

# Theory and simulations of nonlinear whistler-mode chorus waves in the magnetosphere

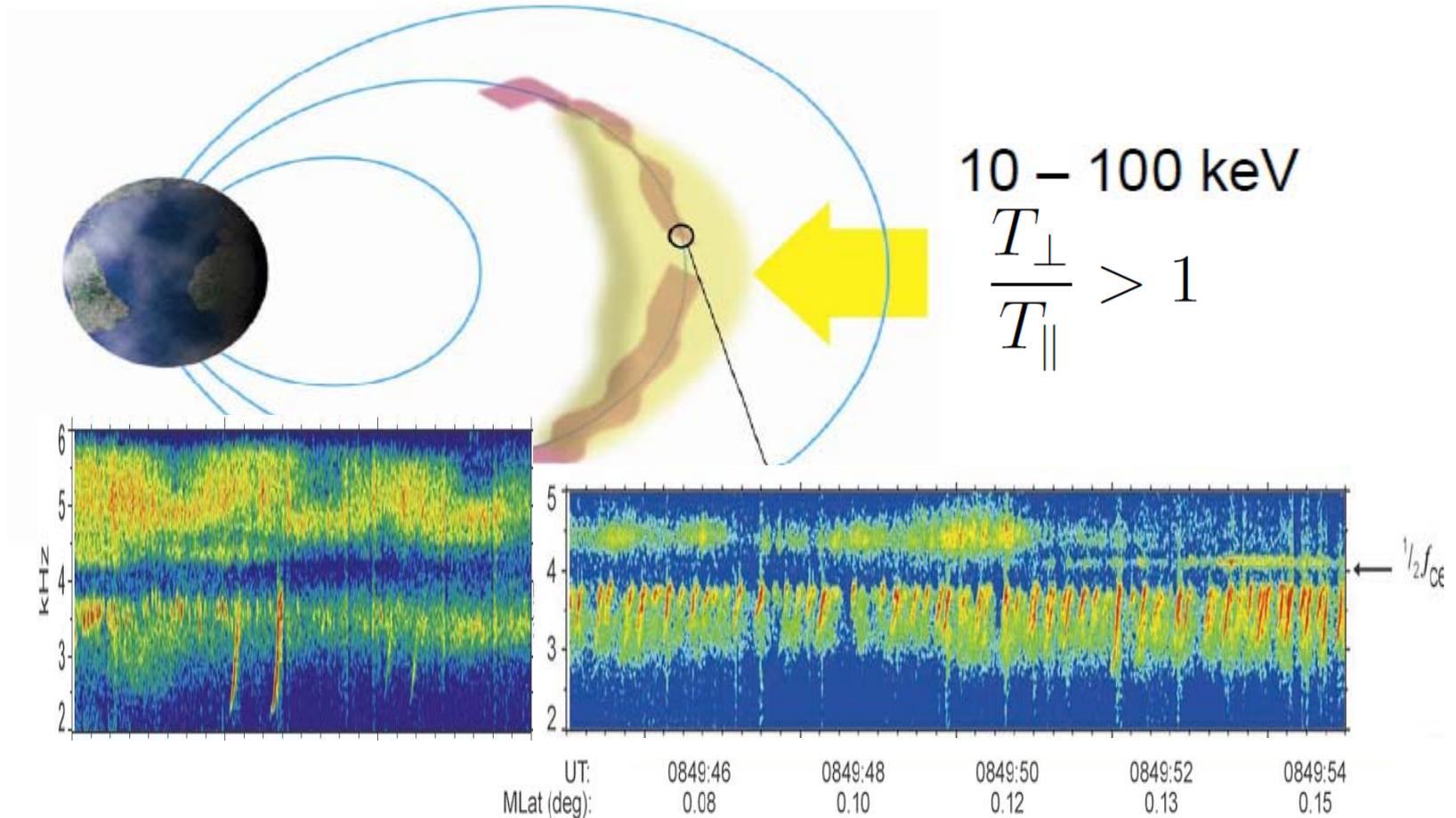
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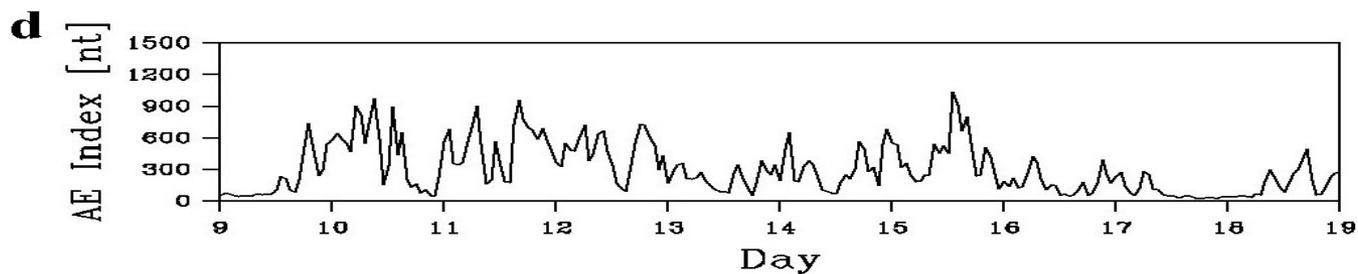
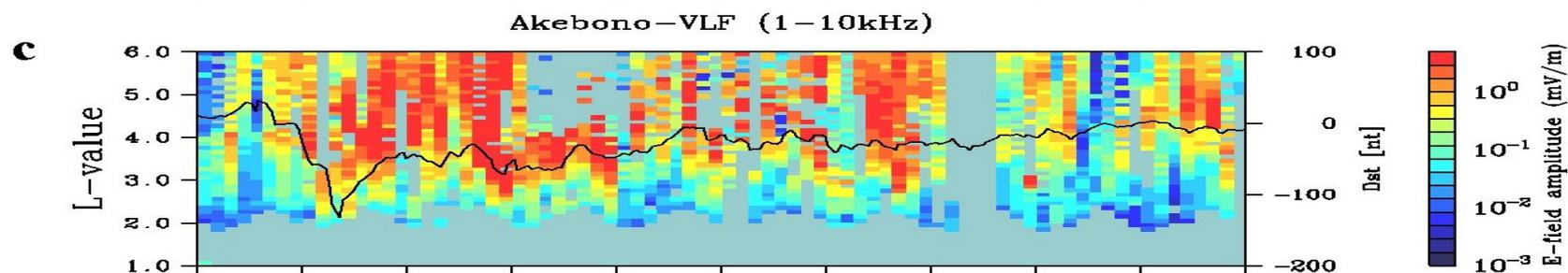
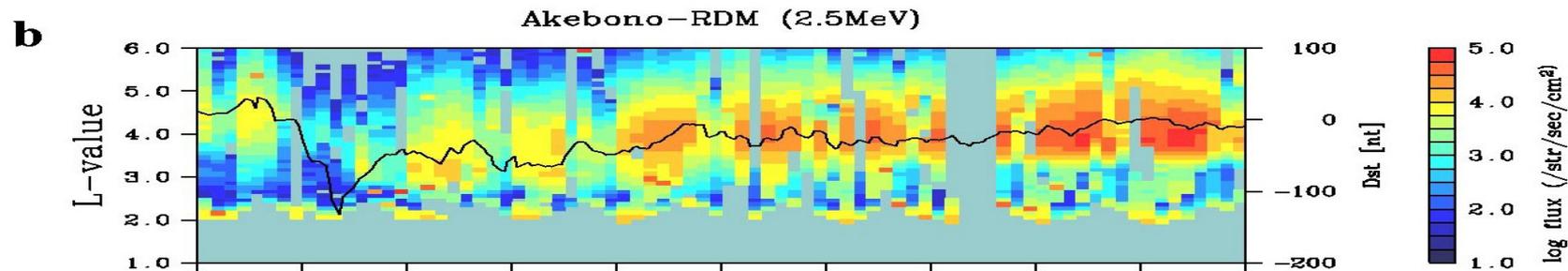
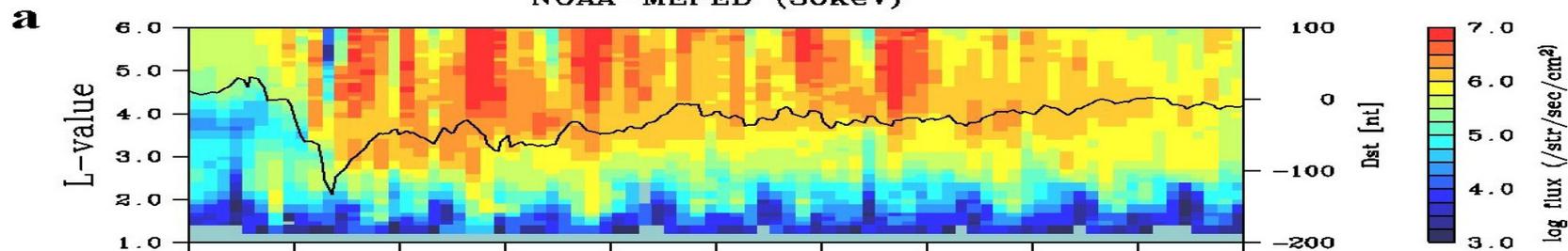
Modern Challenges in Nonlinear Plasma Physics, A Conference honoring the Career of  
Dennis Papadopoulos, Sani Resort, Halkidiki, Macedonia, Greece, June 15- 19, 2009

# Chorus Emission due to Injection of Energetic Electrons



[Cluster, Santolik and Gurnett,]

1990 / 10  
NOAA-MEPED (30keV)



[Kasahara et al., GRL, 2008]

# Overview

## Constant Frequency Wave

$$\frac{\partial \omega}{\partial t} = 0$$

- Linear growth rate
- Quasi-linear diffusion
- Nonlinear trapping and acceleration in an inhomogeneous medium: RTA and URA

## Rising-tone Chorus Element

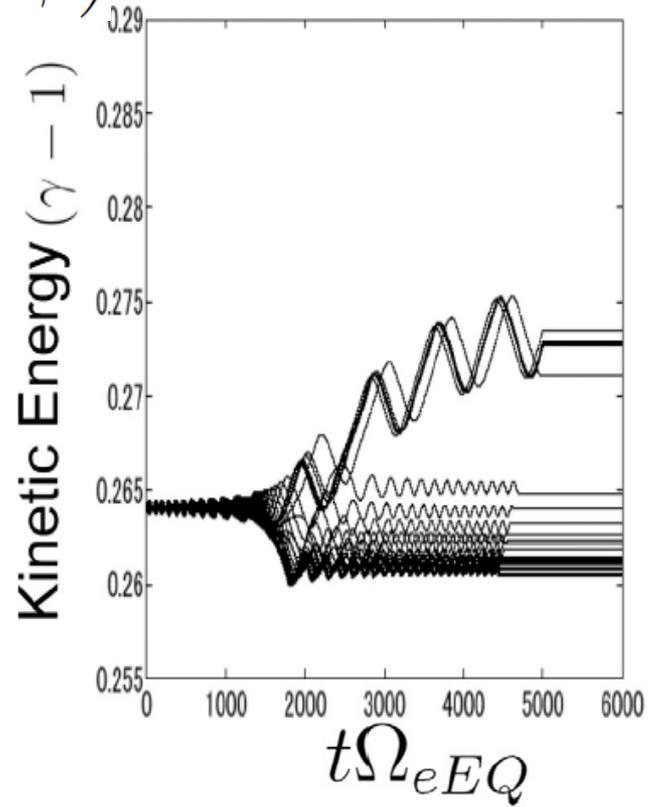
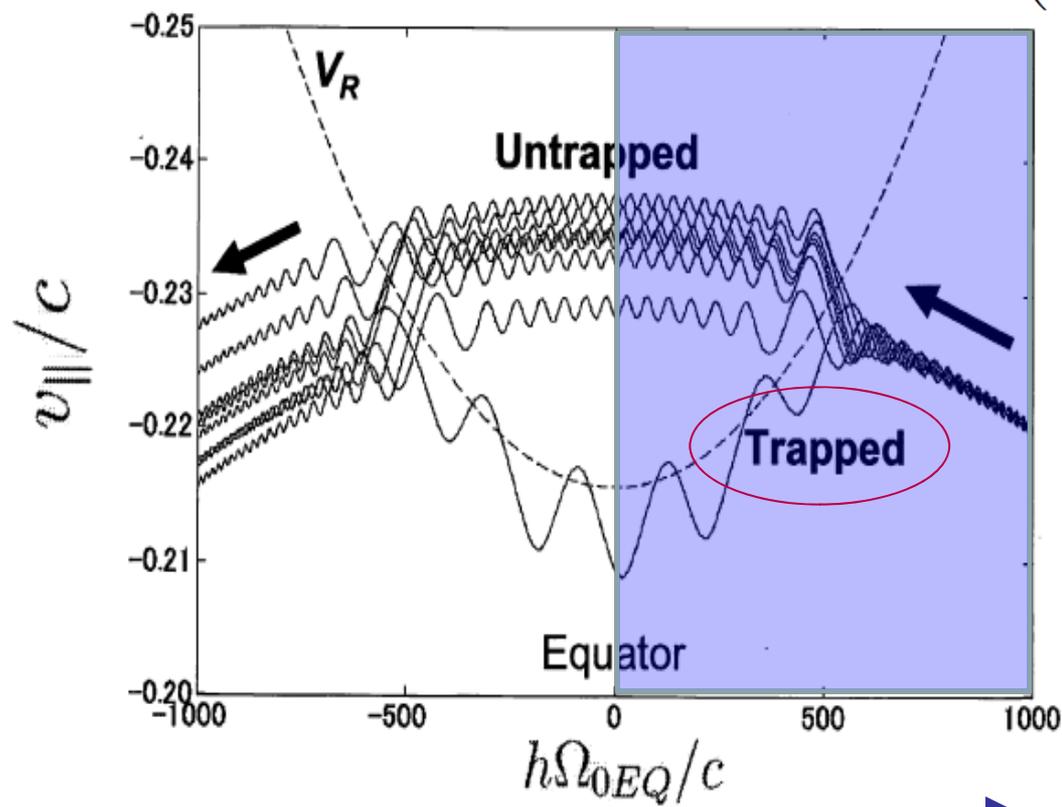
$$\frac{\partial \omega}{\partial t} > 0$$

