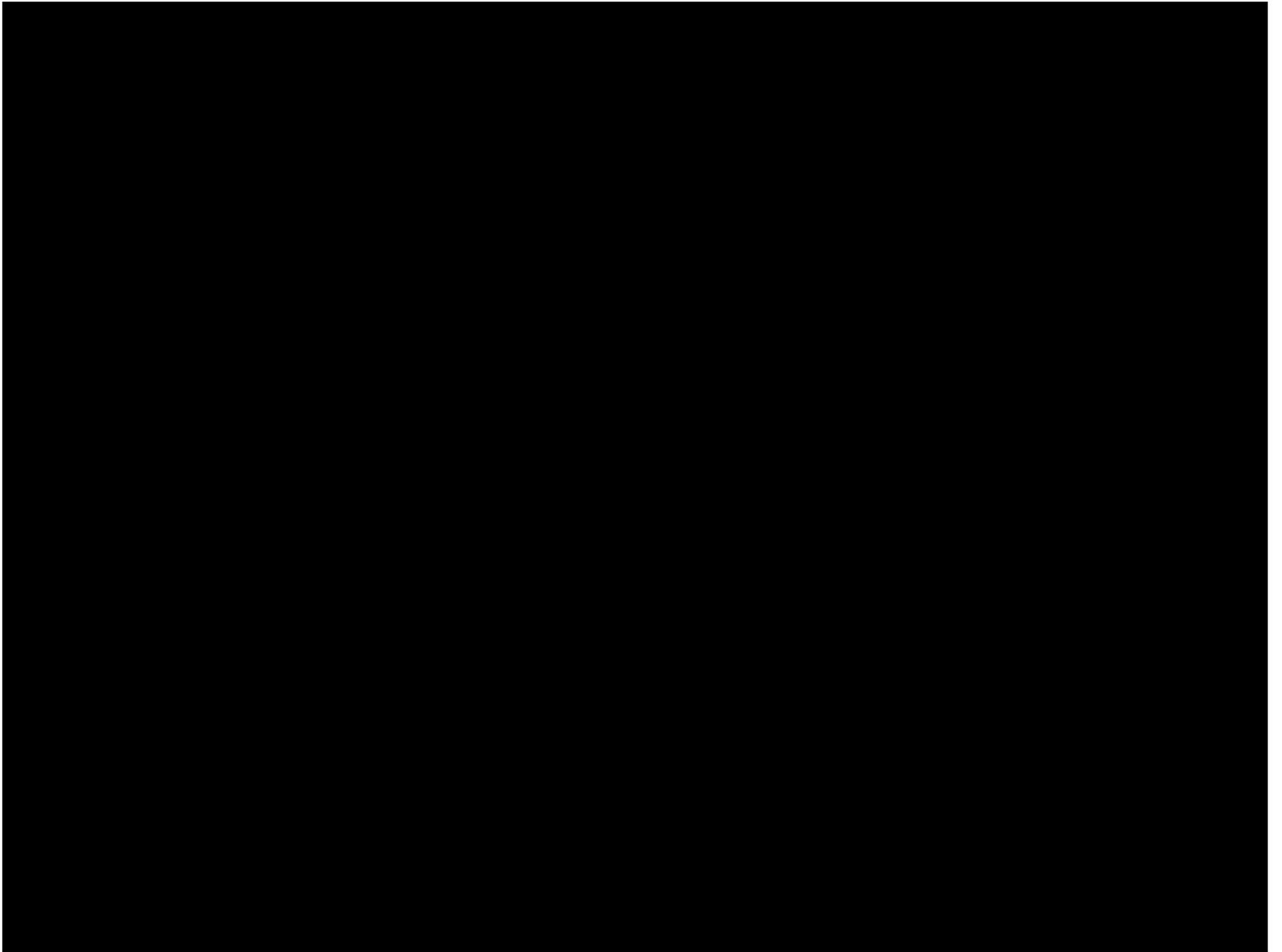
High Speed Cusp-Cleft Outflow

- In this case we increased the outflow velocity to 50 km/s while maintaining the same flux
- The initial substorm at is not dramatically affected by the outflow since it does not reach the X line
- Since the outflow lands tailward of the reconnection region the simulation enters the same quasi-steady state with no clear entry of O⁺ into the magnetotail



Ionospheric Outflow Summary

- We have clearly shown that ionospheric outflow dramatically impact the evolution of the magnetosphere
 - Cusp outflow which hits the near Earth reconnection site results in O^+ entering the plasma sheet creating conditions for a second substorm
- Initial results indicate that the impact of outflow on the evolution of the magnetosphere is sensitive to the parameters of the outflow
- Need to determine the relative importance of changes in magnetospheric shape and conductance are playing in the reduction of the cross polar cap potential
- More accurate and causally driven models of ionospheric outflow are needed



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