- Nonlinear-wave growth
- Generation of upper-band chorus emissions
- Efficient acceleration by nonlinear trapping and formation of the pancake distribution

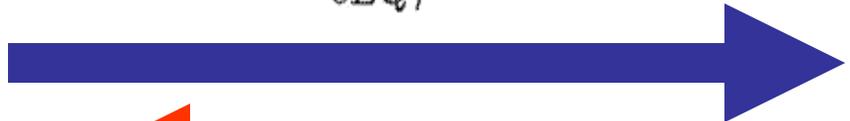
# Test Particle Simulation of Nonlinear Wave Trapping and Acceleration of Trapped Particles (10 – 100 keV)

$$\frac{\partial \omega}{\partial t} = 0$$

$$\omega - kv_{\parallel} = \Omega_e(h) / \gamma \quad \Rightarrow \quad V_R = \frac{1}{k} \left( \omega - \frac{\Omega_e}{\gamma} \right)$$



Wave

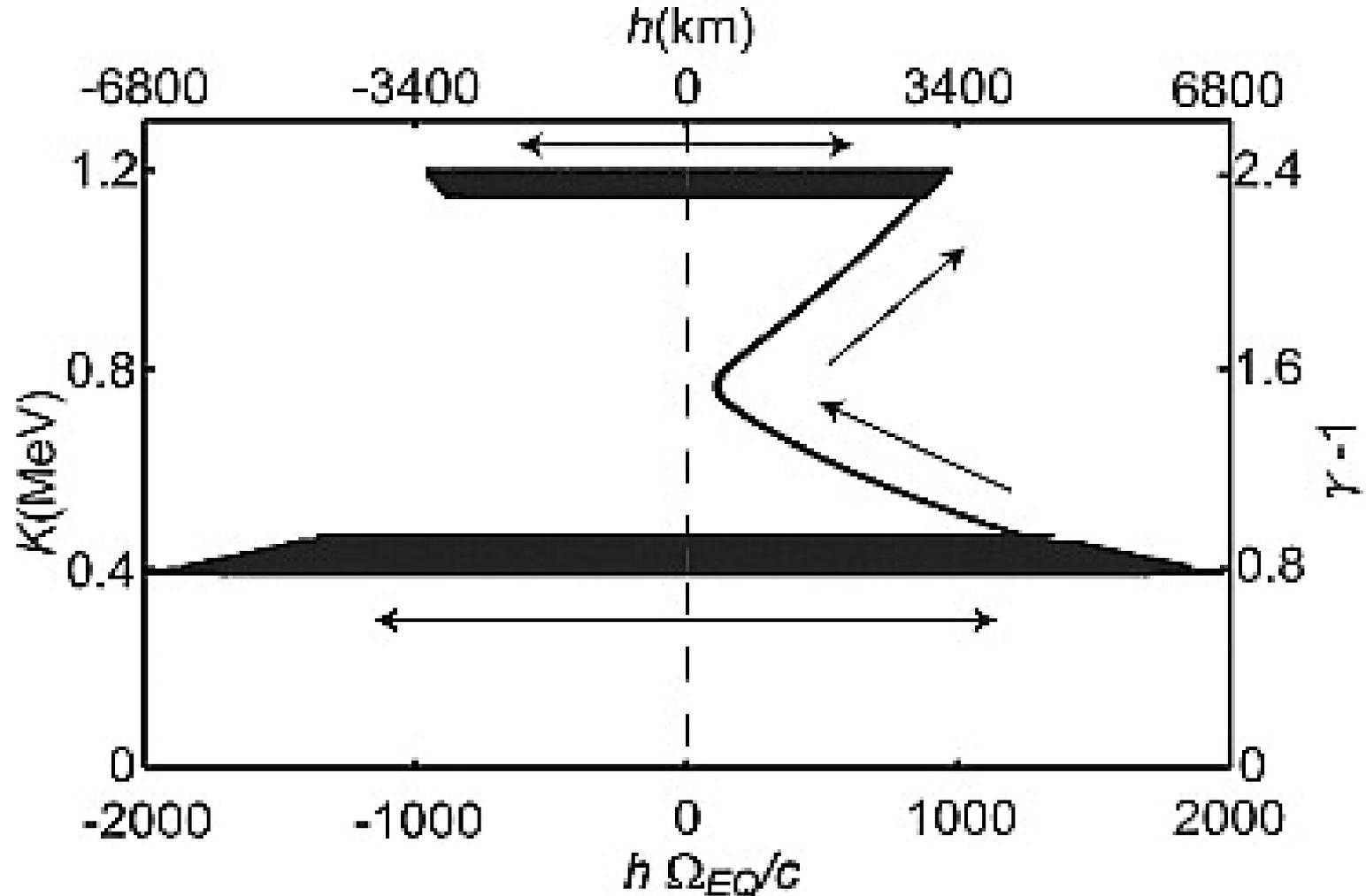


Electrons

# Trajectories of Resonant Electrons (400 keV)

$$\frac{\partial \omega}{\partial t} = 0$$

Relativistic Turning Acceleration (RTA)



**Bw = 125 pT**

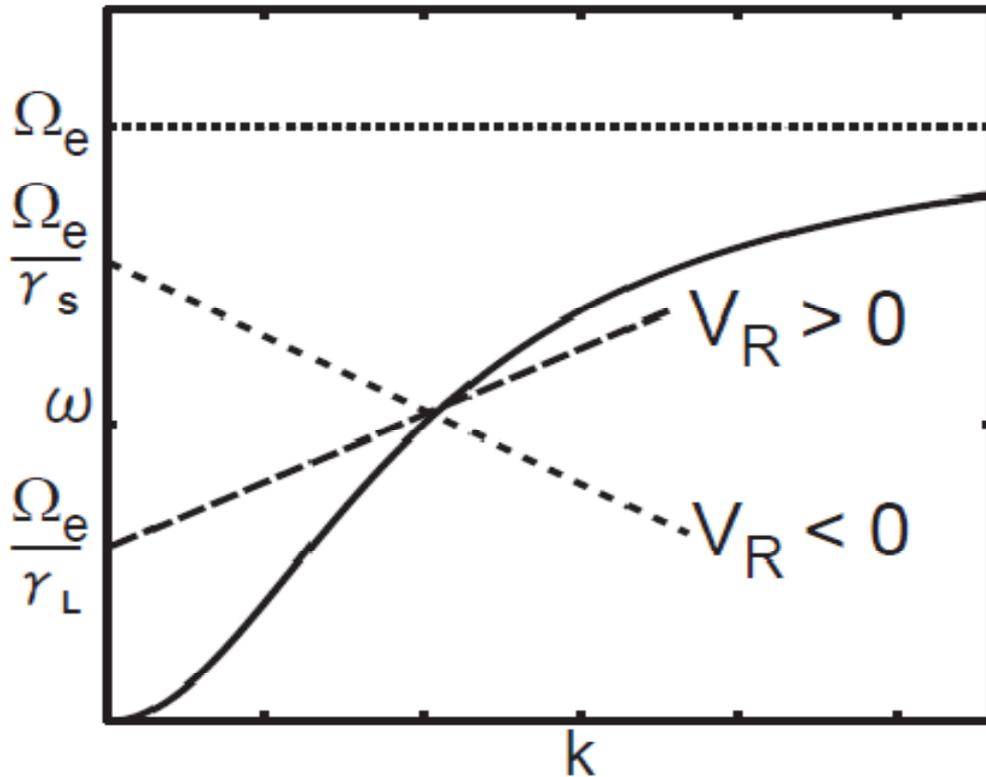
$$\omega_p = 2.0 \Omega_{e0} \quad \omega = 0.4 \Omega_{e0}$$

[Omura, Furuya, and Summers, JGR, 2007]

## Resonance Velocity

$$V_R = \frac{\omega}{k} \left( 1 - \frac{\Omega_e(h)}{\omega \gamma} \right) \quad \gamma = [1 - (v_{\parallel}^2 + v_{\perp}^2)/c^2]^{-1/2}$$

$$\gamma \geq \Omega_e(h)/\omega \Rightarrow V_R \geq 0$$



$$K = m_0 c^2 (\gamma - 1)$$

$$\frac{dK}{dt} = -e E_w v_{\perp} S$$

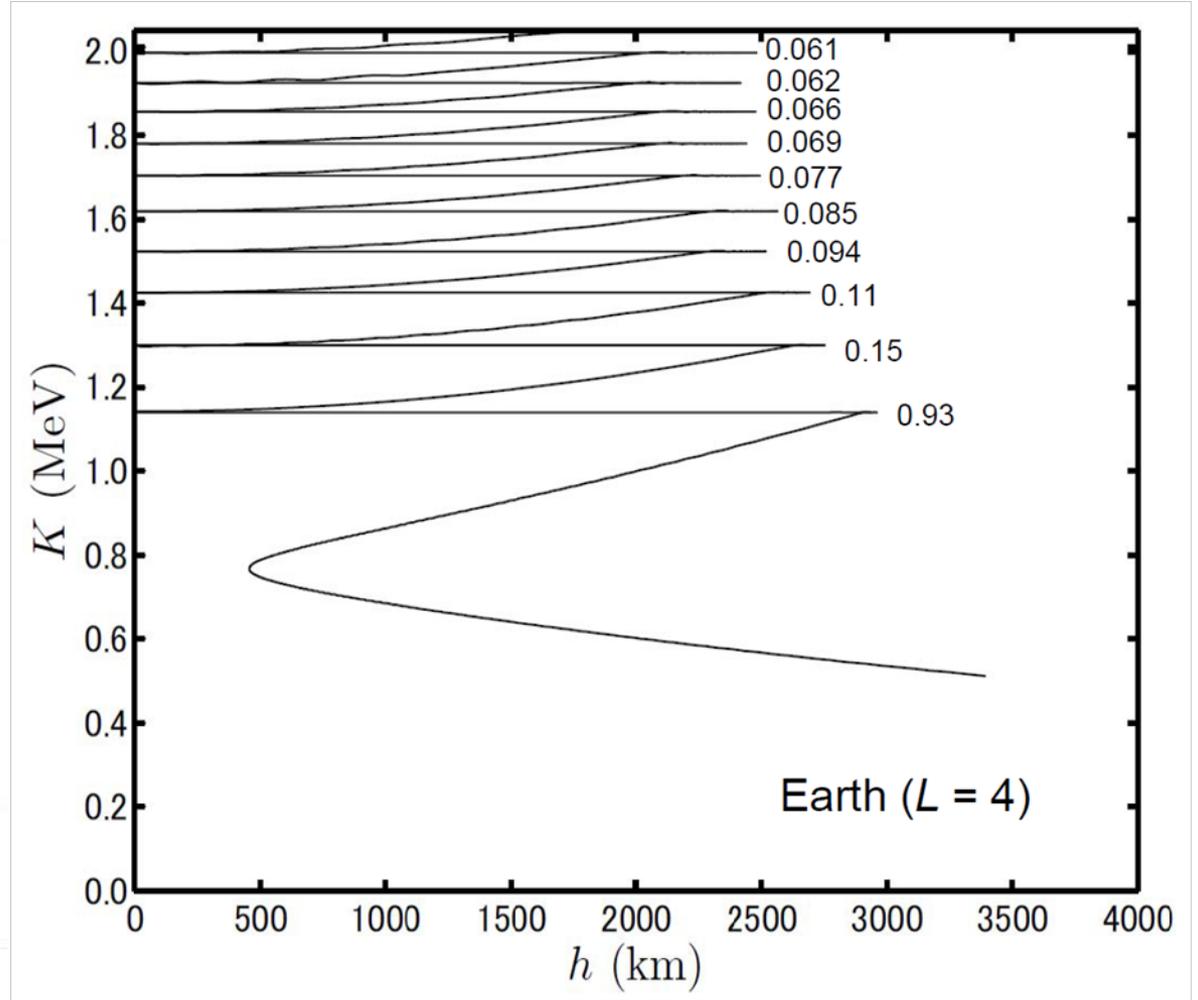
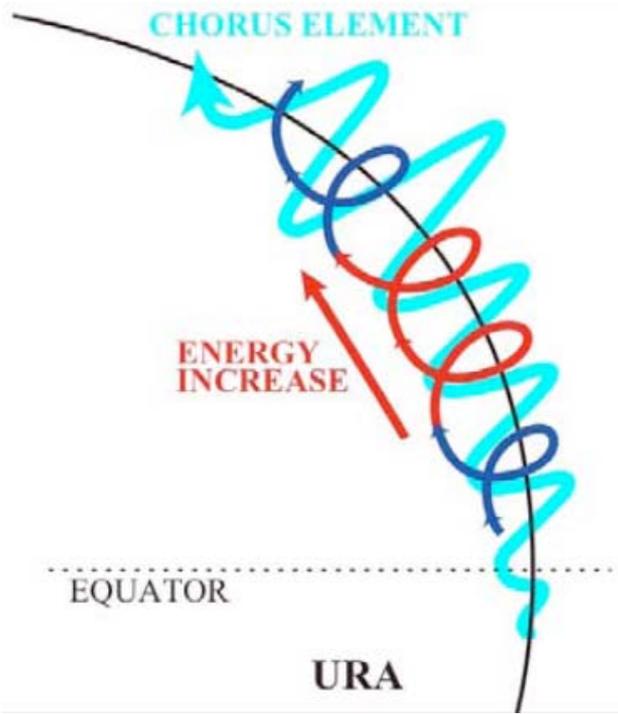
$$\frac{\partial \omega}{\partial t} = 0$$

# Trajectories of Resonant Electrons (> 1MeV)

## Ultra-Relativistic Acceleration (URA)

$$\gamma_0 > \Omega_{EQ} / \omega$$

$$V_{R0} > 0$$



$$\frac{\partial \omega}{\partial t} = 0$$

$$\frac{d\zeta}{dt} = \theta$$

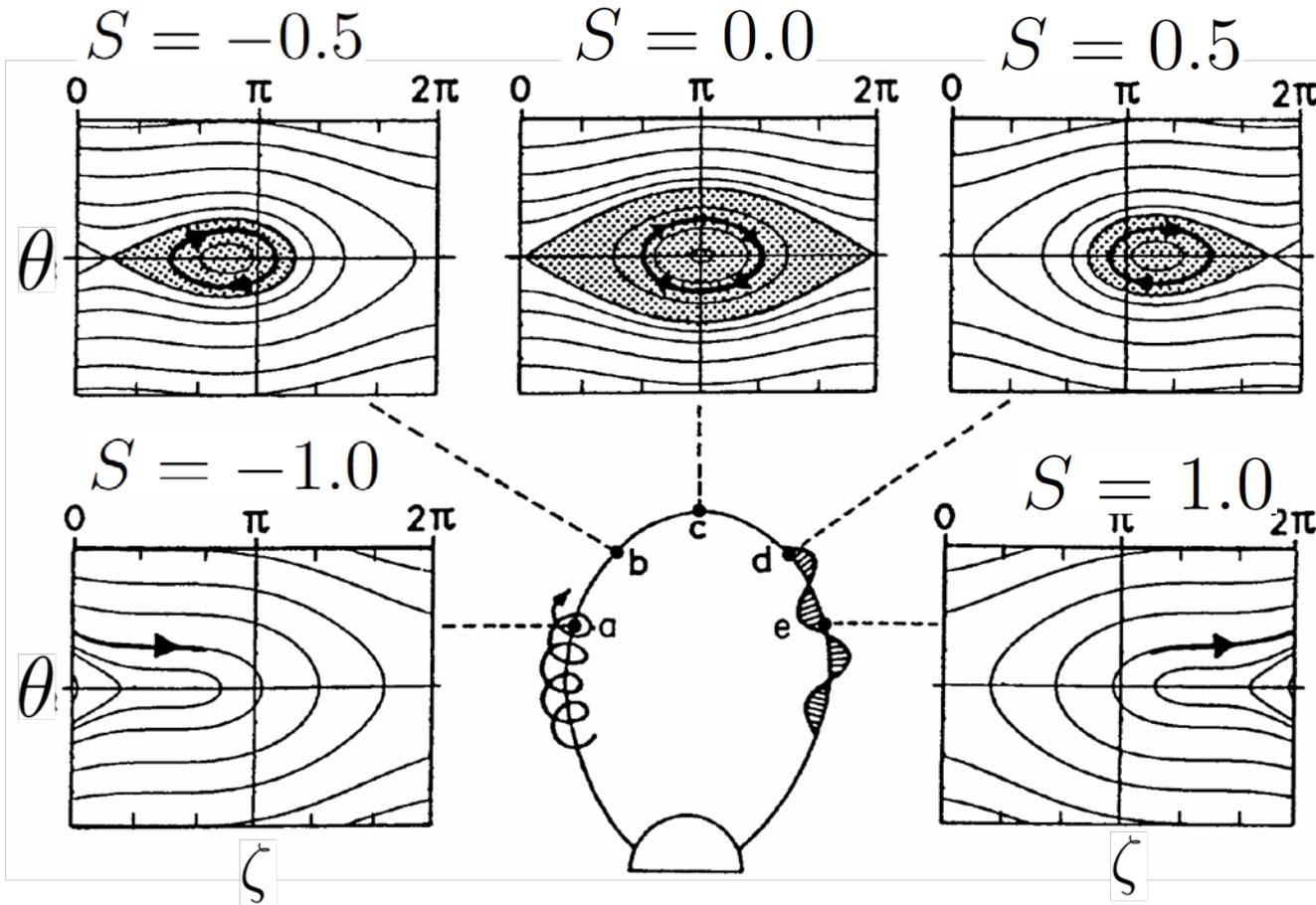
$$\frac{d\theta}{dt} = \omega_t^2 (\sin \zeta + S)$$

$$\theta = k(v_{\parallel} - V_R)$$

for Whistler-mode Wave

$$\theta = -k(v_{\parallel} - V_p)$$

for Longitudinal Wave



$$\frac{\partial \omega}{\partial t} > 0$$

## Inhomogeneity Ratio

$$S = -\frac{1}{s_0 \omega \Omega_w} \left( s_1 \frac{\partial \omega}{\partial t} + c s_2 \frac{\partial \Omega_e}{\partial h} \right)$$

$$s_0 = \frac{\delta V_{\perp 0}}{\xi c} \quad s_1 = \gamma \left( 1 - \frac{V_R}{V_g} \right)^2$$

$$s_2 = \frac{1}{2\xi\delta} \left\{ \frac{\gamma\omega}{\Omega_e} \left( \frac{V_{\perp 0}}{c} \right)^2 - \left[ 2 + \Lambda \frac{\delta^2 (\Omega_e - \gamma\omega)}{\Omega_e - \omega} \right] \frac{V_R V_p}{c^2} \right\}$$

$\Lambda = \omega / \Omega_e$  for inhomogeneous density model

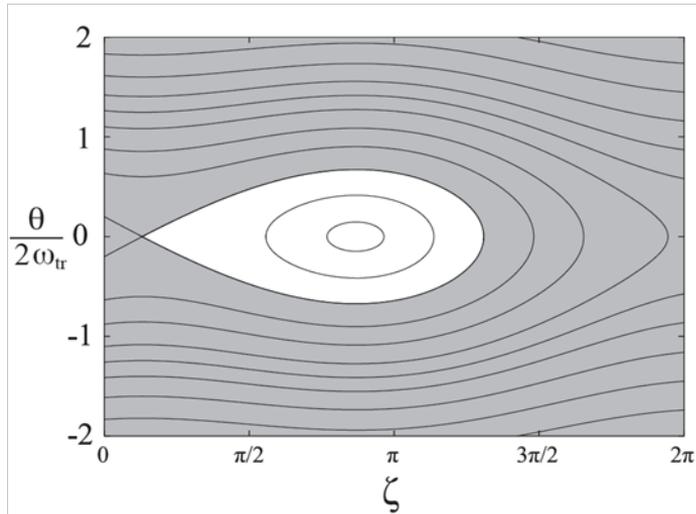
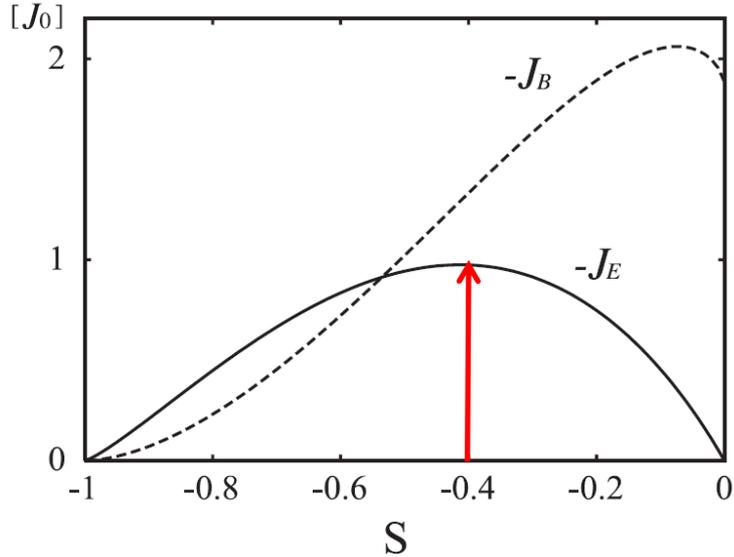
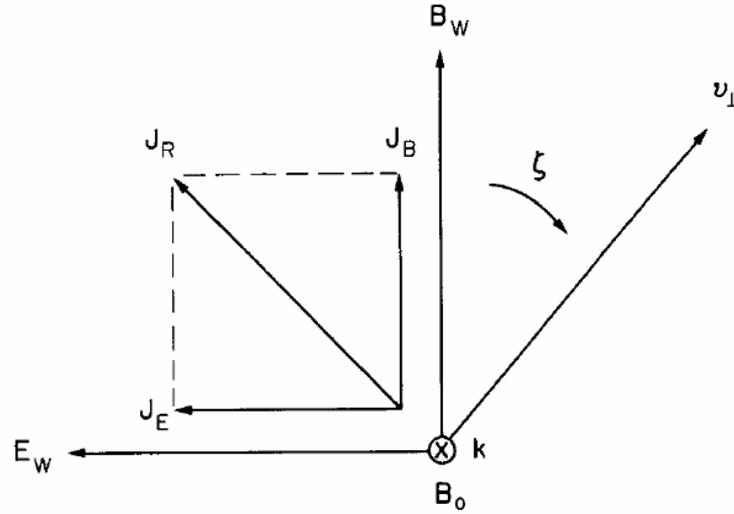
$\Lambda = 1$  for constant density model

# Nonlinear Wave Growth due to Formation of Electromagnetic Electron Hole

$$\frac{\partial \omega}{\partial t} > 0$$

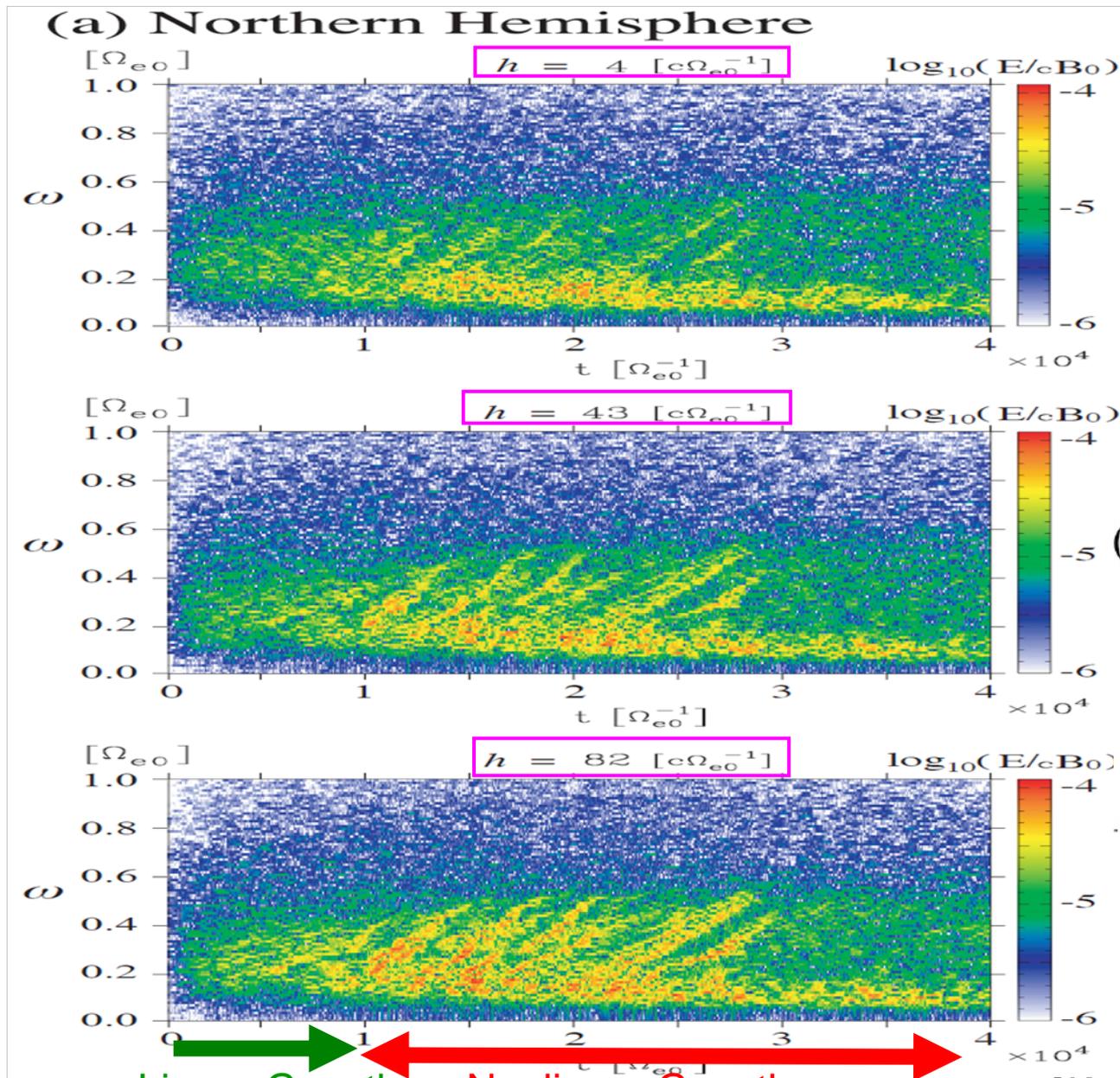
$$\frac{\partial B_w}{\partial t} + V_g \frac{\partial B_w}{\partial h} = -\frac{\mu_0 V_g}{2} J_E$$

$$\frac{\partial \omega}{\partial t} + V_g \frac{\partial \omega}{\partial h} = 0$$

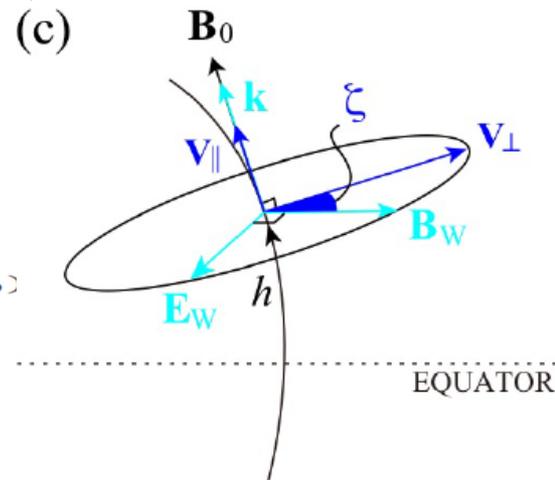
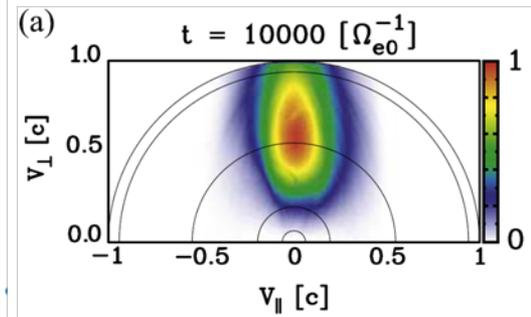


Maximum  $J_E$   $S_{EQ} = -0.4$

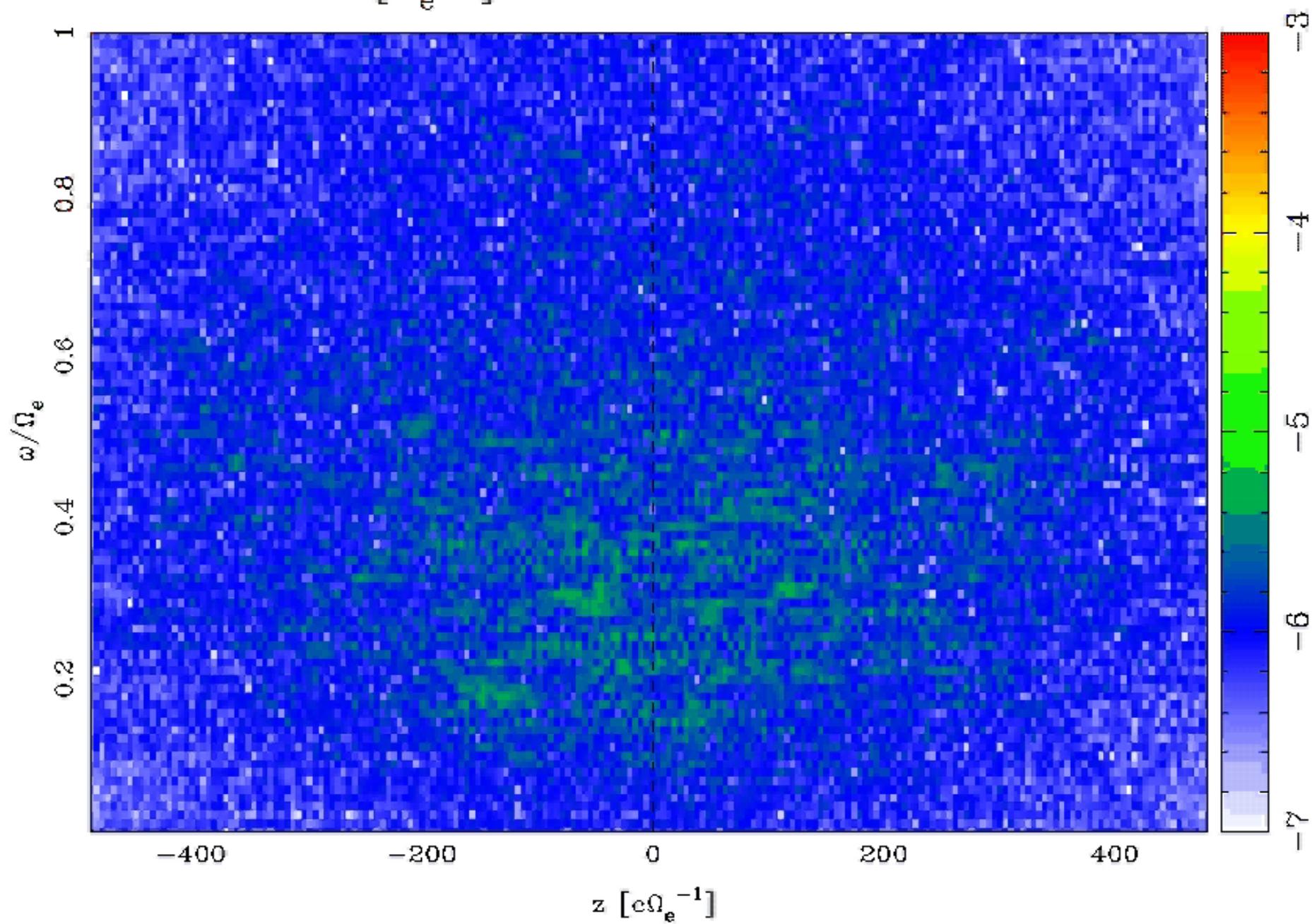
# Electron Hybrid Simulation (cold electrons : fluid)



$$T_{\perp}/T_{\parallel} = 9$$

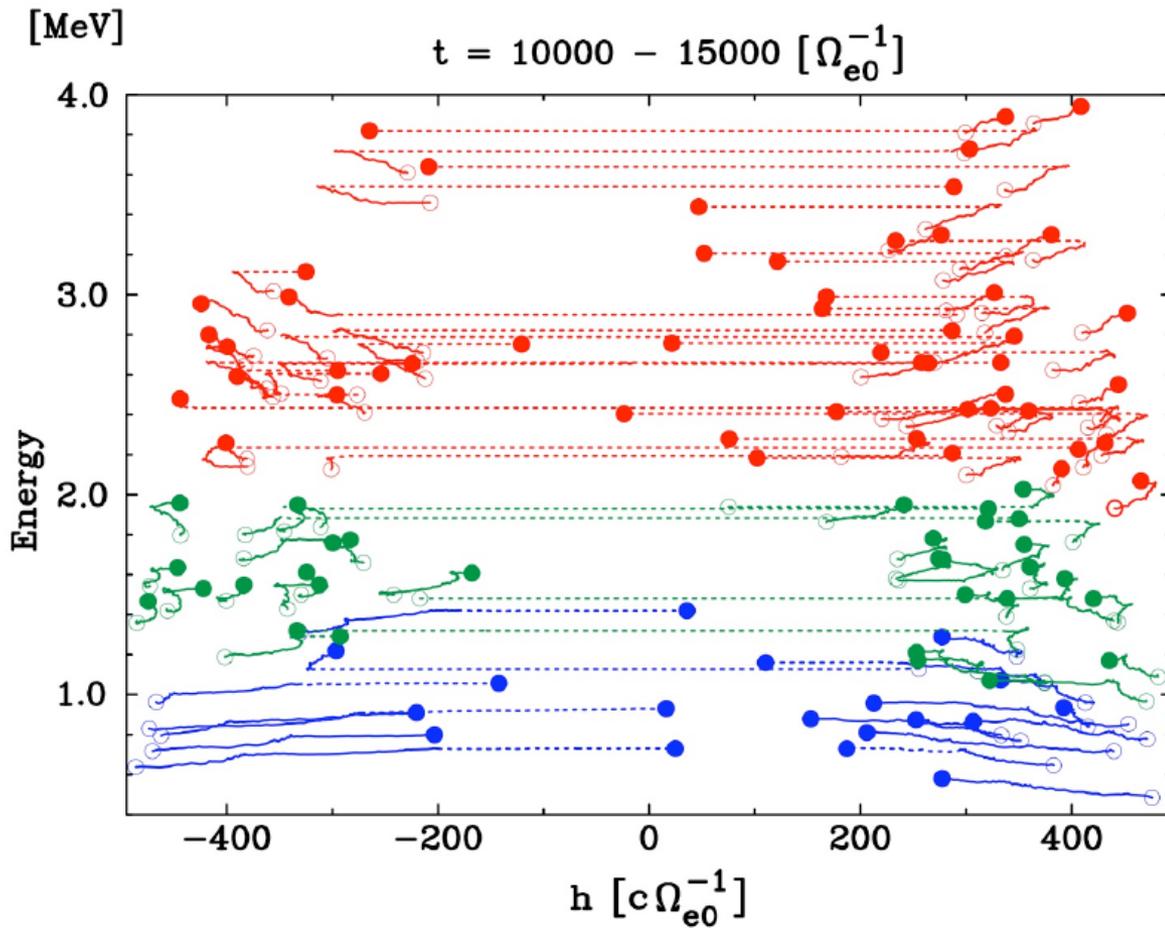


$t = 599.04 [\Omega_e^{-1}]$

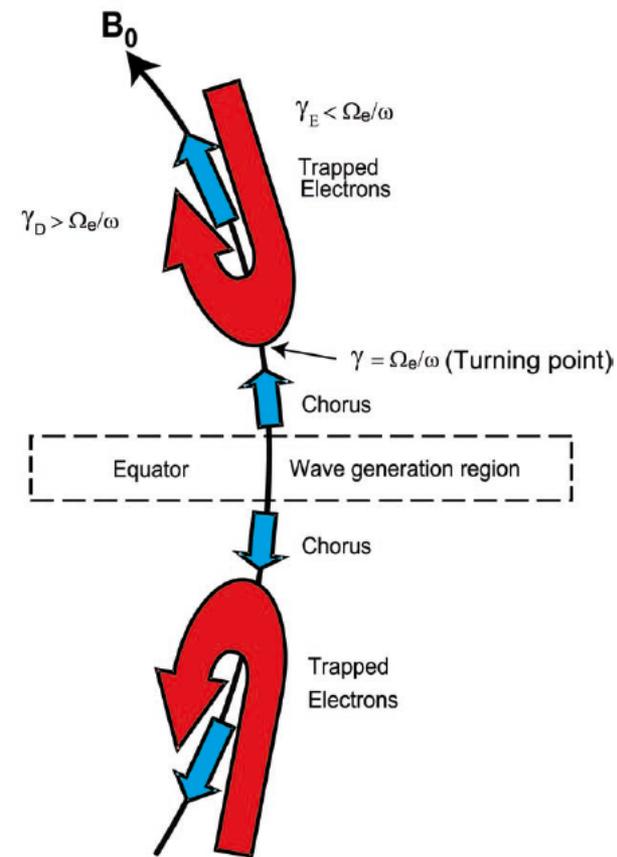


$$\frac{\partial \omega}{\partial t} > 0$$

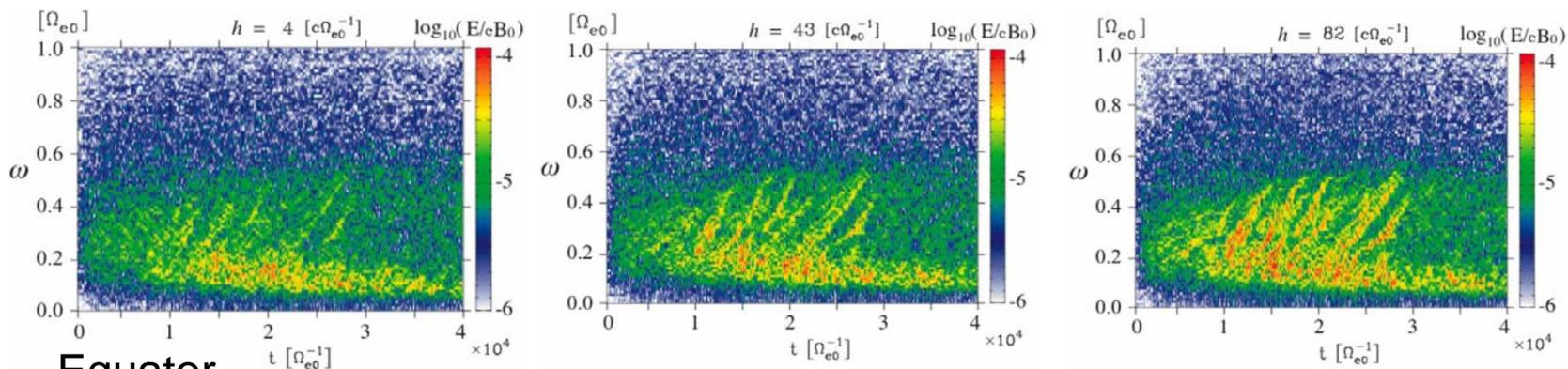
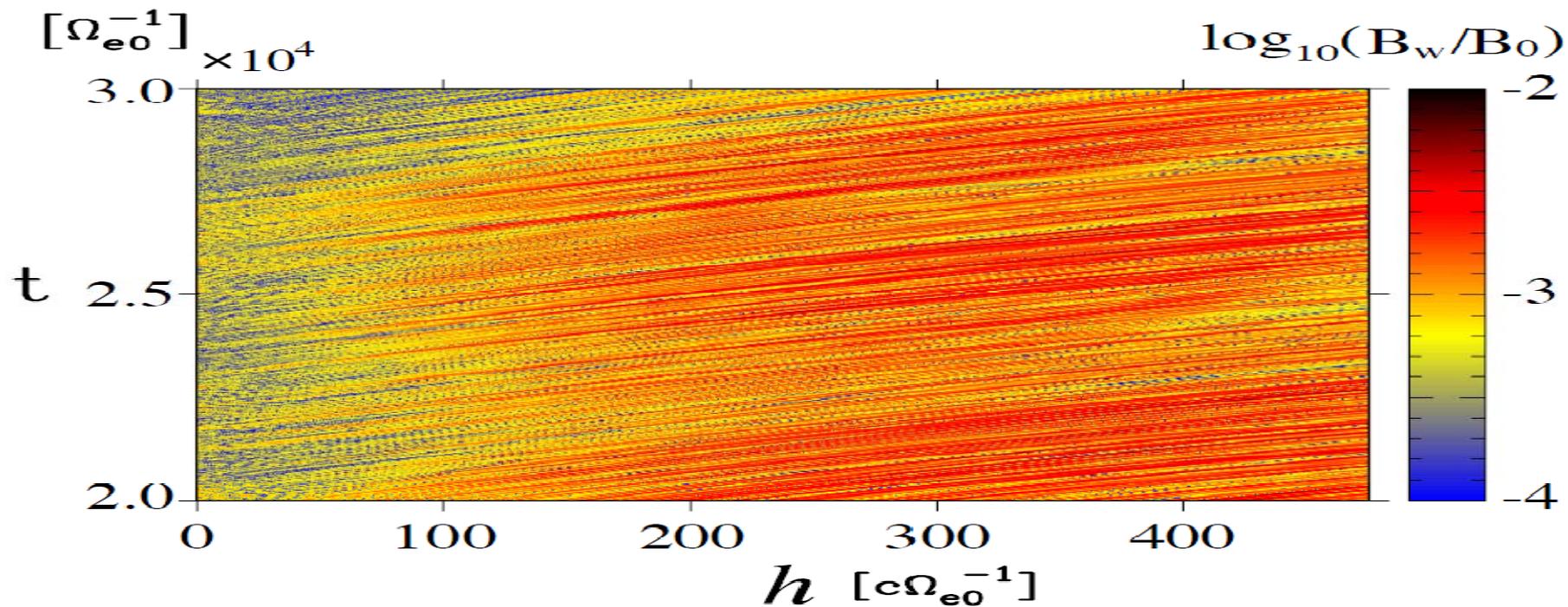
# RTA and URA in the chorus generation process



## RELATIVISTIC TURNING ACCELERATION



[Katoh, Omura and Summers, *Ann. Geophys.* 2008]



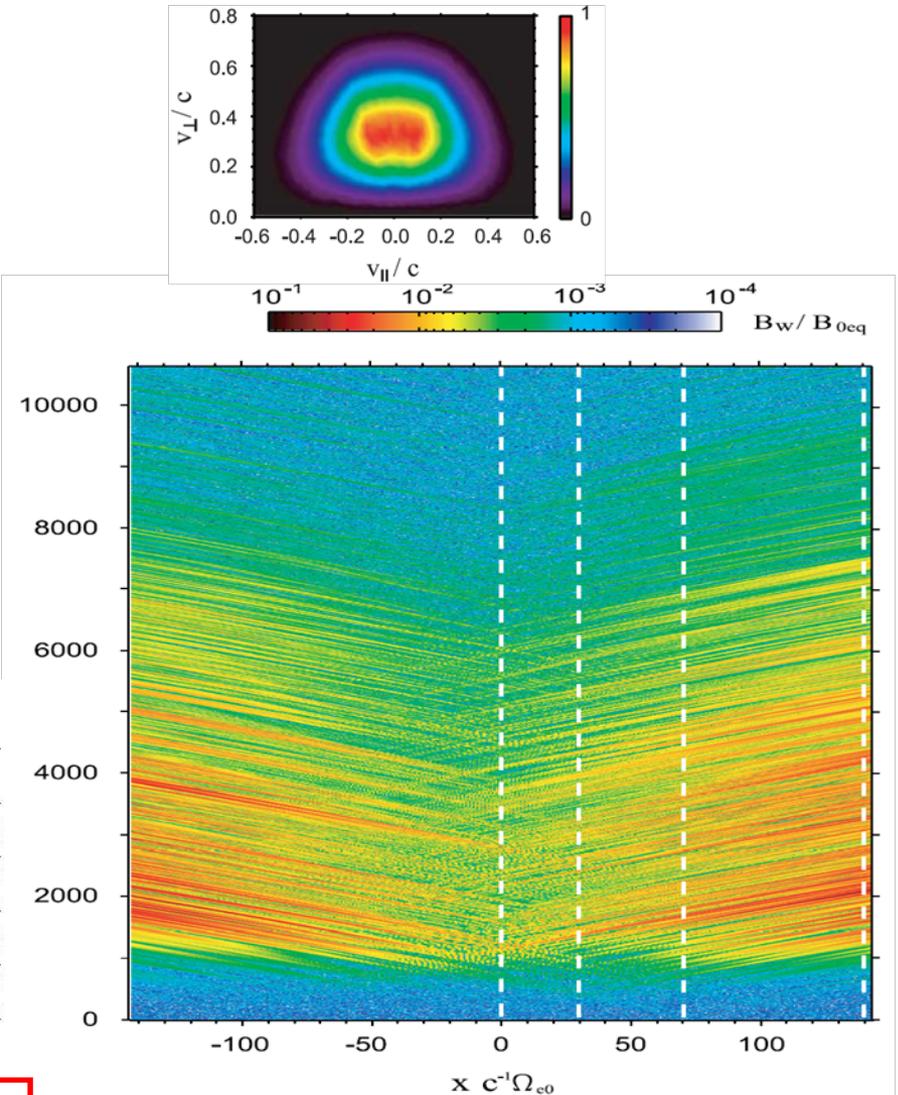
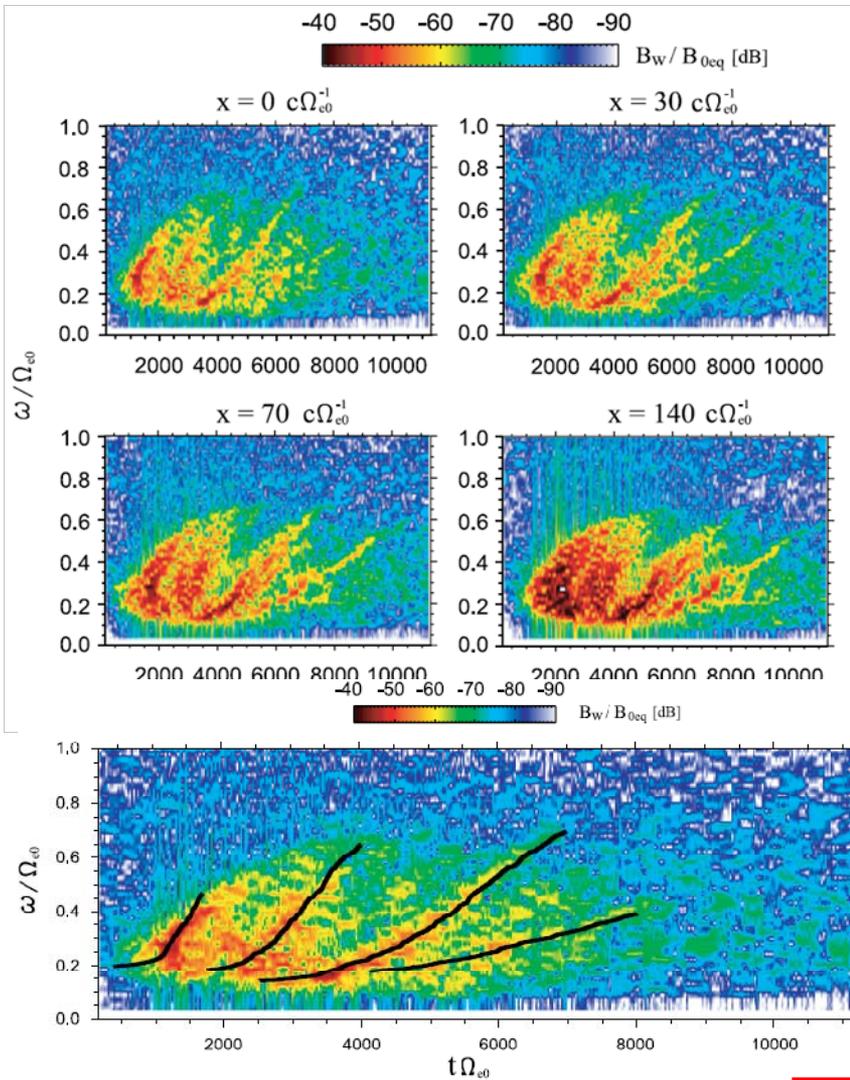
Equator



$$S = \frac{1}{B_w \omega_t^2 \delta^2} \left[ \gamma \left( 1 - \frac{V_R}{V_g} \right)^2 \frac{\partial \omega}{\partial t} + \left\{ \frac{k \gamma v_{\perp}^2}{2 \Omega_e} - \left( 1 + \frac{\delta^2 \Omega_e - \gamma \omega}{2 \Omega_e - \omega} \right) V_R \right\} \frac{\partial \Omega_e}{\partial h} \right]$$

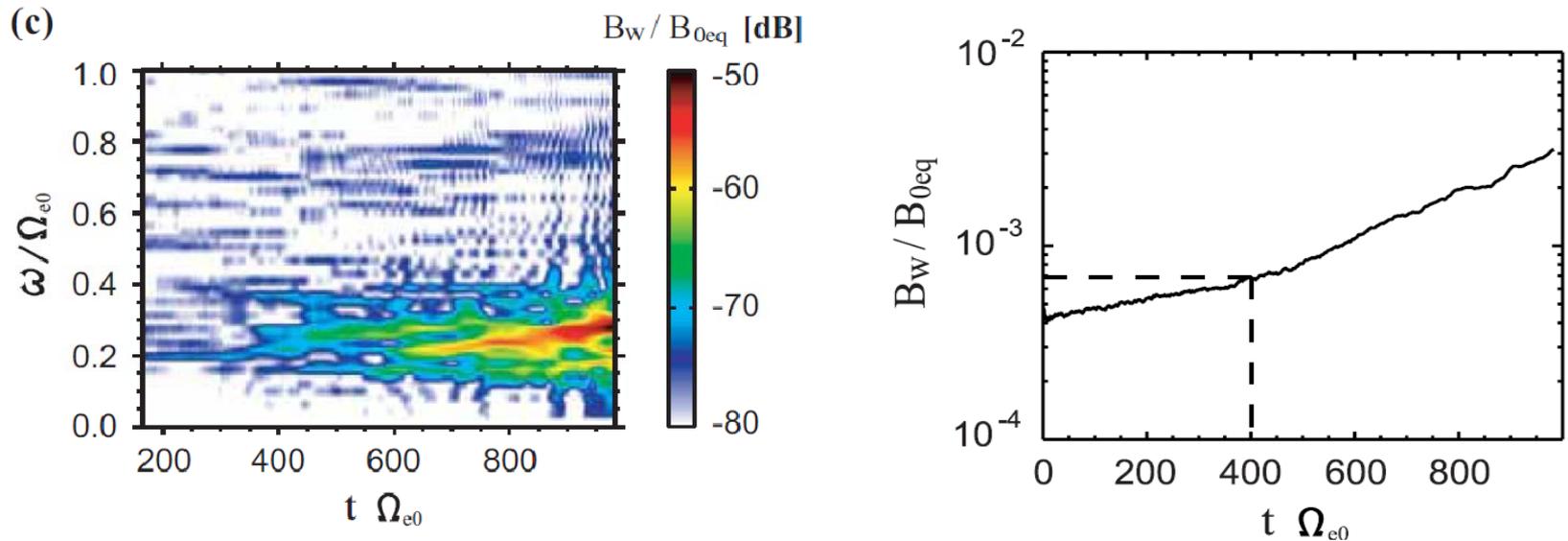
# Full-Particle Simulation of Chorus Emission

$$T_{\perp} / T_{\parallel} = 3$$



$$\frac{\partial \omega}{\partial t} = \frac{0.4 \delta V_{\perp 0}}{\gamma \xi} \frac{\omega}{c \Omega_{e0}} \left( 1 - \frac{V_R}{V_g} \right)^{-2} \frac{B_w}{B_{0eq}} \Omega_{e0}^2$$

[Hikishima et al., JGR, 2009]



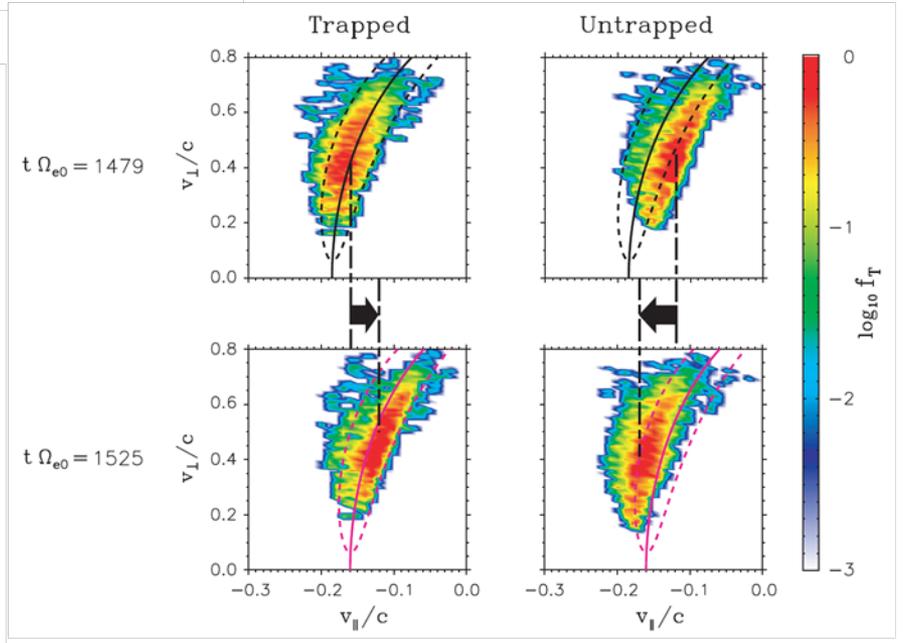
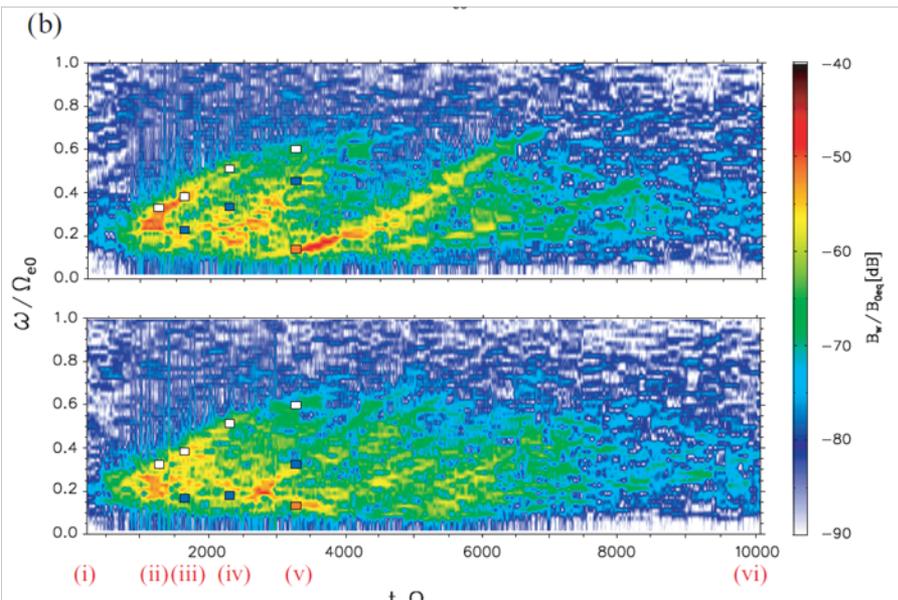
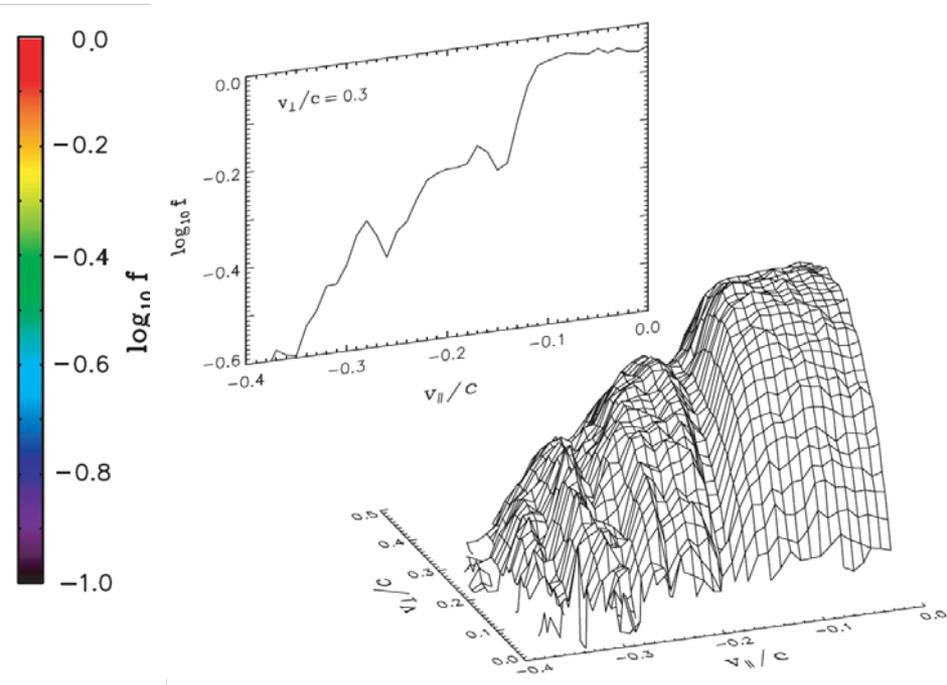
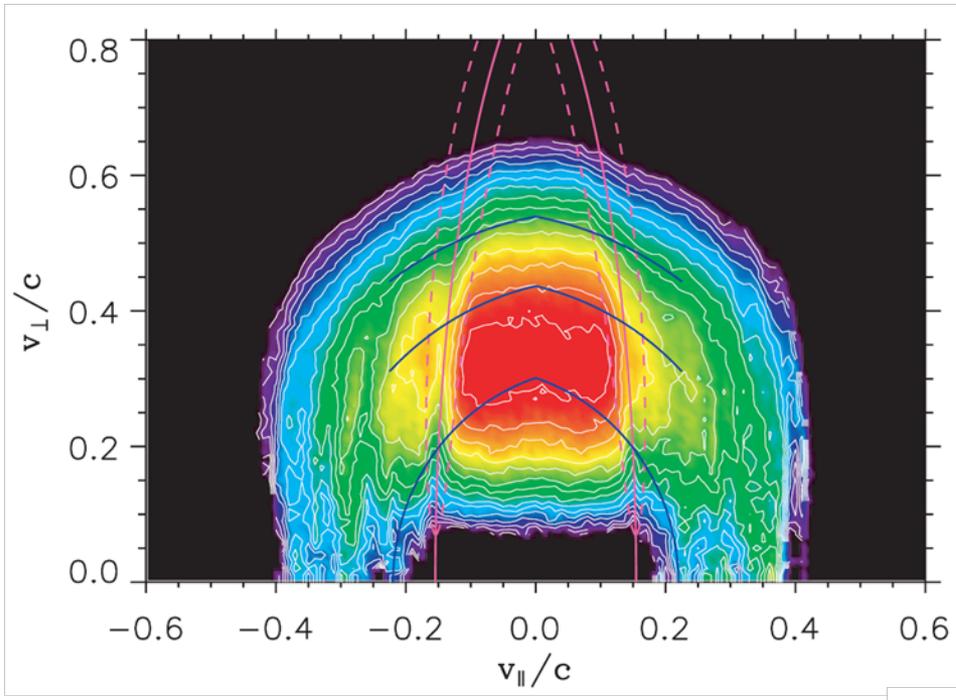
The condition for **absolute instability** at the equator

$$\frac{\partial B_w}{\partial t} = -\frac{\mu_0 V_g}{2} J_E - V_g \frac{\partial B_w}{\partial h} > 0$$

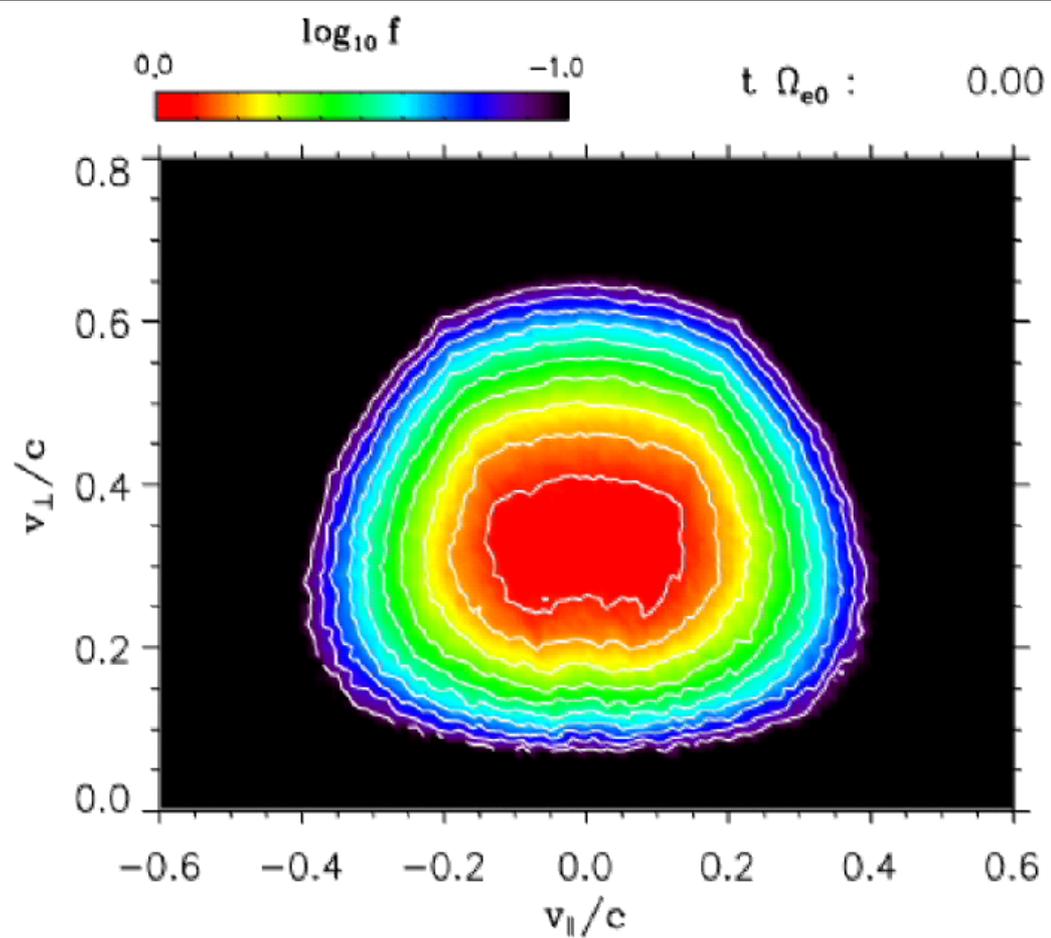
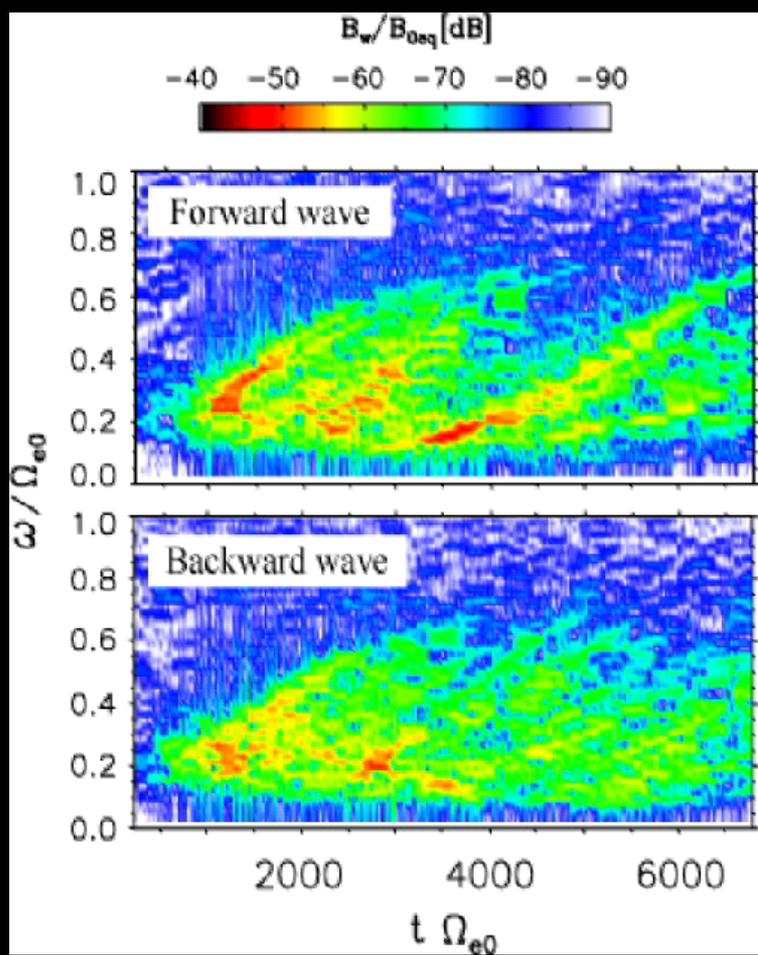
gives the **amplitude threshold** for nonlinear wave growth.

$$\tilde{\Omega}_w = \frac{B_w}{B_0} > \tilde{\Omega}_{th} = \frac{100\pi^3 \gamma^3 \xi}{\tilde{\omega} \tilde{\omega}_{ph}^4 \tilde{V}_{\perp 0}^5 \delta^5} \left( \frac{\tilde{a} s_2 \tilde{U}_{t\parallel}}{Q} \right)^2 \exp \left( \frac{\gamma^2 \tilde{V}_R^2}{\tilde{U}_{t\parallel}^2} \right)$$

The nonlinear growth rate can be greater than the linear growth rate.



[Hikishima et al., submitted to JGR]

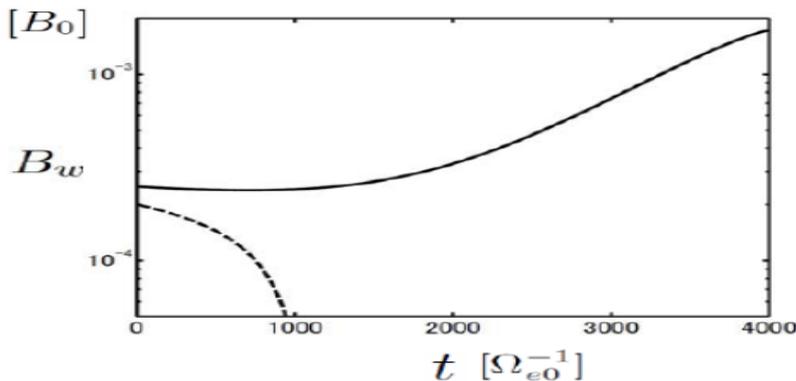
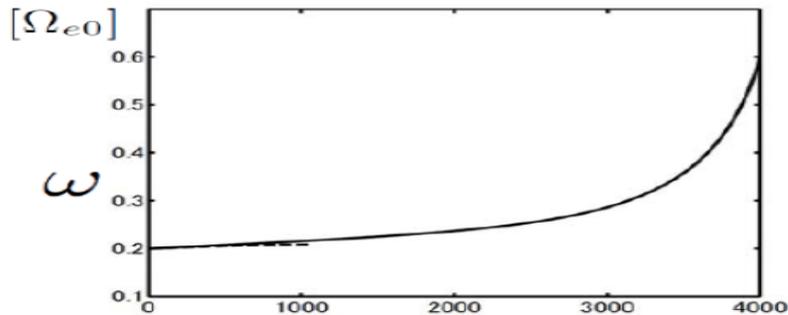


# Chorus Equations

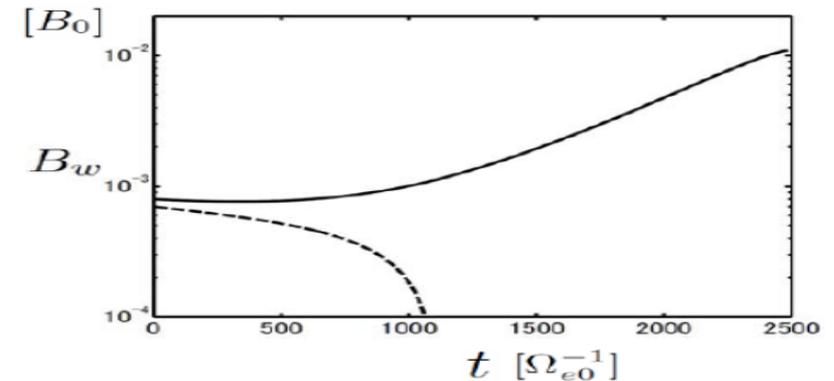
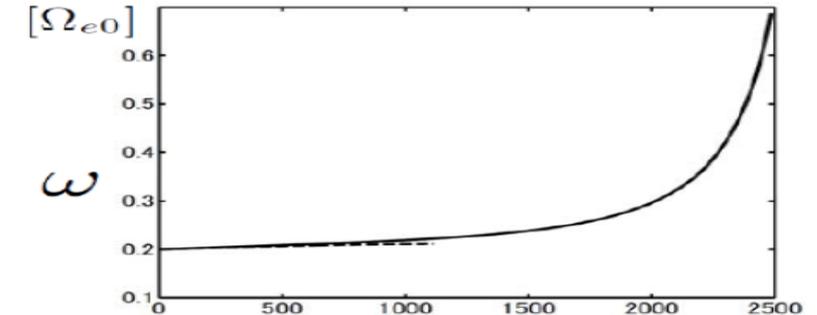
$$\frac{\partial \tilde{\Omega}_w}{\partial \tilde{t}} = \tilde{V}_g \left[ \frac{Q \tilde{\omega}_{ph}^2}{2 \tilde{U}_{t\parallel}} \left( \frac{\tilde{V}_{\perp 0} \delta}{\pi \gamma} \right)^{3/2} \left( \frac{\xi \tilde{\Omega}_w}{\tilde{\omega}} \right)^{1/2} \exp \left( -\frac{\gamma^2 \tilde{V}_R^2}{2 \tilde{U}_{t\parallel}^2} \right) - \frac{5 s_2 \tilde{a}}{s_0 \tilde{\omega}} \right]$$

$$\frac{\partial \tilde{\omega}}{\partial \tilde{t}} = \frac{2 s_0}{5 s_1} \tilde{\omega} \tilde{\Omega}_w$$

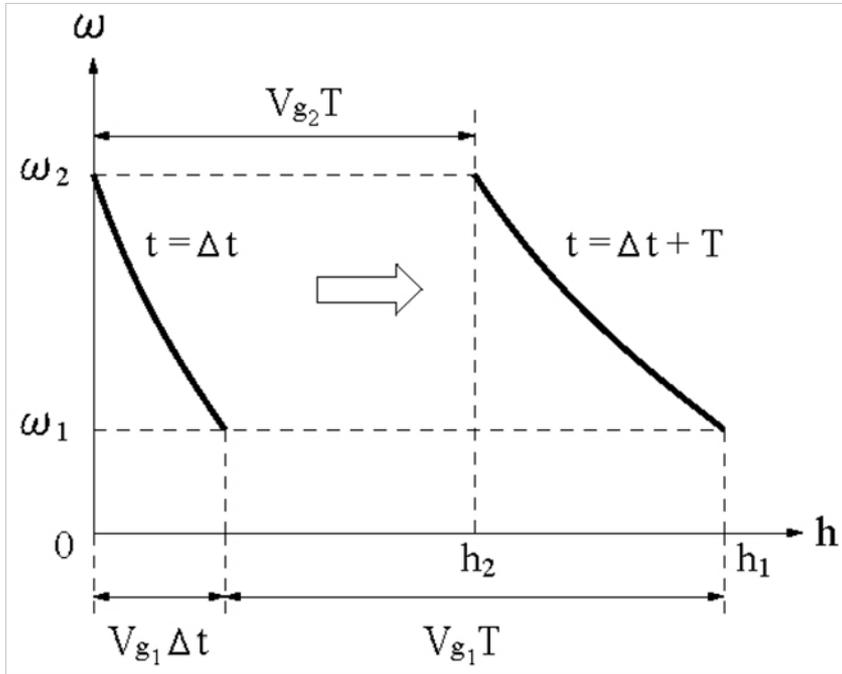
**(a) Simulation A**



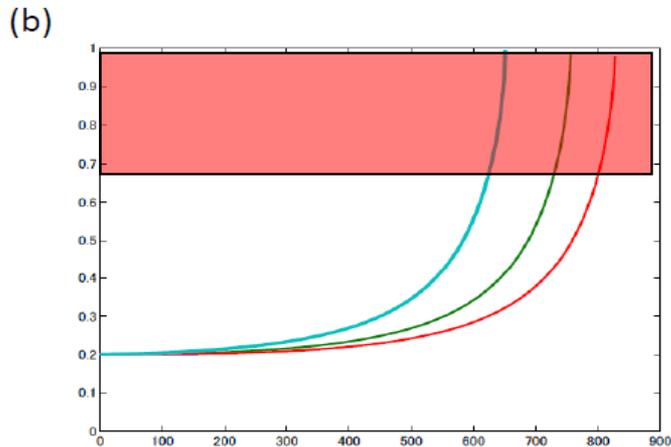
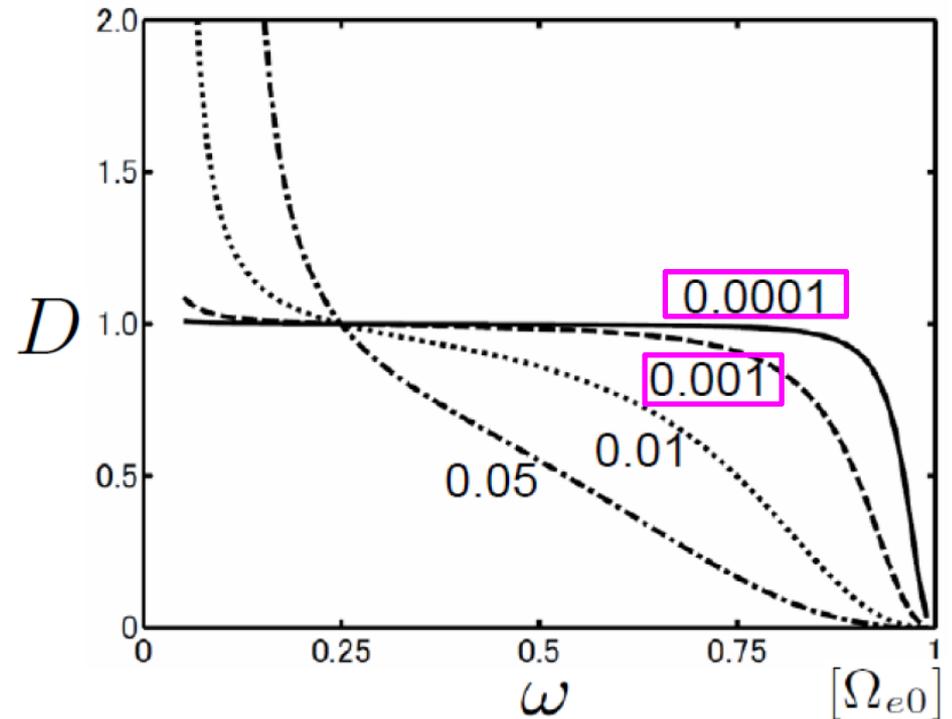
**(b) Simulation B**



# Spatial Variation of Frequency Sweep Rate (Dispersion Effect)



$$\left( \frac{\partial \omega}{\partial t} \right)_{h=h_T} = D \left( \frac{\partial \omega}{\partial t} \right)_{h=0}$$



$$D = \left[ 1 - \frac{\delta^3 (\Omega_e^2 - 4\omega\Omega_e - 4\xi^2\omega^2)}{4c\xi\omega(\Omega_e - \omega)^2} h_T \left( \frac{\partial \omega}{\partial t} \right)_{h=0} \right]^{-1}$$

# Quasi-parallel Propagation (Oblique)

$$\sin^2 \Psi \ll 1$$

$\Psi$ : Wavenormal Angle

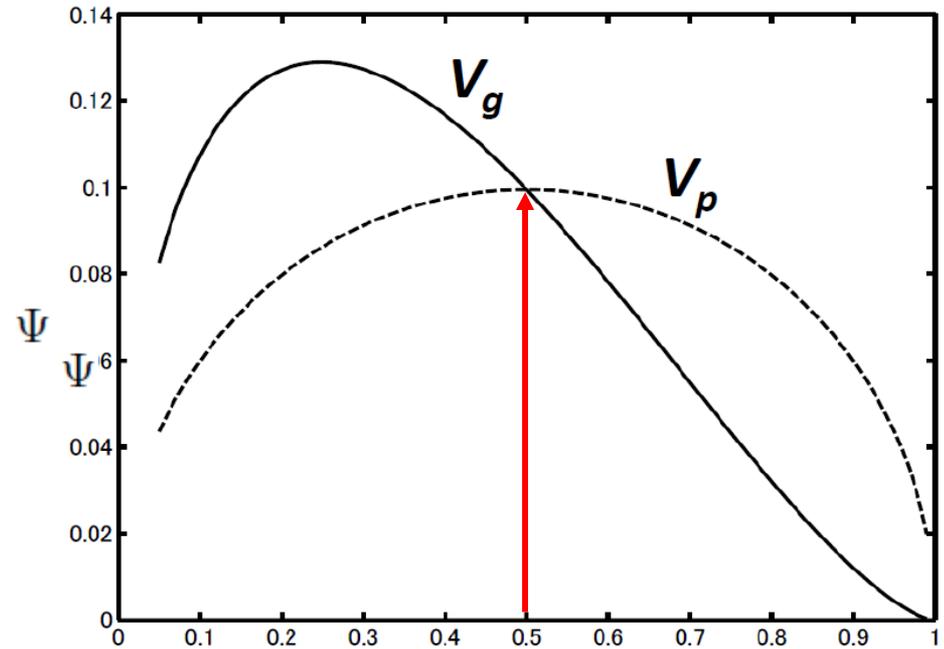
$$E_{w\parallel} = \frac{\omega \sin \Psi}{\delta^2 \Omega_e - \omega} E_w$$

With  $\omega = 0.5 \Omega_e$

$$V_g = V_p$$

$$\tilde{v}_{\parallel} = v_{\parallel} - V_p$$

$$\frac{d\tilde{v}_{\parallel}}{dt} = -\frac{eE_{w\parallel}}{\gamma m_0} \sin \phi - \frac{v_{\perp}^2}{2\Omega_e} \frac{\partial \Omega_e}{\partial h}$$



# Nonlinear Resonant Damping in the dipole magnetic field

$$\theta = -k(v_{\parallel} - V_p)$$

$$\frac{d\phi}{dt} = \theta$$

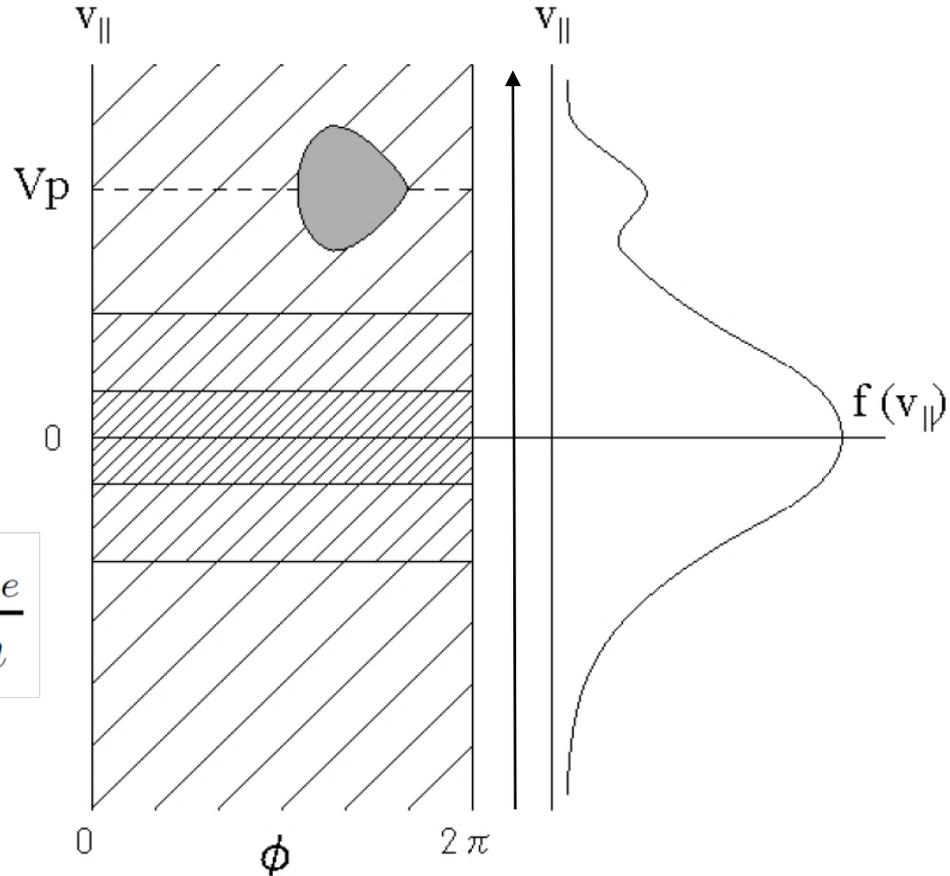
$$\frac{d\theta}{dt} = \omega_{t\parallel}^2 (\sin \phi + S_{\parallel})$$

$$\omega_{t\parallel}^2 = \frac{ekE_{w\parallel}\delta^2}{\gamma m_0}$$

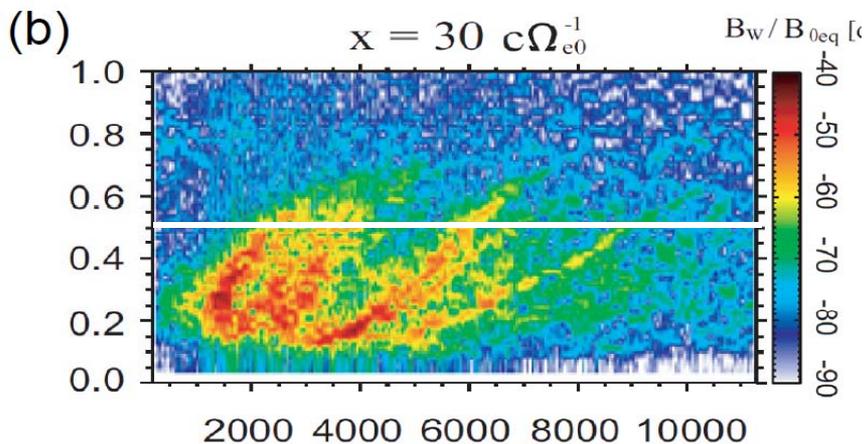
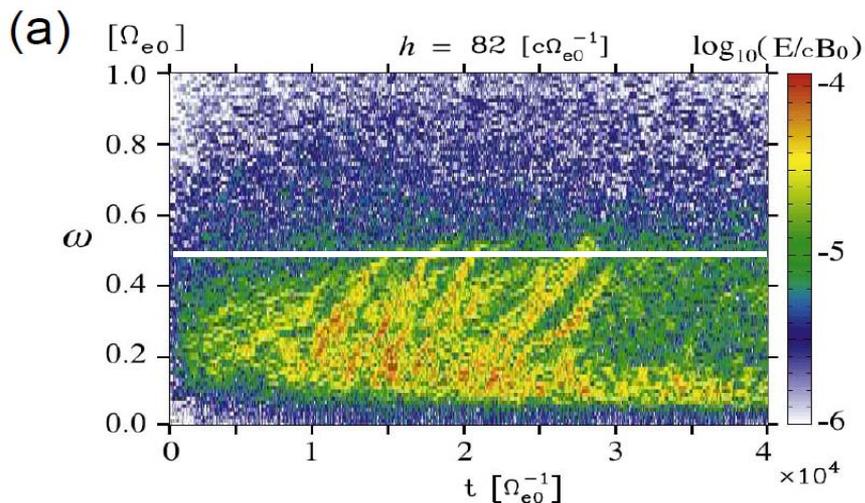
$$S_{\parallel} = \frac{kv_{\perp}^2}{2\omega_{t\parallel}^2 \Omega_e} \frac{\partial \Omega_e}{\partial h}$$

$$S_{\parallel} > 0 \quad \sin \phi_c < 0$$

$$J_{\parallel} E_{w\parallel} = -eE_{w\parallel} \int \int v_{\parallel} g_t(v_{\parallel}, \phi) \sin \phi \, d\phi dv_{\parallel} > 0$$



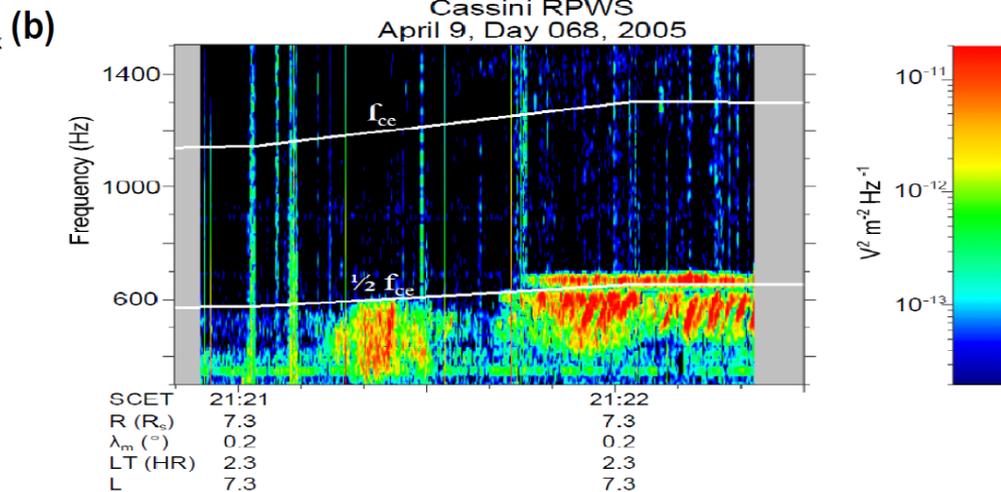
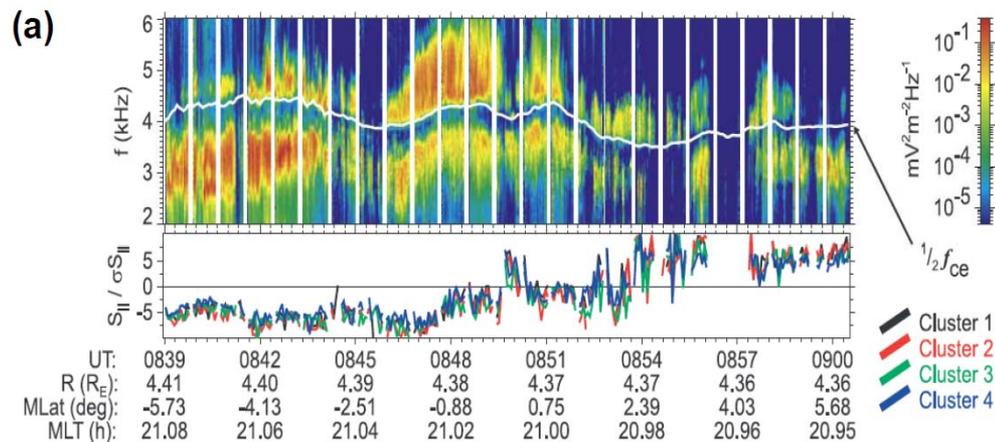
# Simulations : parallel propagation



[Omura et al., JGR, 2008]

[Hikishima et al., JGR, 2009]

# Observations: oblique propagation



[Santolik, et al., JGR, 2003]

[Hosphodarsky et al., JGR, 2008]

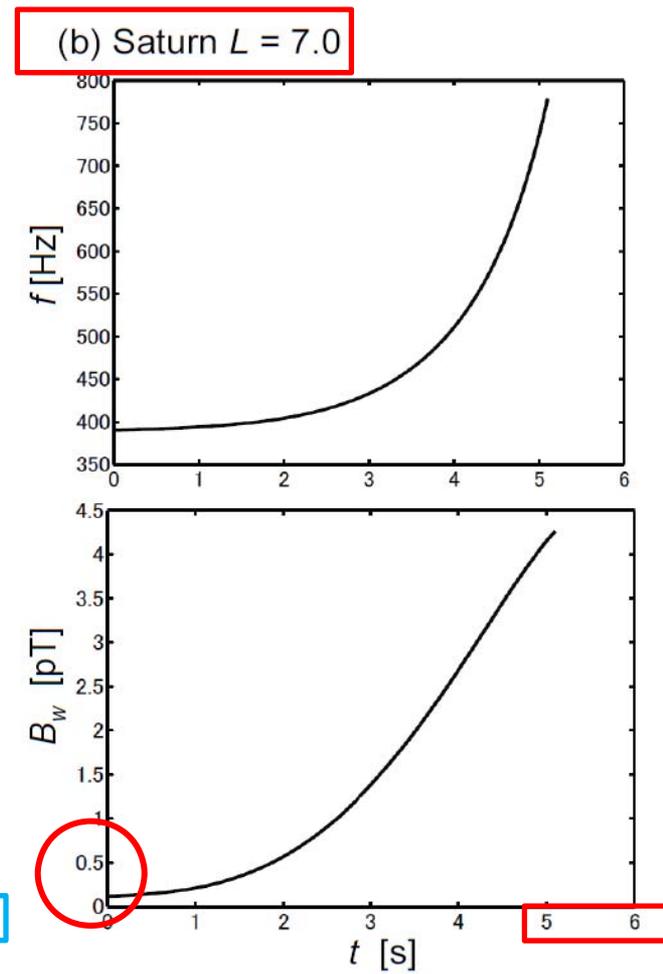
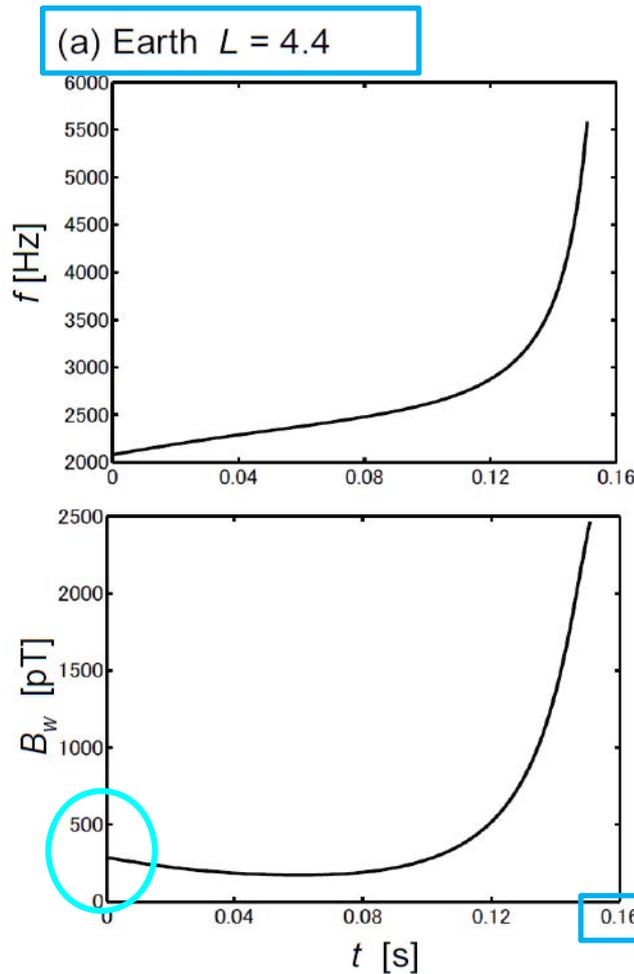
# Chorus Equations

$$\frac{\partial \tilde{\Omega}_w}{\partial \tilde{t}} = \tilde{V}_g \left[ \frac{Q \tilde{\omega}_{ph}^2}{2 \tilde{U}_{t\parallel}} \left( \frac{\tilde{V}_{\perp 0} \delta}{\pi \gamma} \right)^{3/2} \left( \frac{\xi \tilde{\Omega}_w}{\tilde{\omega}} \right)^{1/2} \exp \left( -\frac{\gamma^2 \tilde{V}_R^2}{2 \tilde{U}_{t\parallel}^2} \right) - \frac{5 s_2 \tilde{a}}{s_0 \tilde{\omega}} \right]$$

$$\frac{\partial \tilde{\omega}}{\partial \tilde{t}} = \frac{2 s_0}{5 s_1} \tilde{\omega} \tilde{\Omega}_w$$

Threshold of Wave Amplitude  $B_w$

Time Scale of Chorus  $T_c$



# Summary

- The nonlinear wave growth takes place near the equator through formation of an electromagnetic electron hole for a seed rising tone with a wave amplitude above a threshold.
- The self-sustaining mechanism of nonlinear wave growth results in a rising tone limited only by the dispersion effect operating near the gyro-frequency.
- Oblique propagation in a dipole field induces strong nonlinear resonant damping at half the gyro-frequency due to trapping by the parallel wave electric field, resulting in the lower and upper band chorus emissions.
- Resonant electrons trapped by chorus emissions are accelerated to higher pitch angles, forming a pancake distribution. Some of them are effectively accelerated to relativistic energy by RTA and URA.

# Origin of Plasmaspheric Hiss: Chorus?

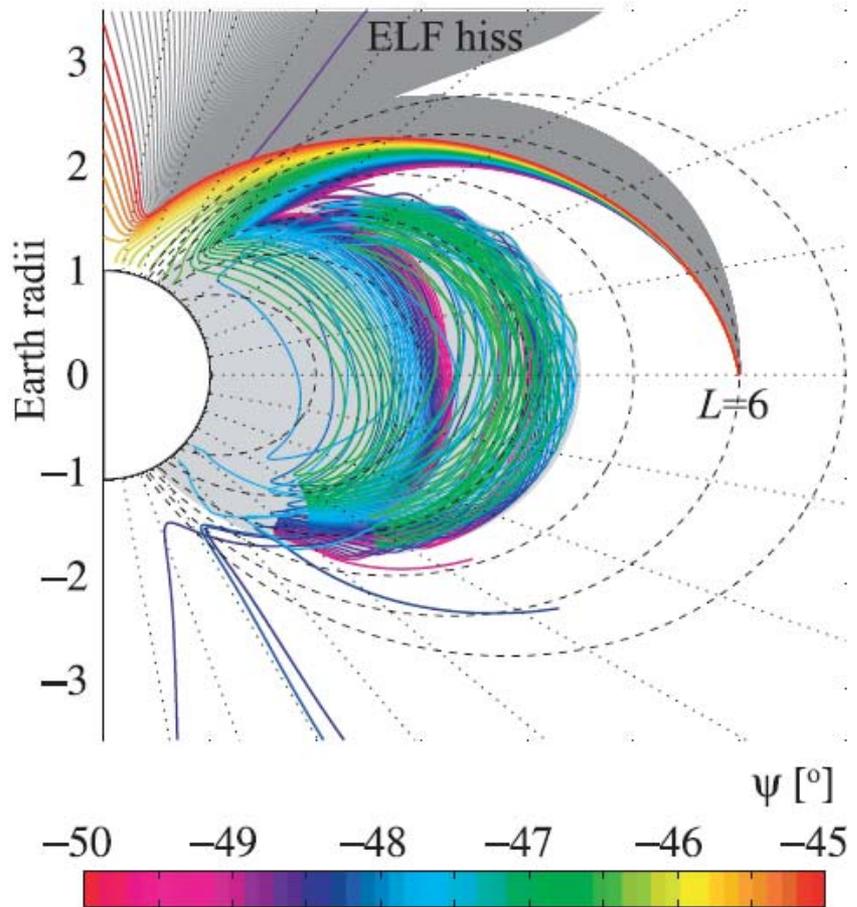
Yes, but not by direct conversion from chorus emissions propagating to high latitudes.

Chorus is generated in parallel propagation at the equator, not at wave normal angles 45 – 50 degrees.

Alternative interpretation:

Chorus emissions effectively precipitate a wide range of energetic electrons into the auroral regions. The precipitated electrons can generate the lower-hybrid waves through Landau resonance. The lower-hybrid waves propagate into the plasmasphere as whistler-mode HISS emissions.

**A study suggested by Dennis 18 years ago.**



[Bortnik et al., Science, 2009]

# Theory and Simulations on Chorus Emissions

1. Y. Omura and D. Summers, Dynamics of high-energy electrons interacting with whistler mode chorus emissions in the magnetosphere, *JGR*, Vol. 111, A09222, 2006.
2. Y. Katoh and Y. Omura, Computer simulation of chorus wave generation in the Earth's inner magnetosphere, *GRL*. Vol. 34, L03102, 2007.
3. Y. Omura, N. Furuya, D. Summers, Relativistic turning acceleration of resonant electrons by coherent whistler-mode waves in a dipole magnetic field, *Journal Geophysical Research*, Vol. 112, A06236, 2007.
4. Y. Katoh and Y. Omura, Relativistic particle acceleration in the process of whistler-mode chorus wave generation, *GRL*, 34, L13102, 2007.
5. D. Summers and Y. Omura, Ultra-relativistic acceleration of electrons in planetary magnetospheres, *GRL.*, 34, L24205, 2007.
6. Y. Omura, Y. Katoh, and D. Summers, Theory and simulation of the generation of whistler-mode chorus, *JGR*, vol. 113, A04223, 2008.
7. N. Furuya, Y. Omura, and D. Summers, Relativistic turning acceleration of radiation belt electrons by whistler mode chorus, *JGR*, vol. 113, A04224, 2008.
8. Y. Katoh, Y. Omura, and D. Summers, Rapid energization of radiation belt electrons by nonlinear wave trapping, *Ann. Geophys.*, 26, 3451, 2008
9. M. Hikishima, S. Yagitani, Y. Omura, and I. Nagano , Full particle simulation of whistler-mode rising chorus emissions in the magnetosphere, *JGR*, 114, A01203, 2009.
10. Y. Omura, M. Hikishima, Y. Katoh, D. Summers, S. Yagitani, Nonlinear mechanisms of lower band and upper band VLF chorus emissions in the magnetosphere, *JGR*, *in press*.